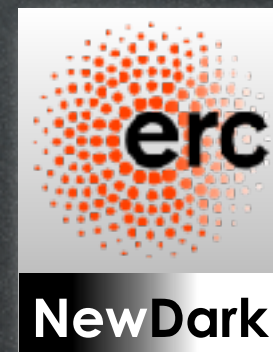


13 November 2017
DESY Theory, Hamburg

*How we fell in love with WIMPs
and should not dump them (yet)*

Marco Cirelli
(CNRS LPTHE Jussieu Paris)

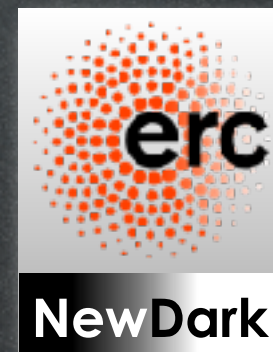


Based on: Cirelli, Fornengo, Strumia 'Minimal Dark Matter', NPB 2006 +...
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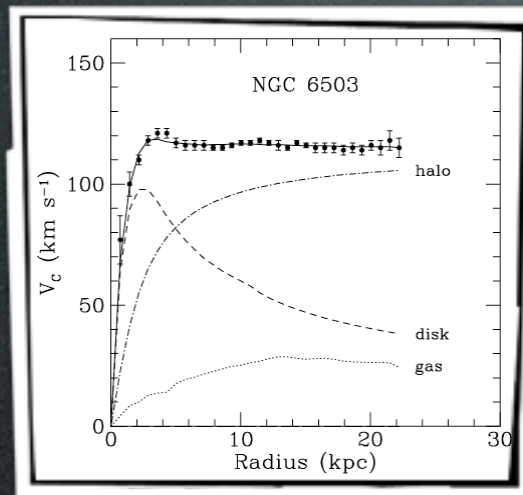
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Introduction

DM exists

Introduction

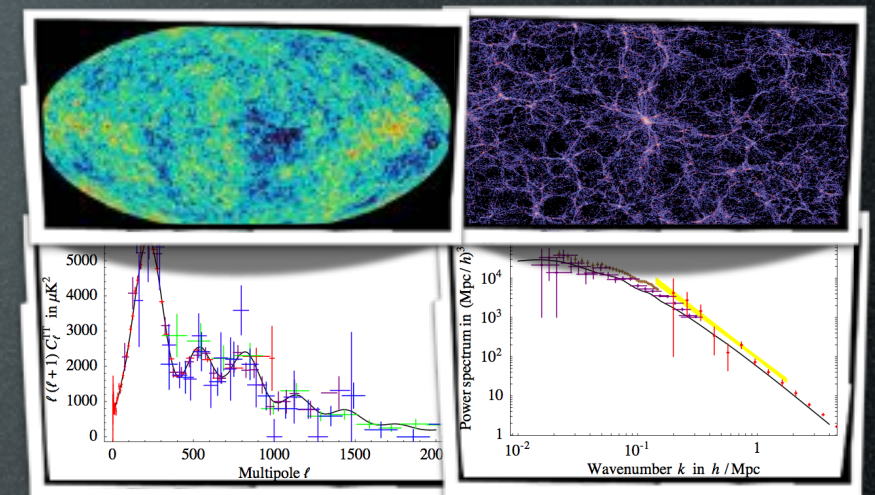
DM exists



galactic rotation curves



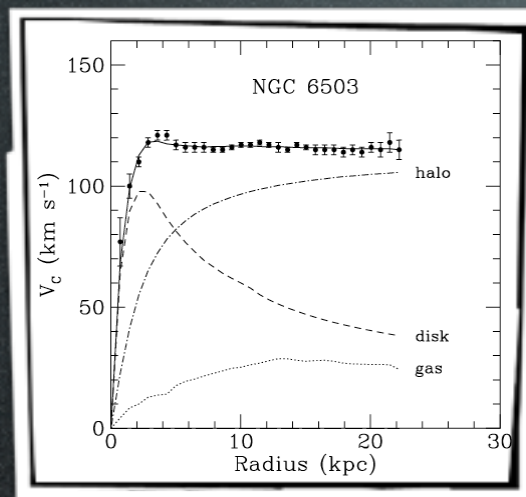
weak lensing (e.g. in clusters)



'precision cosmology' (CMB, LSS)

Introduction

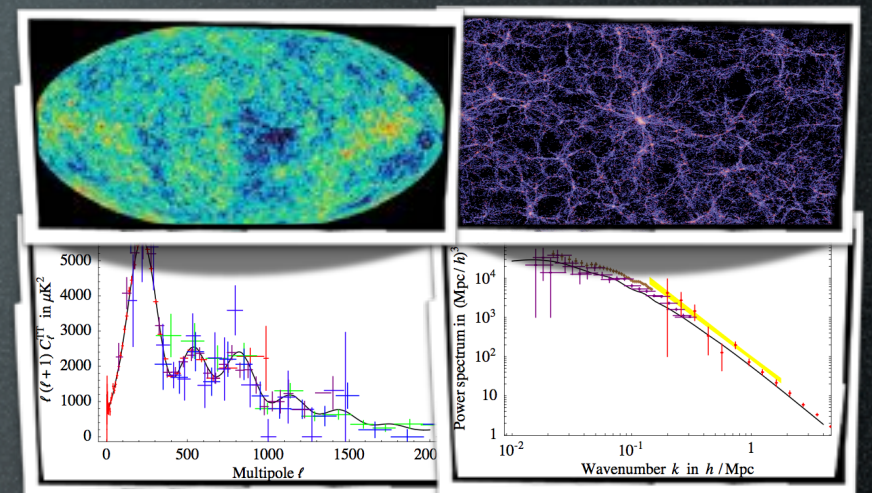
DM **exists**



galactic rotation curves



weak lensing (e.g. in clusters)



‘precision cosmology’ (CMB, LSS)

But: **what** is it?

