IFAE, Pavia, 19 Aprile 2006

Cosmologia e neutrini, con massa fissa e variabile

Marco Cirelli (Yale)

reviews:

Lesgourgues, Pastor astro-ph/0603494, Giunti et al., Phys.Rept.379 hep-ph/0211462 Dolgov, Phys.Rept.370 hep-ph/0202122 MaVaNs: Fardon, Nelson, Weiner JCAP 0410 (2004) Cirelli, Gonzalez-Garcia, Pena-Garay NPB 2005 + many others

OUTLINE

Part (1): Neutrino masses and cosmology: bird's eye view
- current bounds
- future sensitivities

Part (2): Mass Varying Neutrinos:

- the basic idea
- constraints from solar neutrino physics
- future developments

Conclusions

Light, Standard Model Neutrinos in cosmology are significant because:

- neutrinos are a lot (as abundant as photons)

- neutrinos are hot main component of the rel energy density that sets the expansion rate of the Universe
- but not so hot, they have a mass they cool down at an interesting time
- undergo matter effects in the primordial plasma
- determine neutron/proton in BBN, i.e. primordial composition of the Universe
- have energy density similar to Dark Energy

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determine neutron/proton in BBN,

- have energy density similar to Dark Energy

$m_{ u}$ affects CMB and matter spectra

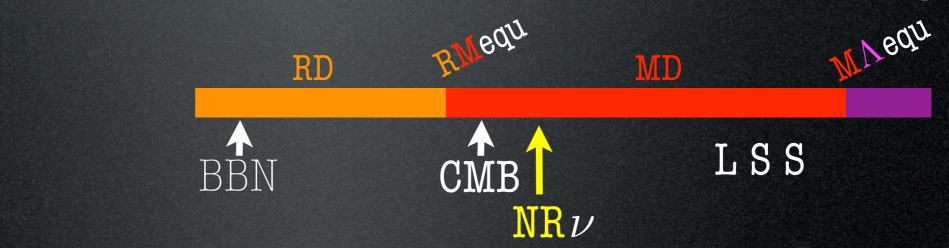
• Neutrinos become NR at

$$c \sim \frac{m_{\nu}c^2}{3 \ k \ T_{\nu,0}} \sim 2 \cdot 10^3 \left(\frac{m_{\nu}}{\mathrm{eV}}\right)$$

(CMB: $z \sim 1100$)

Since
$$m_{\nu} < 0.5 \text{ eV} \left(\sum_{i} m_{\nu_{i}} < 1.5 \text{ eV} \right)$$
,

neutrinos became NR after CMB last scattering:



- indirect effect on CMB

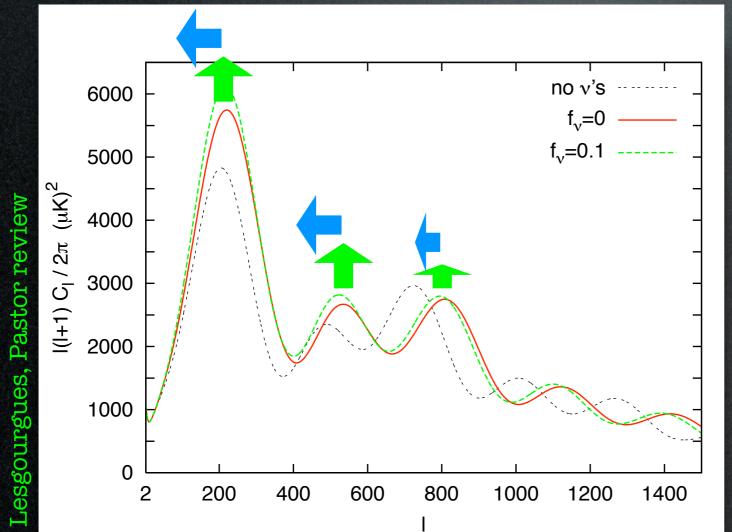
- (indirect and) direct effect on LSS and later stuff

m_{ν} affects CMB and matter spectra

On CMB power spectrum:

Contribution of neutrinos to the total energy density today $\Omega_{\nu}h^2 = \frac{\sum m_{\nu_i}}{93 \text{ eV}}$ at the expenses of other components, e.g. Ω_{CDM}

E.g. $m_{\nu} \angle A$, $\Omega_{\nu} \angle A$, $\Omega_{CDM} \searrow$, rel energy $\angle A$, RM equ delayed.



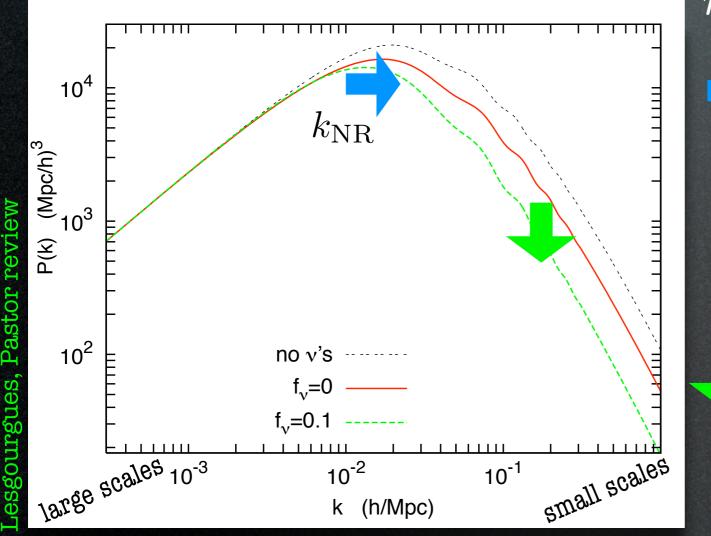
the universe "was larger" at recombination because there was more RD

less ISW effect right after recombination because grav wells decay more

Bottom line: CMB spectrum is (mildly) sensitive to $\sum m_{\nu_i}$.