Candidates

WIMPs

Candidates

WIMPs

$SU(2)_L$



Candidates

new physics at
the TeV scale



thermal
freeze-out



WIMPs

Candidates

new physics at
the TeV scale

thermal
freeze-out

WIMPs

LHC

AMS, Fermi, CTA
Antares, Icecube

Direct
Detection



Candidates

new physics at
the TeV scale

thermal
freeze-out

WIMPs

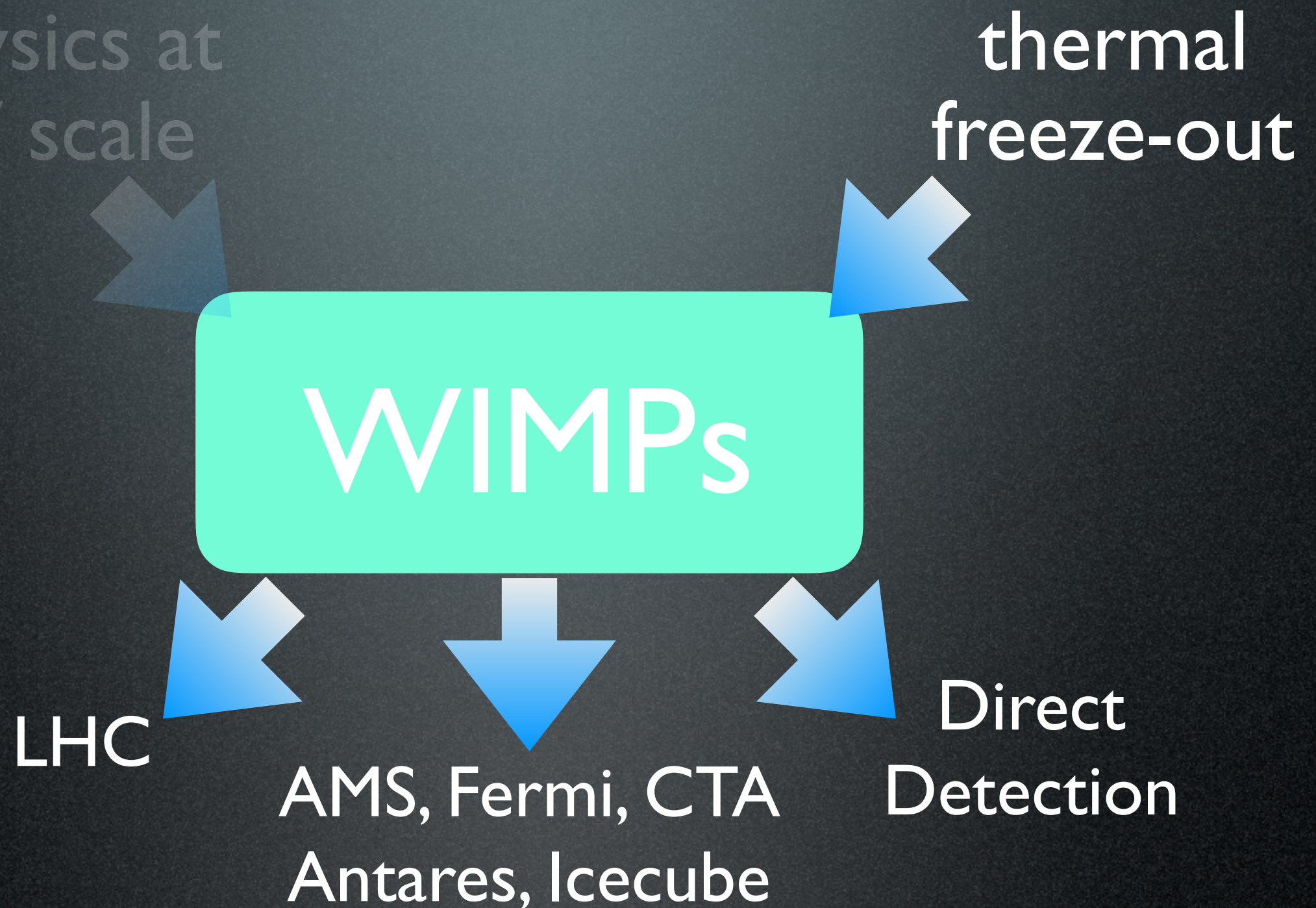
LHC

AMS, Fermi, CTA
Antares, Icecube

Direct
Detection

- 1.
- 2.
- 3.

Candidates



1. even without a larger framework, WIMPs are **still appealing**
- 2.
- 3.

Candidates

new physics at
the TeV scale



thermal
freeze-out

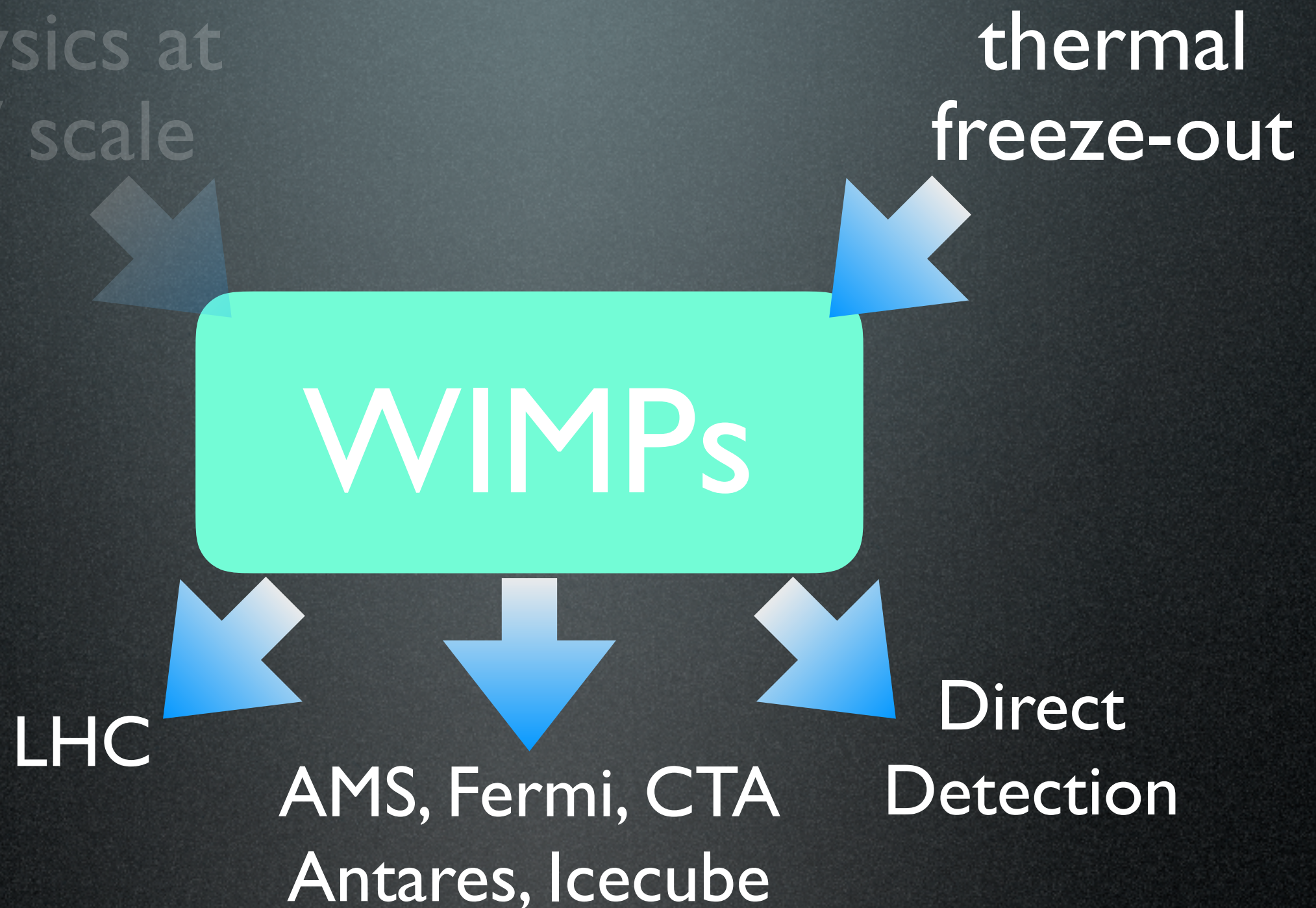


WIMPs

Aside:

Asking ‘Is New Physics at the TeV (still) exciting?’ is a bit like asking ‘Are french movies nice?’. If consistently, over many years, the large majority of movies from France are nice, this is strong motivation to build a new movie theatre and invest a lot in them. If, after many years, most french movies turn out to be horribly boring and self-centered, it’s OK to lose interest in them. That does not imply that there can be one beautiful french movie one day that wins the Oscar (or a point in the SuSy parameter space that turns out to be The Truth.)

Candidates



1. even without a larger framework, WIMPs are **still appealing**
2. the frontier is **multi-TeV**
- 3.

Candidates

new physics at
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thermal
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WIMPs

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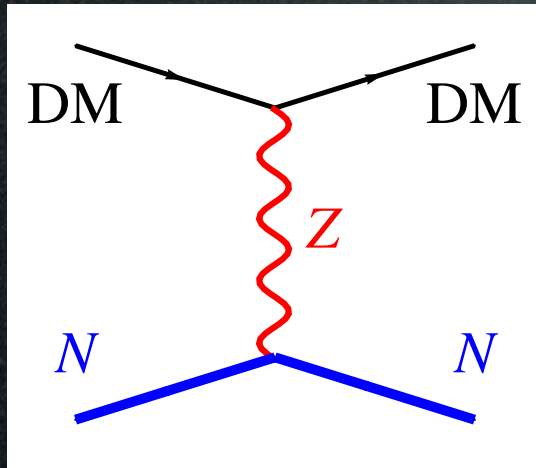
1. even without a larger framework, WIMPs are **still appealing**
2. the frontier is **multi-TeV**
3. searches are **complementary** and still have **ground to cover**

WIMP DD: **'theory'**

SM weak scale SI interactions

WIMP DD: 'theory'

SM weak scale SI interactions

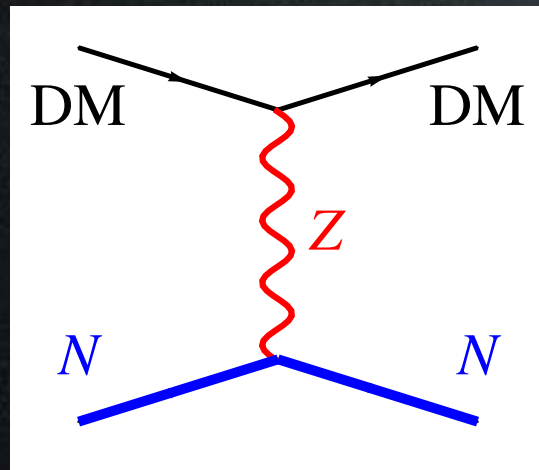


tree level,
vector

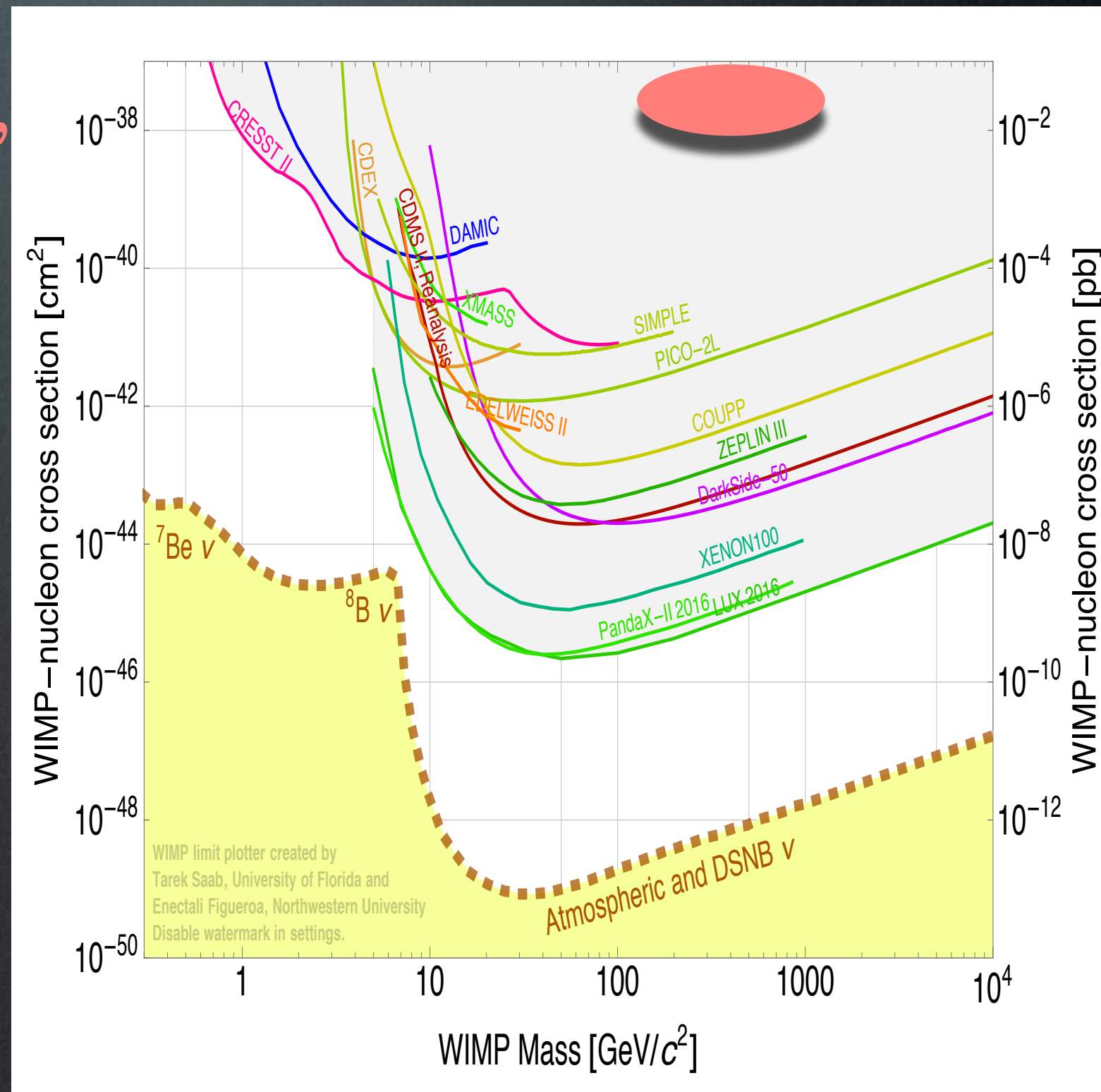
$$\sigma_{\text{SI}} \sim \frac{\alpha^2 m_N^2}{M_Z^4}$$

WIMP DD: 'theory'