$m_{ u}$ affects CMB and matter spectra

On matter power spectrum: neutrinos are not trapped (free stream), counteracting the clustering of galactic structures. Massive neutrinos become NR and travel $\lambda_{\rm FS}$ < Hubble radius. Small scales are affected.

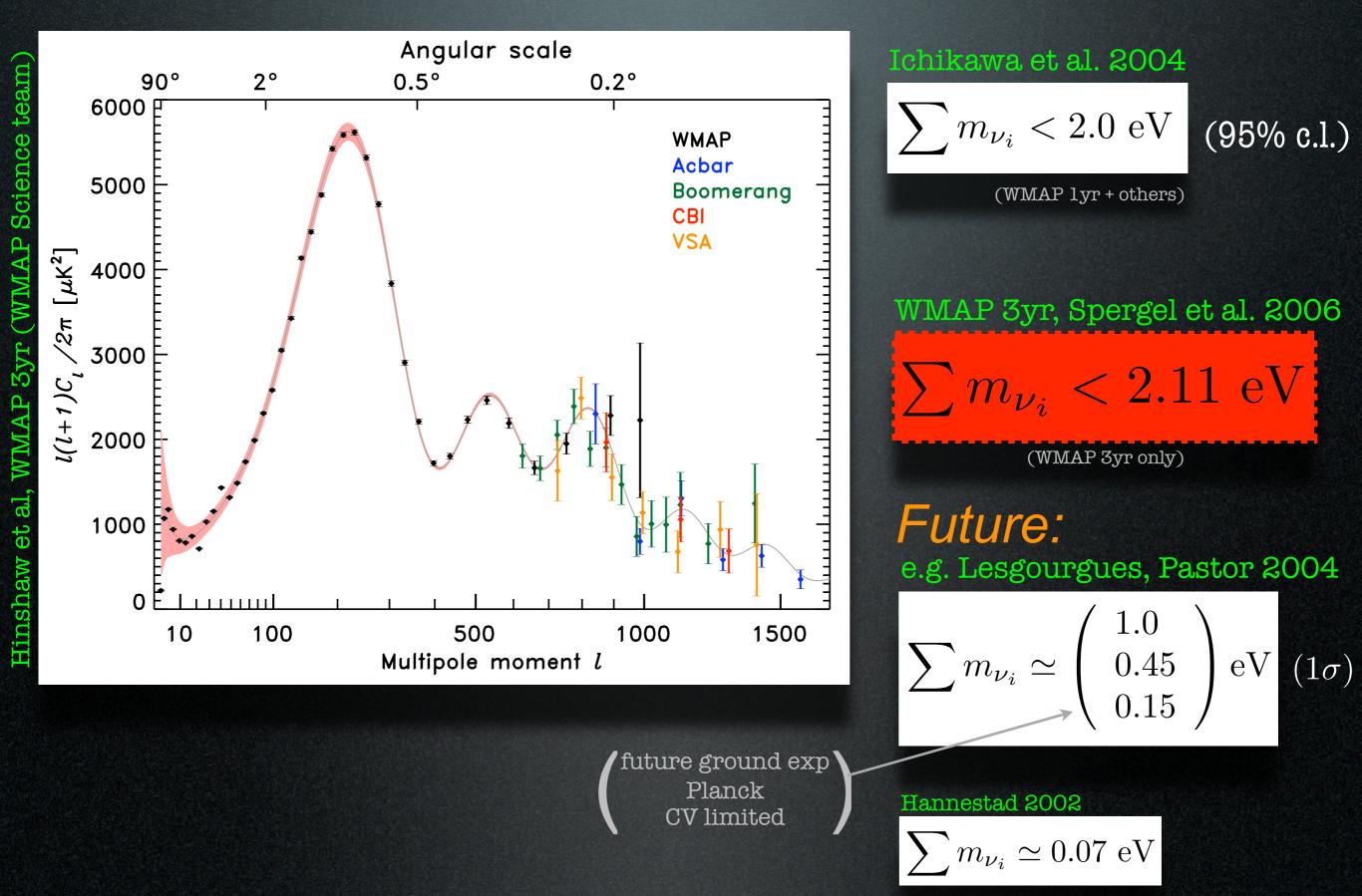


 m_{ν} determines 2 things:

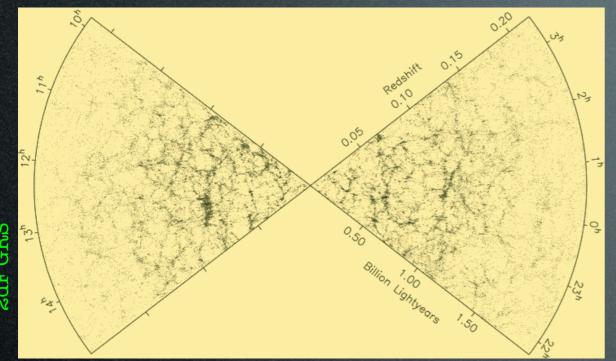
time (\leftrightarrow scale) of NR: $k_{\rm NR} = 0.03 \left(\frac{m_{\nu}}{1 {\rm eV}}\right)^{1/2} \Omega_m^{1/2} h \,{\rm Mpc}^{-1}$ $m_{\nu} \checkmark$, "slowed down" earlier, could reach smaller scales

amount of suppression: $\frac{\Delta P(k)}{P(k)} = -8 \frac{\Omega_{\nu}}{\Omega_{m}} \qquad \Omega_{\nu} h^{2} = \frac{\sum m_{\nu_{i}}}{93 \text{ eV}}$ $m_{\nu} \swarrow, \Omega_{\nu} \checkmark, \text{ suppression } \checkmark$