SM weak scale SI interactions

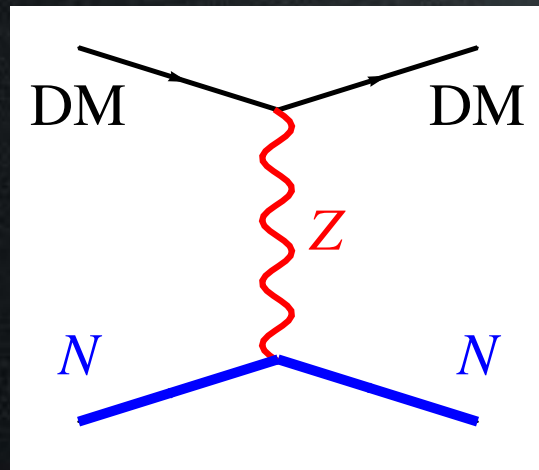


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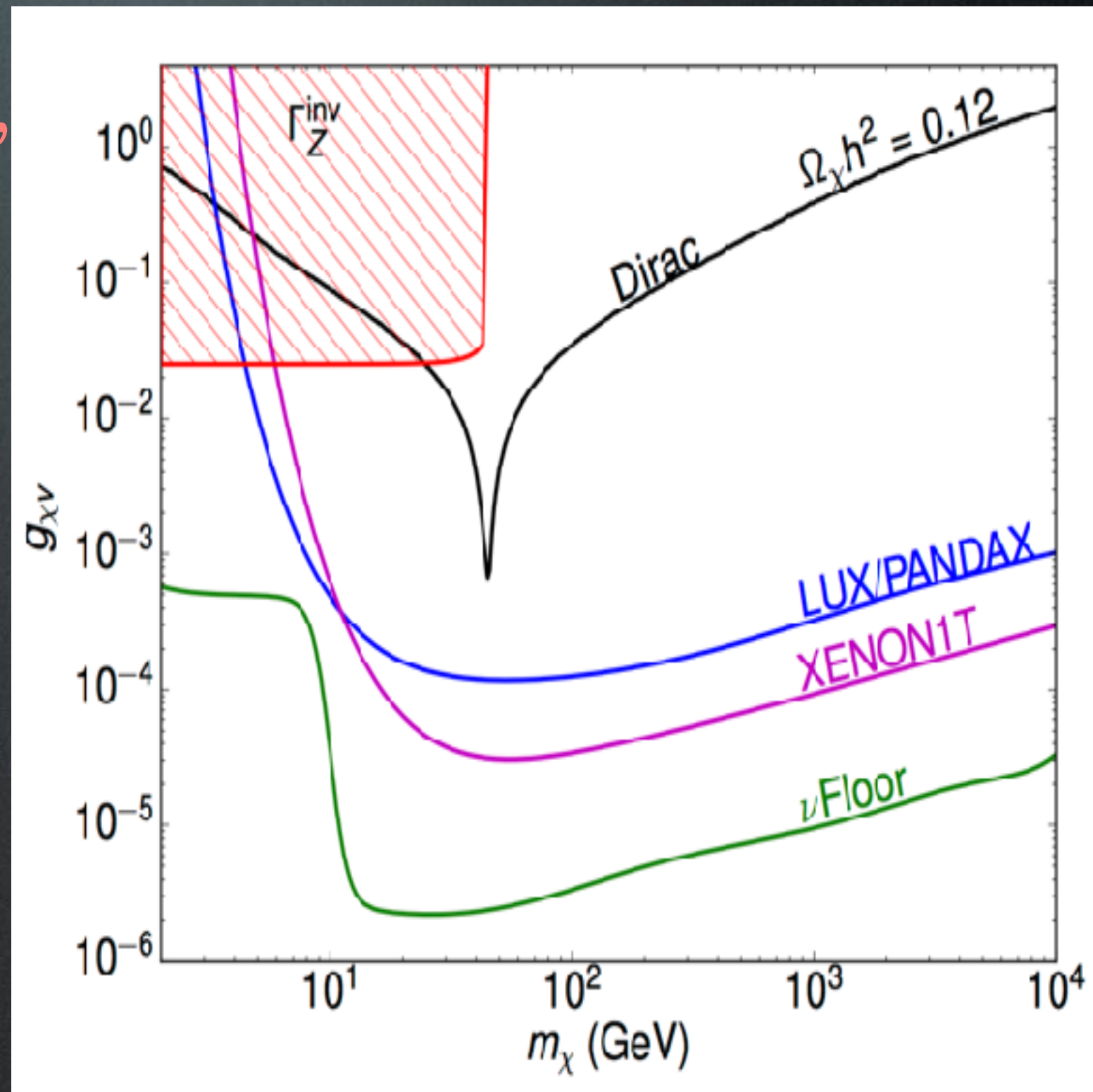


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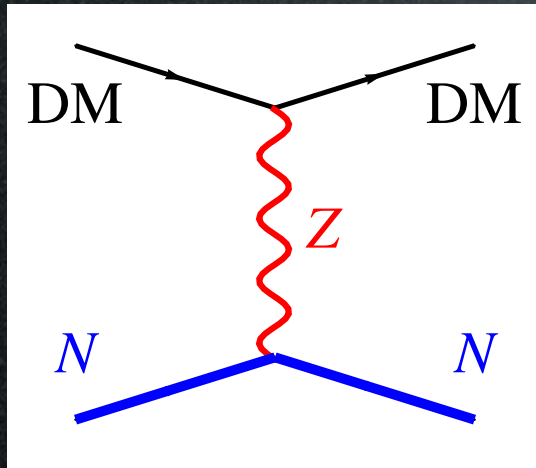


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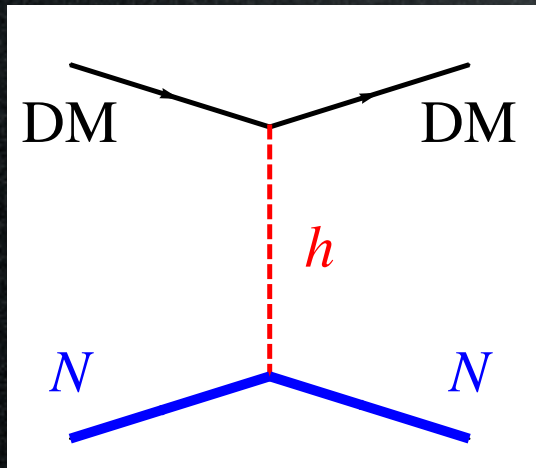
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$$\sigma_{\text{SI}} \sim \frac{\alpha^2 m_N^2}{M_Z^4}$$

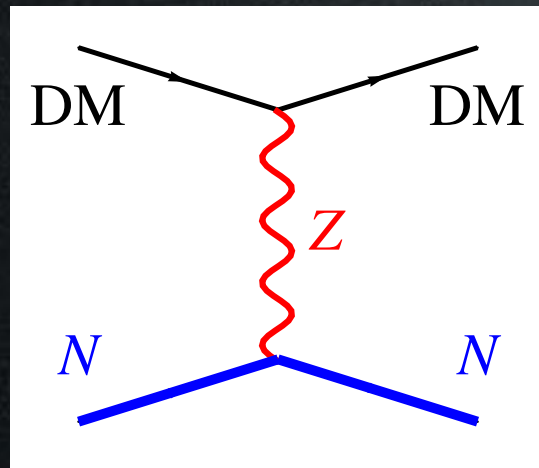


tree level,
scalar

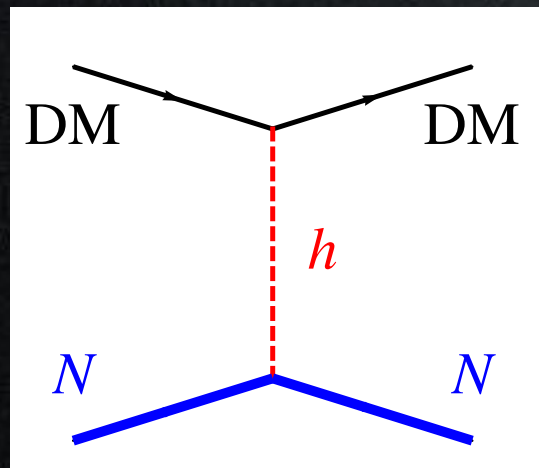
$$\sigma_{\text{SI}} \sim \frac{\alpha^2 m_N^4}{M_h^6}$$

WIMP DD: 'theory'