1.CMB alone



2.LSS galaxy surveys



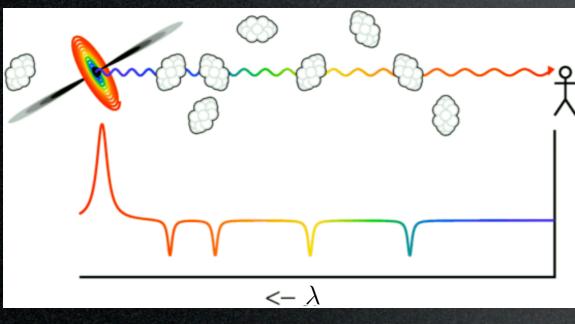
2dF: completed, 222.000 galaxies SDSS: on going, 5th data release, 1M galaxies

Redshift survey of galaxies allows to reconstruct matter distribution.

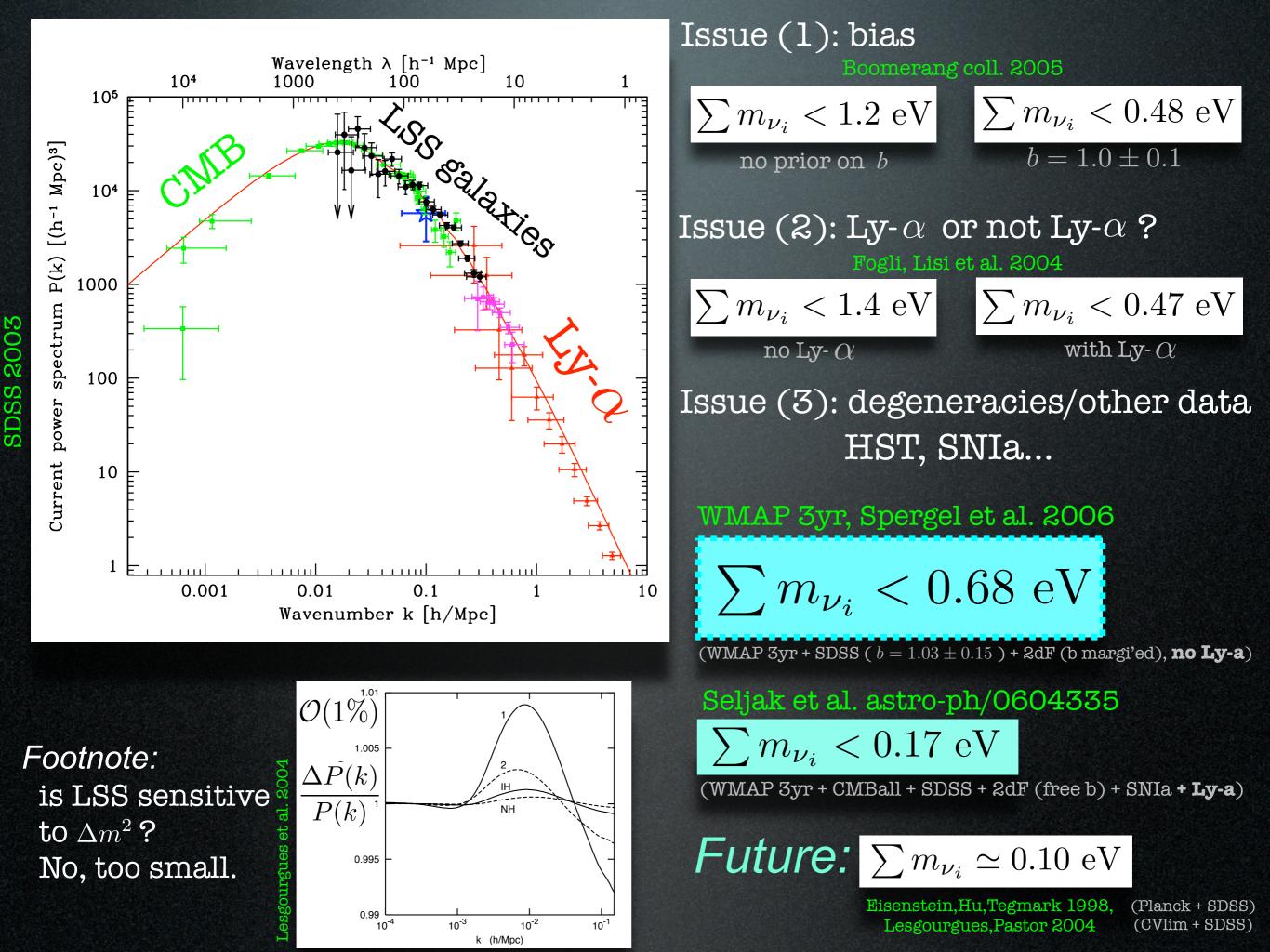
Bias parameter b: how well does light of (non-lin evolved) galaxies trace matter distribution?

Simulations and direct measurements say scale-indep and $b\approx 1.0\pm 0.1$ $_{\rm astro-ph/0406594}^{\rm Seljak~et~al.,~SDSS,}$

2b.Lyman-alpha forest



Distant quasar light, redshifted and absorbed at Ly- α frequency by intervening matter, allows to reconstruct matter z distribution along the line of sight. But: systematics and uncertainties



3.galaxy weak lensing

Weak lensing "ellipticizes" the image of background galaxies, allows to reconstruct intervening matter distribution.

Cooray, astro-ph/9904246

Future: $\sum m_{\nu_i} \simeq 0.1 \text{ eV}$

future lensing surveys: DES, SNAP...