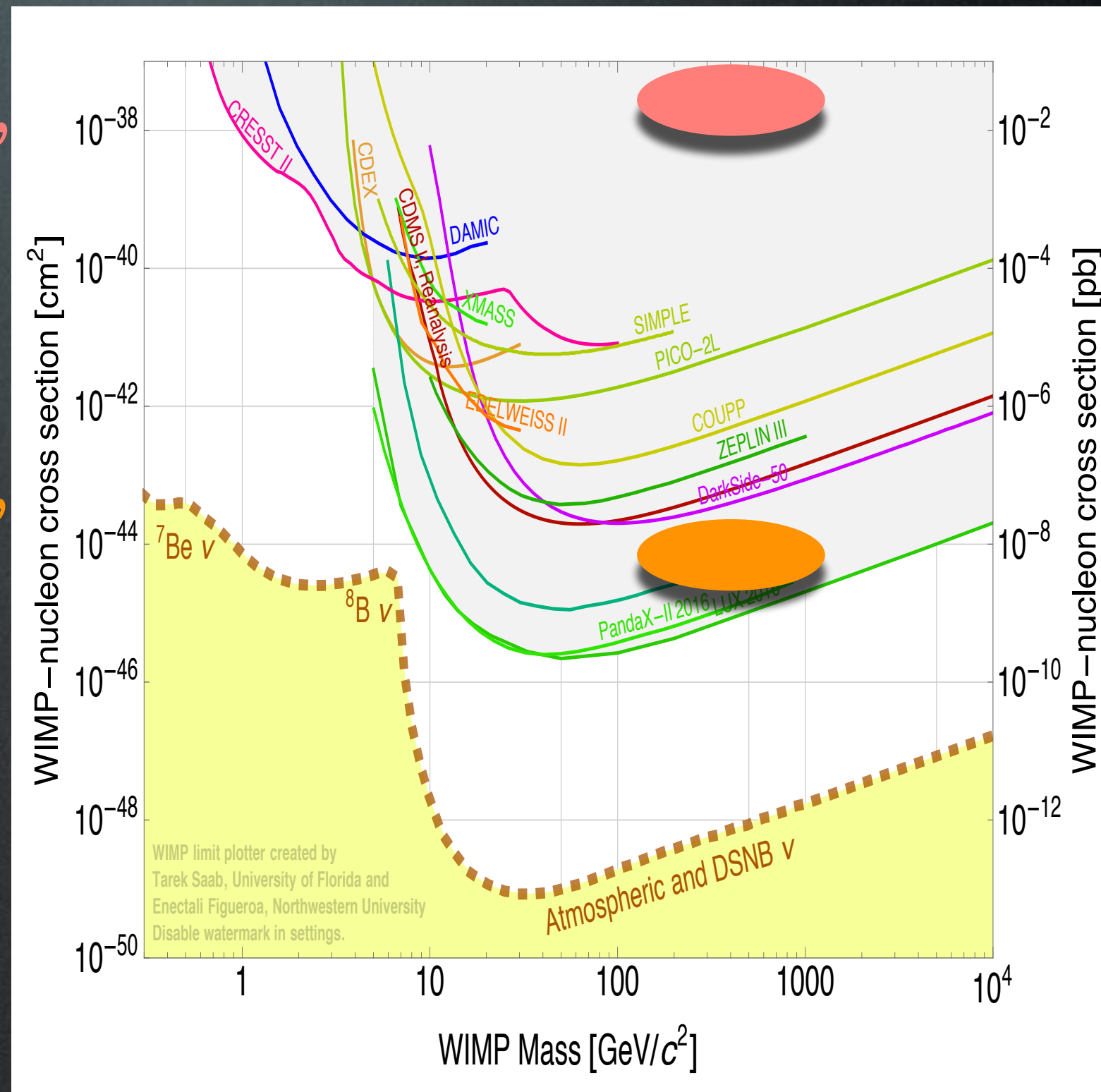
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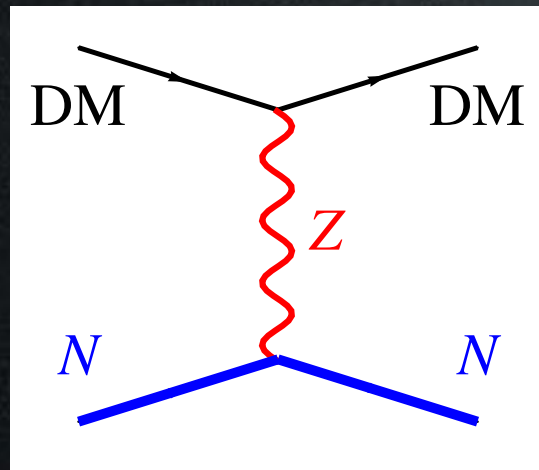


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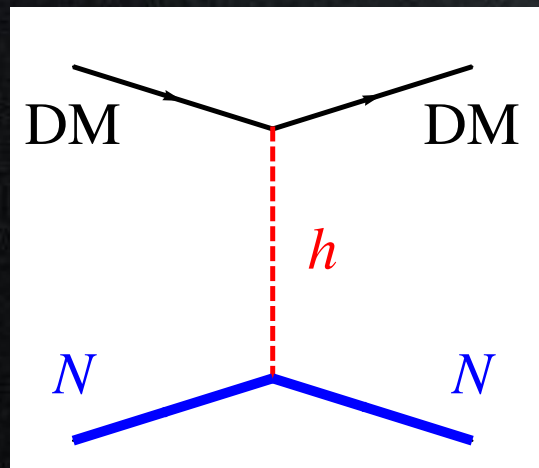