3.CMB weak lensing

Weak lensing "distorts" the CMB, allows to reconstruct intervening matter distribution.

 $\sum m_{\nu_i} \simeq 0.03 \text{ eV}$

Song. Knox 2003

Bernardeau, astro-ph/9611012, Seljak, Zaldarriaga, astro-ph/9810257

Future: $\sum m_{\nu_i} \simeq 0.4 \text{ eV}$ (ground) $\sum m_{\nu_i} \simeq 0.15 \text{ eV}$ (Planck) $\sum m_{\nu_i} \simeq 0.044 \text{ eV}$ (CV limited)

> Kaplinghat et al. 2003. Lesgourgues, Pastor 2006

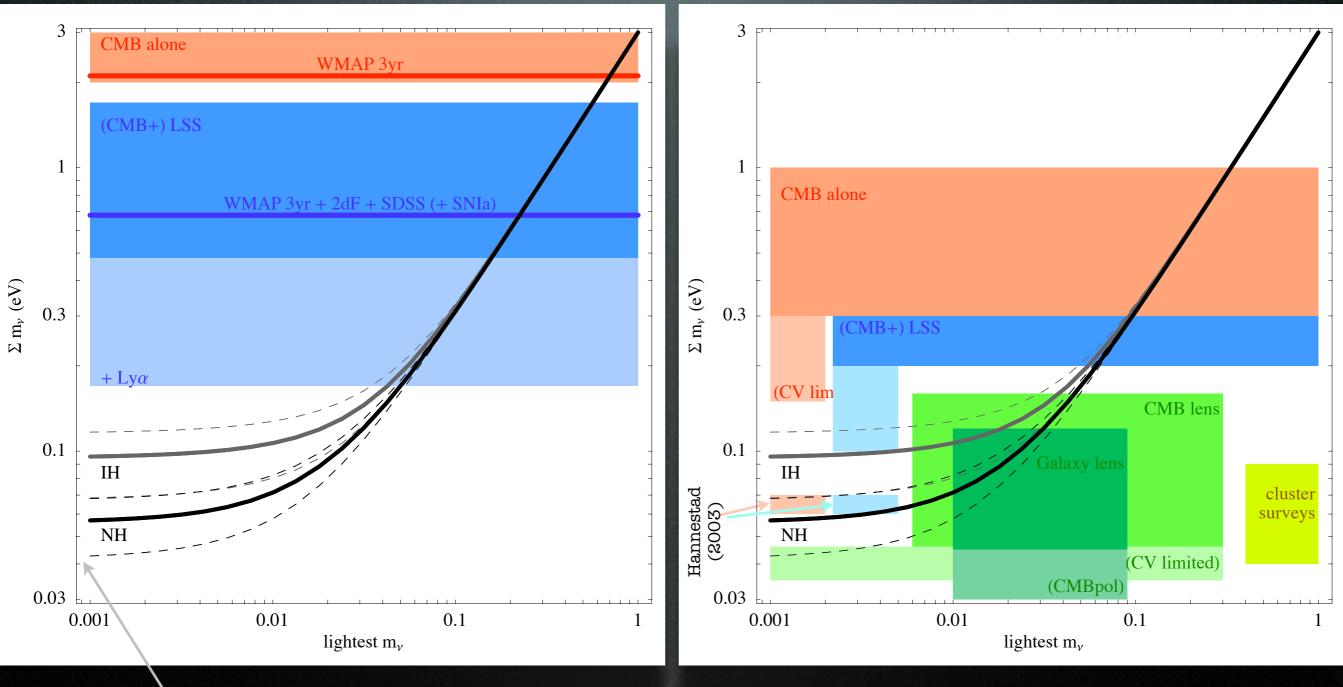
Summary

present bounds

 $m_{
m atm}^2$

22

future sensitivities



Legenda: the bound or measurement will fall somewhere in the colored box; "where it'll fall exactly" depends on the author, the experiment considered, priors, the weather...

Mass Varying Neutrinos

Fardon, Nelson, Weiner JCAP 0410 (2004) Kaplan, Nelson, Weiner PRL93 (2004)

Zurek 2004 Peccei 2005 Bi, Feng, Gu, Li, Wang, Zhang 2003-4 Cirelli, Gonzalez-Garcia, Pena-Garay 2005 Horvat 2005 Afshordi, Zaldarriaga, Kohri 2005 Takahashi, Tanimoto 2005 Fardon, Nelson, Weiner 2005 Barger, Marfatia, Whisnant 2005 Weiner, Zurek 2005 Honda, Takahashi, Tanimoto 2005 **Zhang 2005** Ichikawa, Takahashi 2005 Gu, Bi, Zhang 2005 Gonzalez-Garcia, de Holanda, Funchal 2005 Schwetz, Winter 2005 Gu, Bi, Feng, Young, Zhang 2005 Brookfield, van de Bruck, Mota, Tocchini-Valentini 2005

Mass Varying Neutrinos

Inspiration:

Fardon, Nelson, Weiner, JCAP 2004

- we don't understand Dark Energy, but it's there we don't understand neutrino mass, but it's there
- they have a similar value and they have a similar value today

$$\Omega_{\Lambda} = \frac{\rho_{de}}{\rho_c} \simeq 70\% \quad \Rightarrow \quad \rho_{de} \simeq 3 \ 10^{-11} \ \text{eV}^4 \quad \Rightarrow \quad \sqrt[4]{\rho_{de}} \simeq 2 \ 10^{-3} \ \text{eV}$$

$$m_{\nu} \simeq \sqrt{\Delta m_{atm}^2} = \sqrt{2} \ 10^{-3} \ \text{eV}^2 = 5 \ 10^{-2} \ \text{eV}$$
$$m_{\nu} \simeq \sqrt{\Delta m_{sun}^2} = \sqrt{7} \ 10^{-5} \ \text{eV}^2 = 8 \ 10^{-3} \ \text{eV}$$

$$\rho_{\nu,0} \simeq \sum m_{\nu_i} \ n_{c\nu b,i} \simeq \text{few } 10^{-14} \text{ eV}^4$$

$$m_{c\nu b,i} = \frac{3}{4} \frac{\zeta(3)}{\pi^2} g_{\nu} T^3_{\nu,0} = 8 \ 10^{-13} \ \text{eV}^3$$

 $\simeq 112 \text{cm}^{-3}$

maybe they are related maybe they "track" each other

(A different approach: Barbieri, Hall, Oliver, Strumia, 2005)

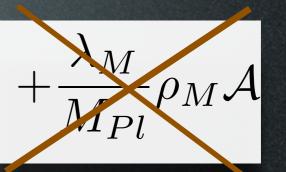
The framework:

Fardon, Nelson, Weiner, JCAP 2004

Ingredients ν_l \mathcal{A} N_r

 $m_{\nu}(\mathcal{A}) = \frac{m_D^2}{M(\mathcal{A})}$

$$\mathcal{L} = m_D \nu_l N_r + M(\mathcal{A}) N_r N_r + V_{tot}(\mathcal{A})$$



Barger et al. hep-ph/0502196

$$\mathcal{L} = m_{\nu}(\mathcal{A}) \ \nu_{l}\nu_{l} + V_{de} + V_{c\nu b} + V_{\nu,\text{medium}}$$

("see-saw")

$$V_{de} = V_{de} \left(m_{\nu}(\mathcal{A}) \right)$$

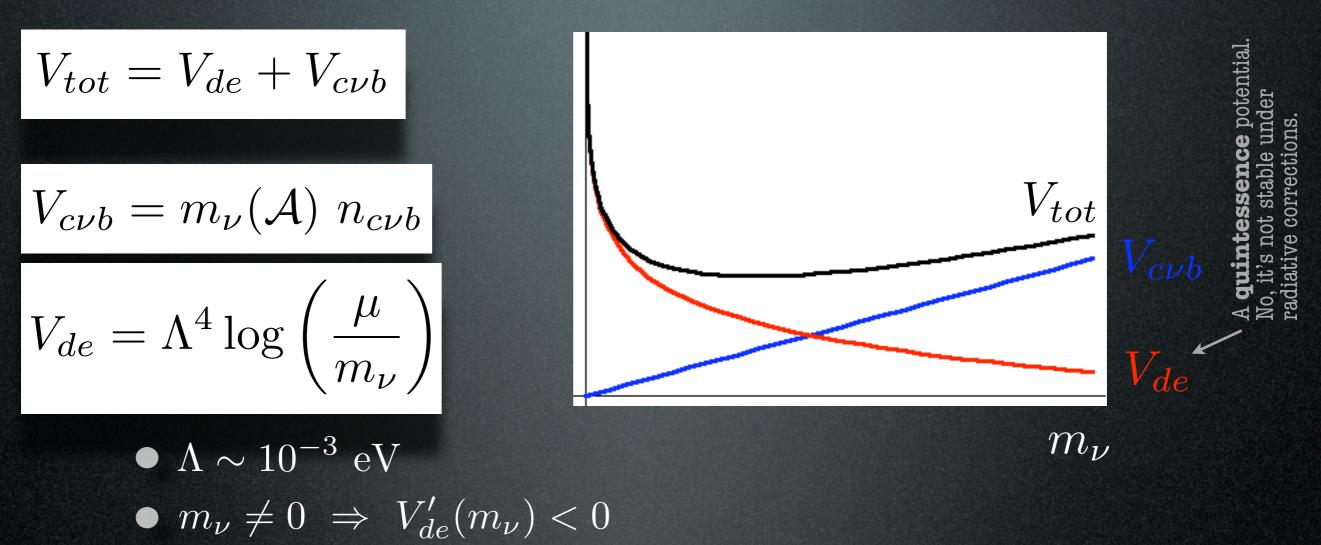
$$V_{c\nu b} = m_{\nu}(\mathcal{A}) \ n_{c\nu b}$$

$$V_{\nu,\text{medium}} = \int \frac{d^3k}{(2\pi)^3} \sqrt{k^2 + m_{\nu}^2} f$$

neutrino energy neutrino distribution fnct

 $\dot{k}(k)$

In the "vacuum":



• $\omega \approx -1 \Rightarrow |V'_{de}(m_{\nu})| \ll 1 \text{ (flat DE potential)} \quad (\omega = -0.97^{+0.07}_{-0.09})$