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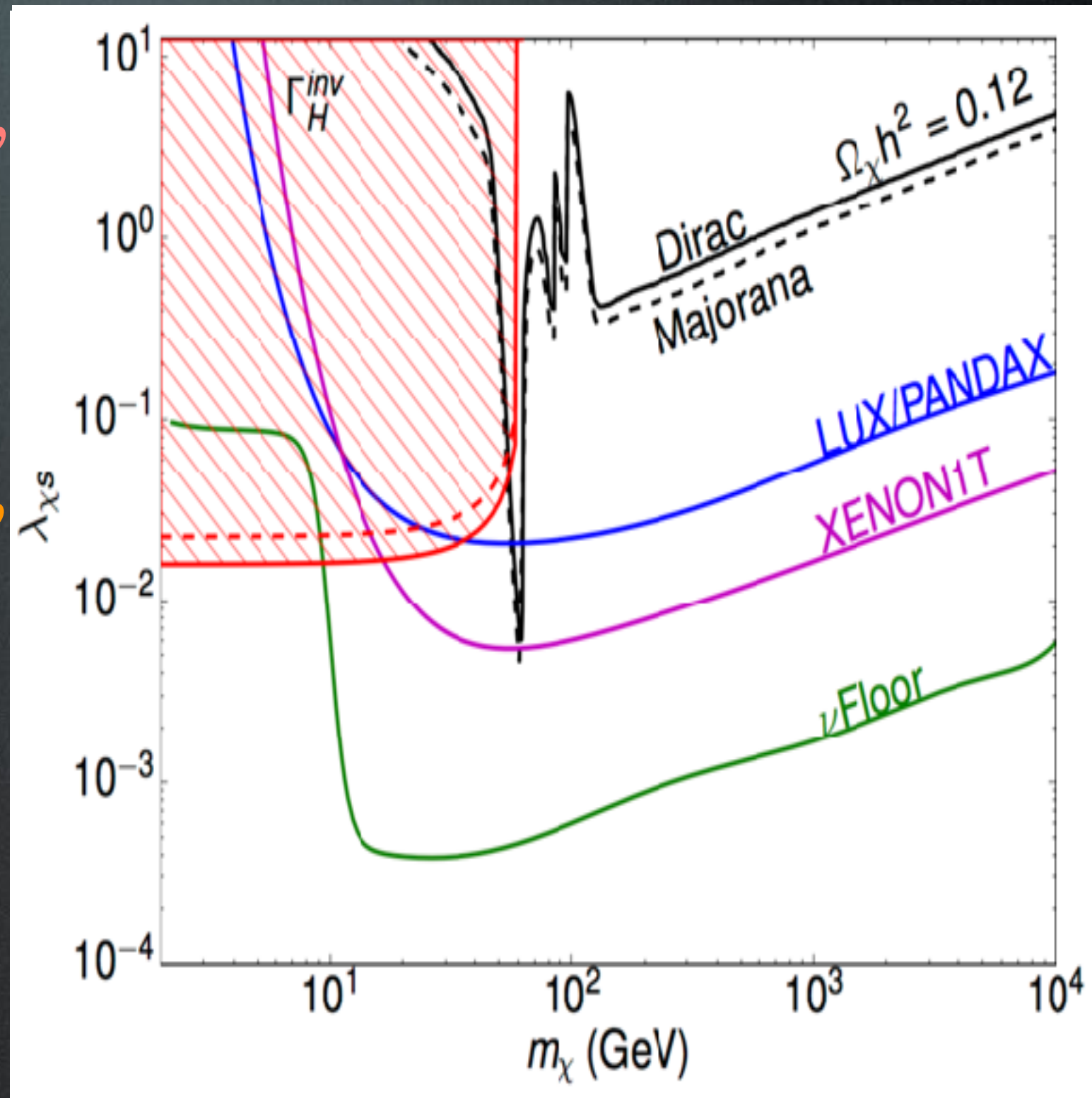
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tree level,
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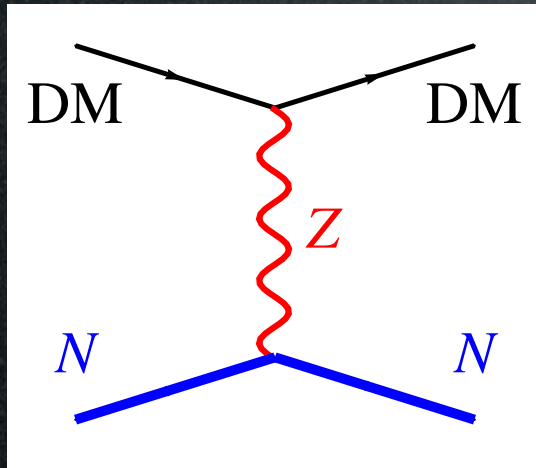


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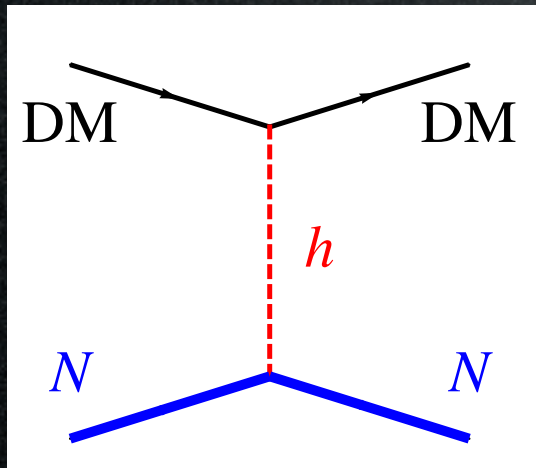
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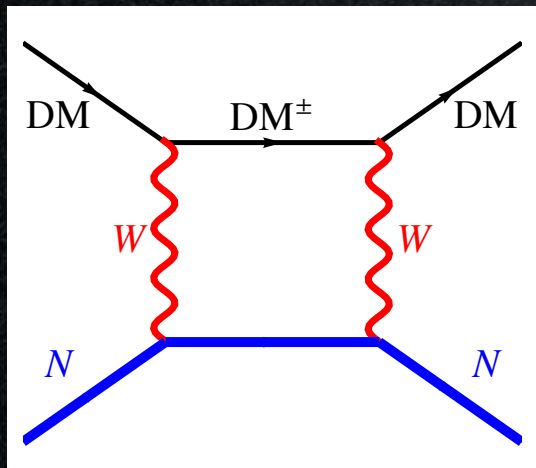
tree level,
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$$\sigma_{\text{SI}} \sim \frac{\alpha^2 m_N^2}{M_Z^4}$$



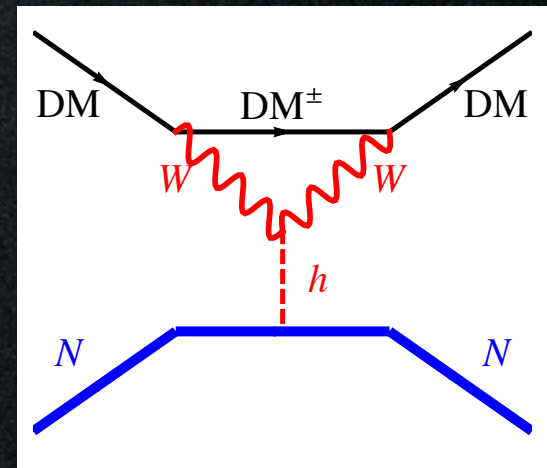
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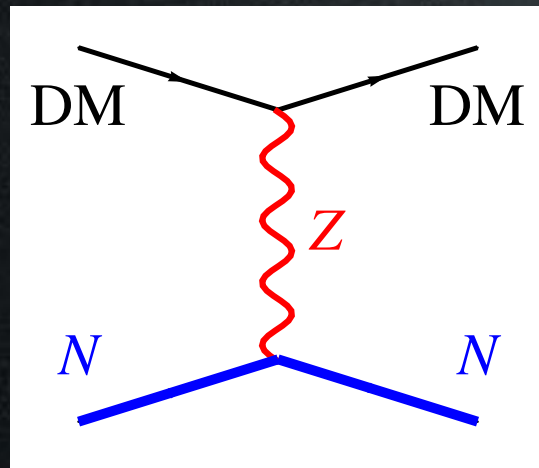
one loop

$$\sigma_{\text{SI}} \sim \frac{\alpha^4 m_N^4}{M_W^6}$$

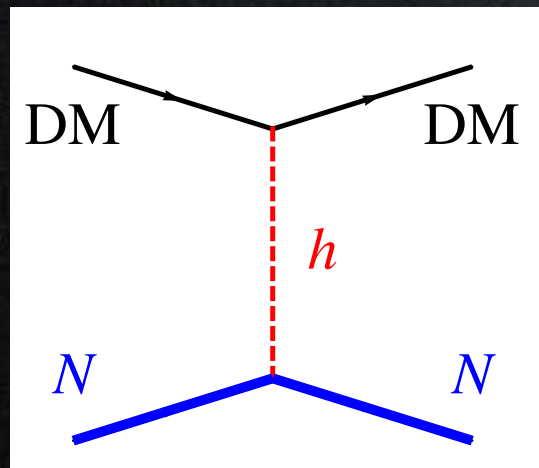


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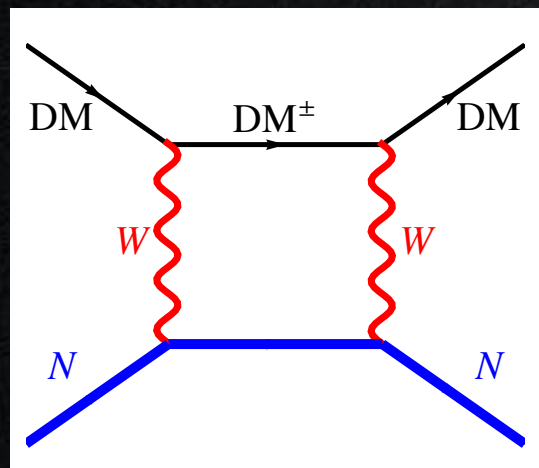
SM weak scale SI interactions



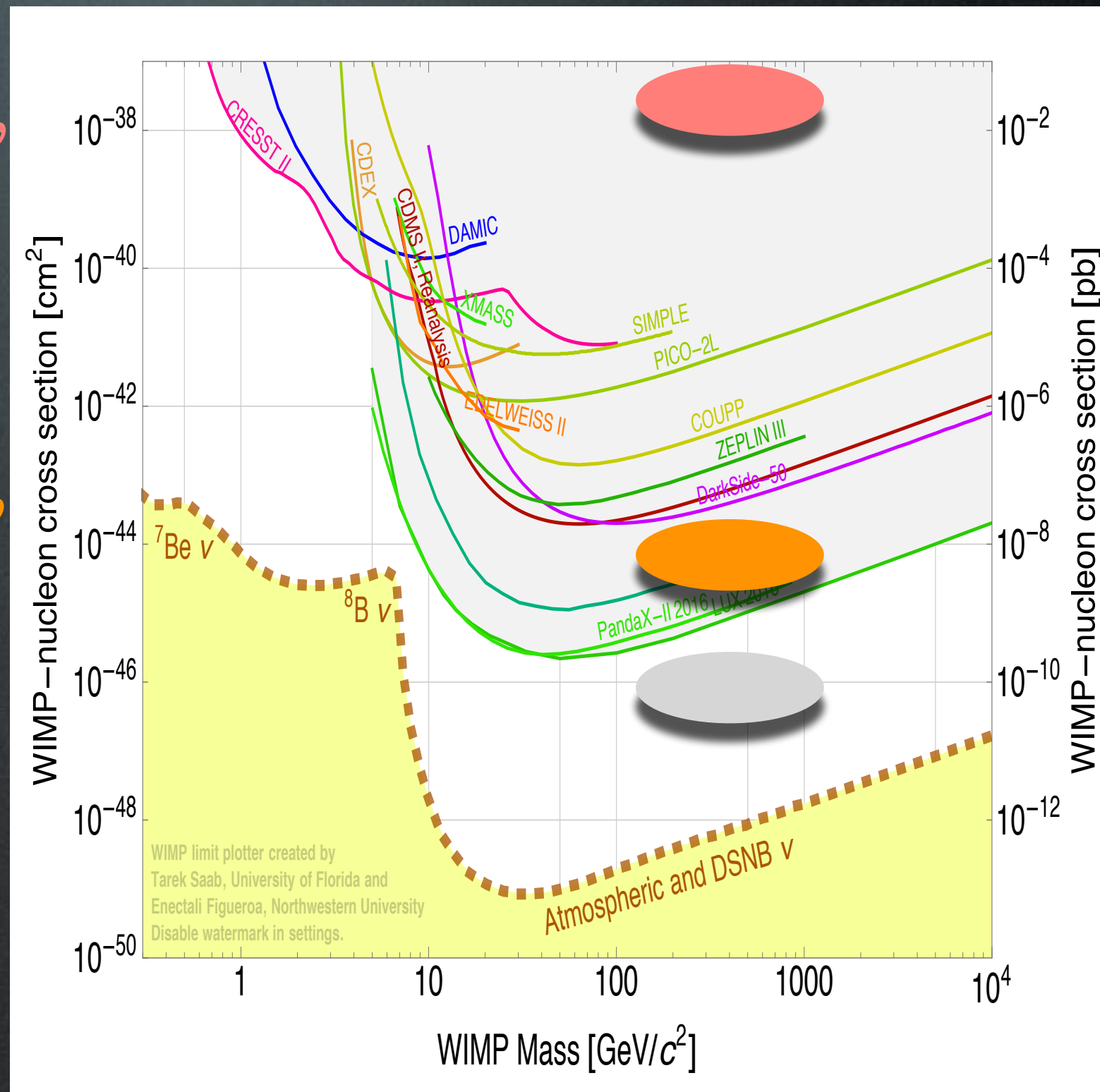
tree level,
vector



tree level,
scalar

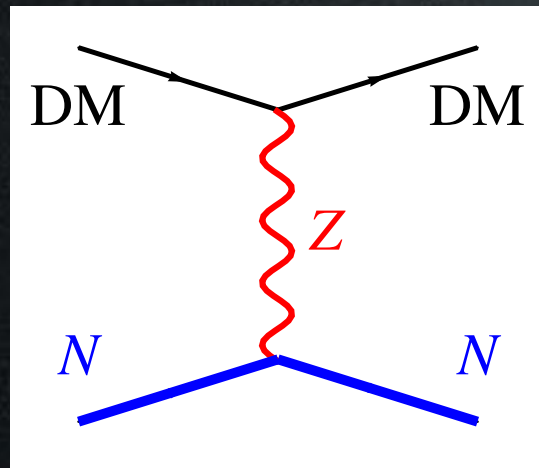


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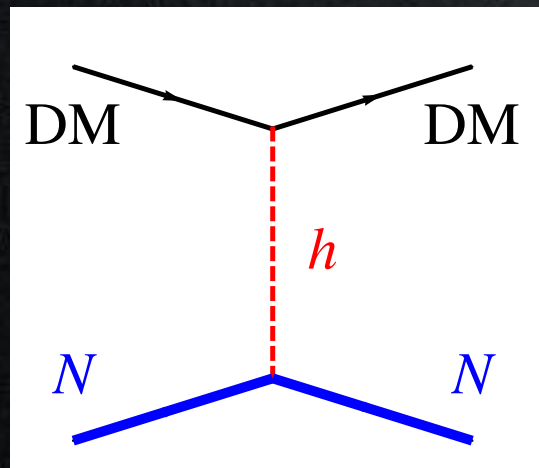
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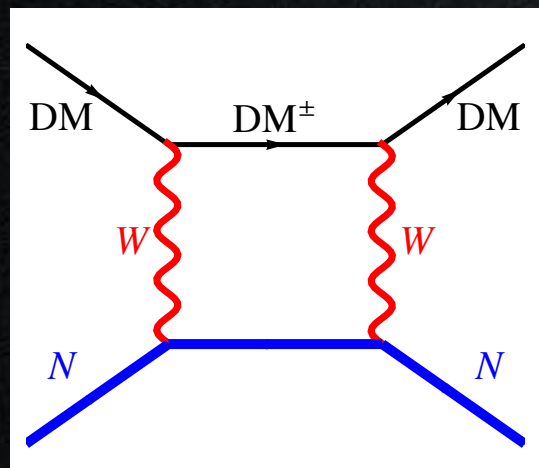


tree level,
vector

Still viable under
which conditions?