WMAP 3yr

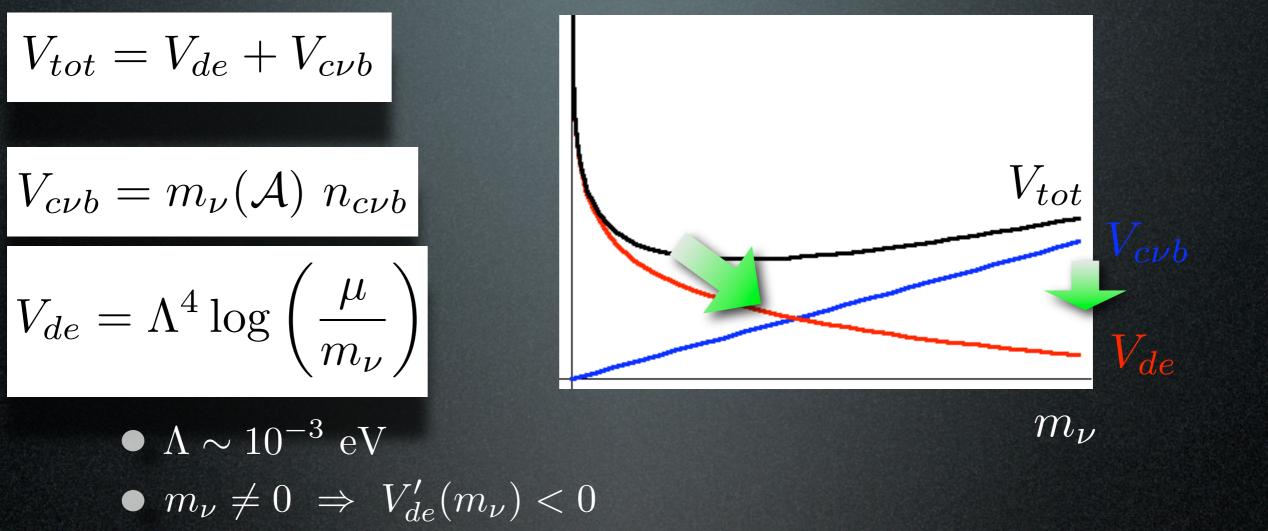
mass

varying!

minimization of the total potential

$$\frac{dV_{tot}(m_{\nu})}{dm_{\nu}} \equiv 0 \Rightarrow \qquad m_{\nu,0} = \frac{\Lambda^4}{n_{c\nu b}} \checkmark$$

In the "vacuum":



• $\omega \approx -1 \Rightarrow |V'_{de}(m_{\nu})| \ll 1$ (flat DE potential) $(\omega = -0.97^{+0.07}_{-0.09})$

WMAP 3yr

mass

minimization of the total potential

$$\frac{dV_{tot}(m_{\nu})}{dm_{\nu}} \equiv 0 \Rightarrow m_{\nu,0} = \frac{\Lambda^4}{n_{c\nu b}} \underset{\text{varying!}}{\text{mass}}$$

In a neutrino rich medium:

Cirelli, Gonzalez-Garcia, Pena-Garay NPB 2005

$$V_{tot} = V_{de} + V_{c\nu b} + V_{\nu,\text{medium}}$$

$$V_{\nu,\text{medium}} = \int \frac{d^3k}{(2\pi)^3} \sqrt{k^2 + m_{\nu}^2} f(k)$$

 $\frac{\text{minimization of the total potential}}{dV_{tot}(m_{\nu})} \equiv 0 \implies V'_{de}(m_{\nu}) + n_{c\nu b} + m_{\nu} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{k^2 + m_{\nu}^2}} f(k) \equiv 0$

$$m_{\nu} = m_{\nu,0} - A \ m_{\nu,0}^2 + \dots - \frac{-if A}{-if A}$$

• if
$$A \to 0$$
, back to vacuum case
• if $A \sim \mathcal{O}(1)$...

$$A = \frac{1}{n_{c\nu b}} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{k^2 + m_{\nu}^2}} f(k) \approx \frac{1}{n_{c\nu b}} \frac{n_{\nu,\text{medium}}}{\langle E_{\nu} \rangle}$$

In the Sun:

given solar $\,
u_e \,$ spectrum and prod regions, compute effective Δm^2

$$\Delta m_{\text{MaVaN}}^2(x) = m_2^2(x) - m_1^2(x)$$

$$\simeq \Delta m_0^2 \left[1 - 3A_2(x)m_{01} \right] + 2 \left[A_1(x) - A_2(x) \right] m_{01}^3 + \dots$$

effective Δm^2 is a function of m_{01} !

Solar + KamLAND fit: results

Cirelli, Gonzalez-Garcia, Pena-Garay NPB 2005

-fit worsens with m_{01} (best fit for $m_{01} = 0$)

-Solar only:

-moves to lower $\Delta m^2_{21,0}$ -but D/N asymmetry -and CC/NC ratio

-Solar + KamLAND: -KamLAND nails $\Delta m^2_{21,0}$ high



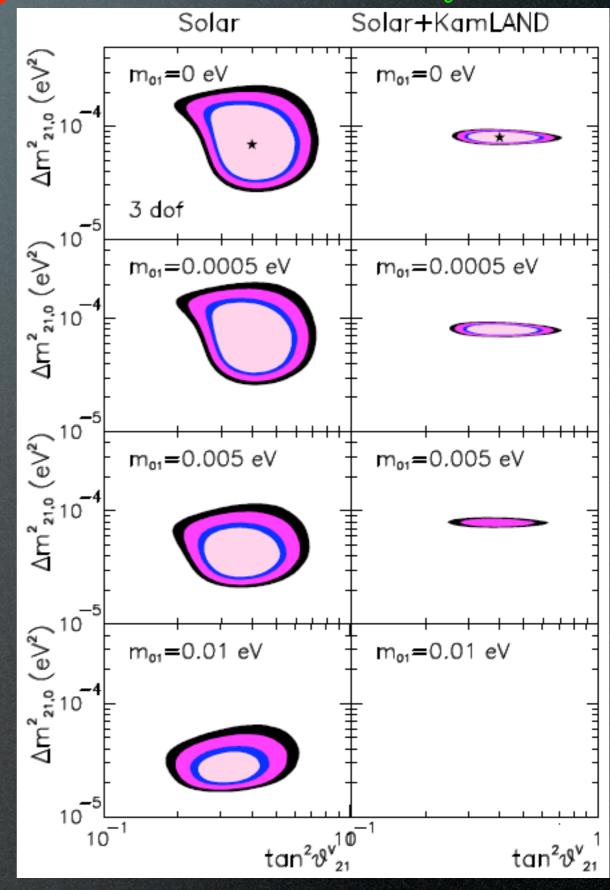
upper bound on m_{01} :

solar only

 $m_{01} \lesssim 0.05 \ \mathrm{eV}$

solar + KamLAND

 $m_{01} \lesssim 0.01 \, \text{eV}$



[more results]

Outlook:

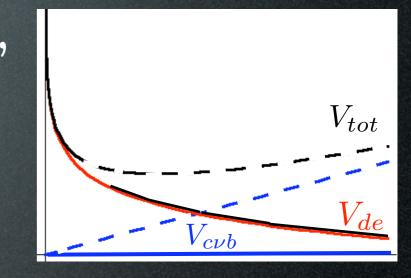
MaVaNs clustering:

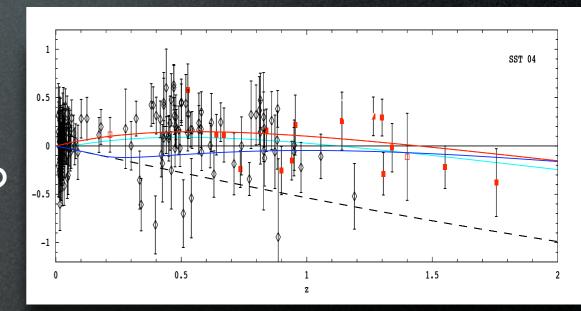
when ν become NR, instabilities collapse, neutrino CDM-like nuggets form.

- connection with $\, m_{
 u}^{
 m today} \, \, {
 m lost}$
- DE disappears (SNIa data uproar)

alternatively: -only the lightest neutrino, Parlon, still R, is coupled to DE? Note -clustering occurred yesterday? (indistinguishable from Λ ?)

Afshordi, Zaldarriaga, Kohri 2005





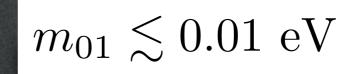
Couple to ordinary matter, environmental mass:

- does not reconcile LSND
- effects on solar and reactor and accelerator neutrinos...

Conclusions

- the bound from cosmology is the dominant bound on m_{ν} : CMB only (CMB +) LSS + Ly- α
 - $\sum m_{\nu_i} < 2.11 \text{ eV}$ $\sum m_{\nu_i} < 0.68 \text{ eV}$ $\sum m_{\nu_i} < 0.17 \text{ eV}$
- future improvements likely "guarantee" positive detection (e.g. lensing surveys)
- Mass Varying Neutrinos models aim to link fruitfully neutrinos and DE: work in progress

solar + KamLAND physics imply



Extra slides

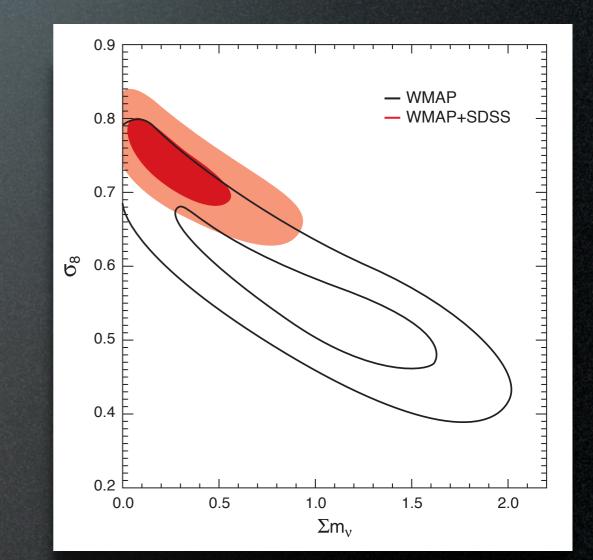
LSS and degeneracies

m_{ν} effect can be cancelled by w < -1. (SNIa data allow less Ω_{Λ} , hence more Ω_{m} , if w < -1; more Ω_{m} brings back up the P(k))

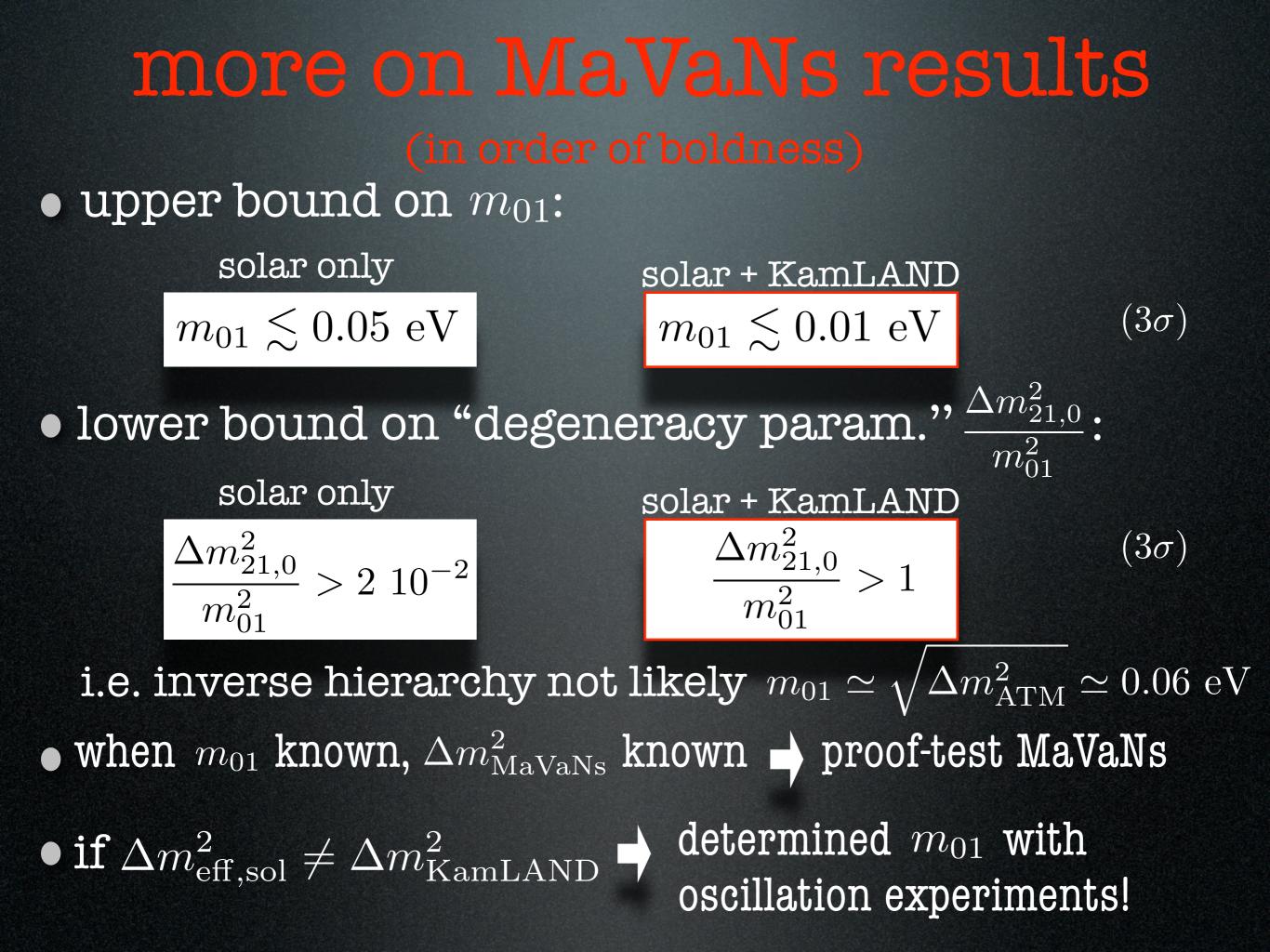
 $-0.8 \\ -1.0 \\ -1.4 \\ -1.4 \\ -1.6 \\ 0.0 \\ 0.5 \\ \Sigma m_{v}$

Still, $\sum m_{\nu_i} \lesssim 1.0 \text{ eV}$.

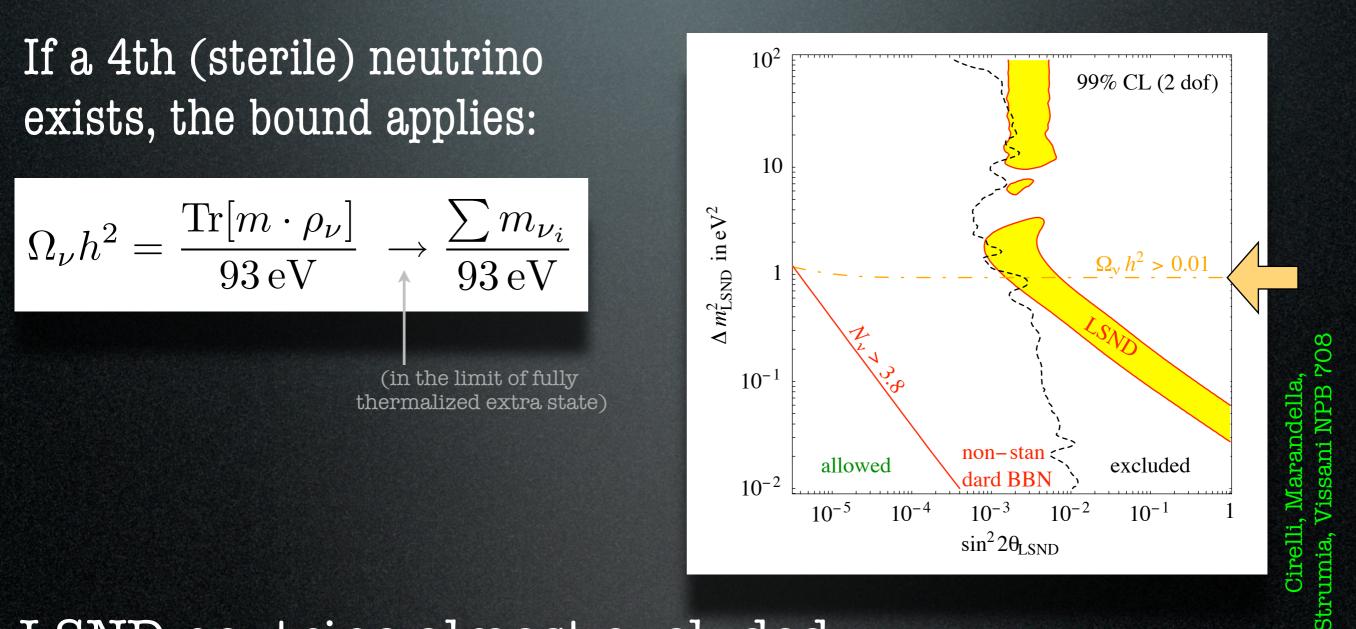
 m_{ν} effect can be cancelled by low σ_8 .



[back to LSS]



LSND sterile and the cosmological mass bound



LSND neutrino almost excluded.

Neutrinos cosmology in one sentence! Wow...

Nature provides for three types of neutrinos, yet scientists know very little about these "ghost particles," which can traverse the entire Earth without interacting with matter. But the abundance of neutrinos in the universe, produced by stars and nuclear processes, may explain how galaxies formed and why antimatter has disappeared. Ultimately, these elusive particles may explain the origin of the neutrons, protons and electrons that make up all the matter in the world around us.

Reactors BBN

??

LSS

CnuB

Leptogenesis

Sun, SNe

MINOS press release, march 2006