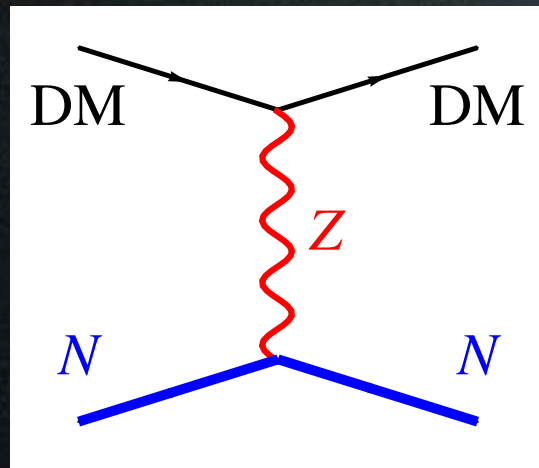
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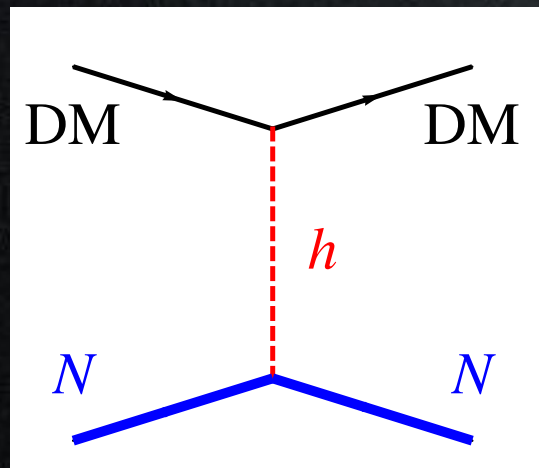
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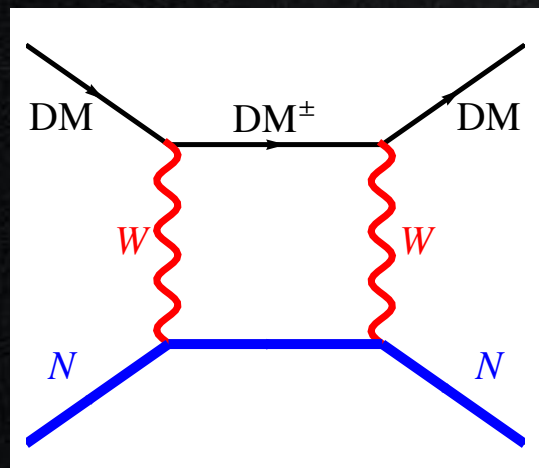
~~tree level,
vector~~

Still viable under
which conditions?

- real particle
(Majorana fermion, real scalar)



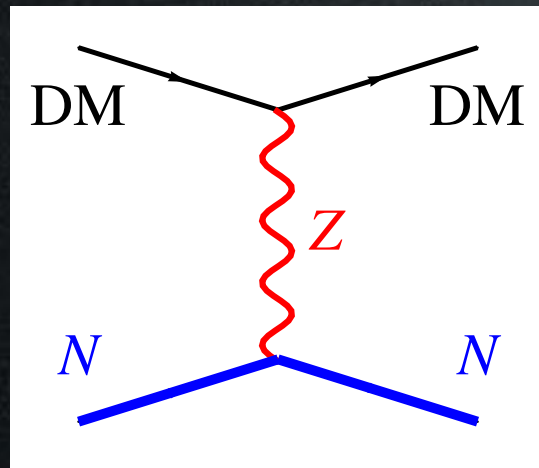
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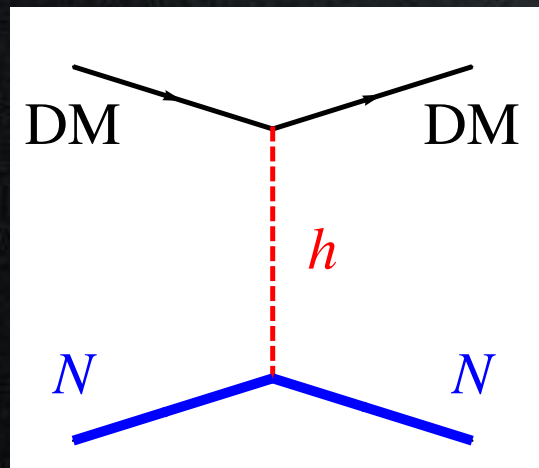
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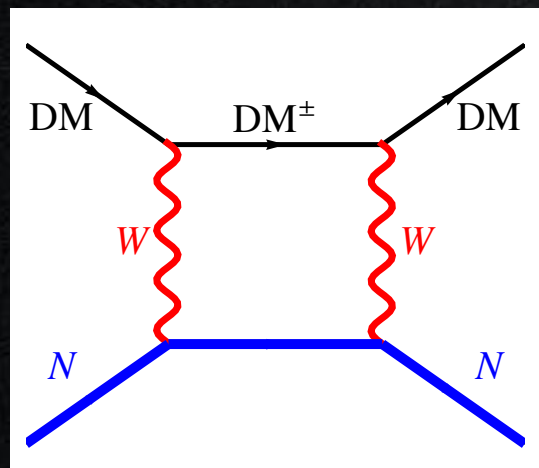
SM weak scale SI interactions



~~tree level,
vector~~



~~tree level,
scalar~~



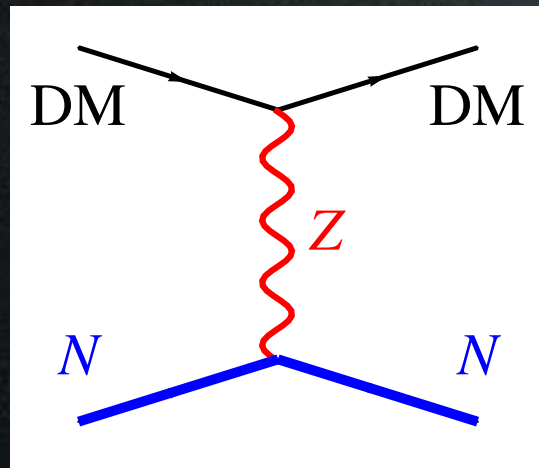
one loop

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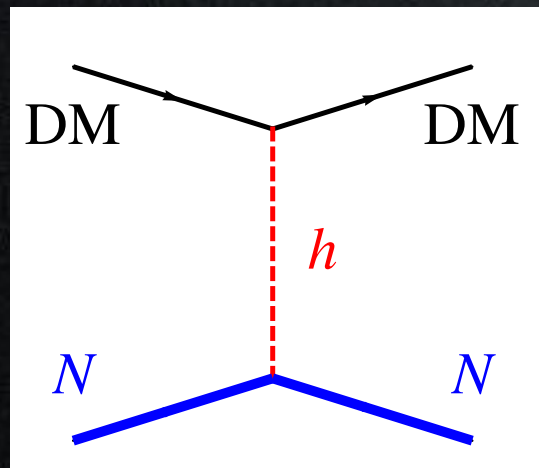
- real particle
(Majorana fermion, real scalar)
- hypercharge $Y = 0$

WIMP DD: 'theory'

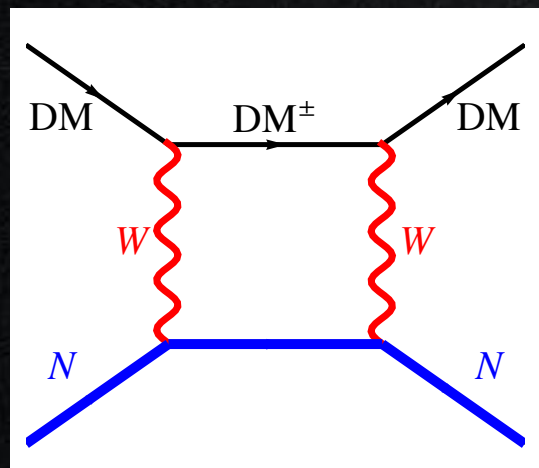
SM weak scale SI interactions



~~tree level,
vector~~



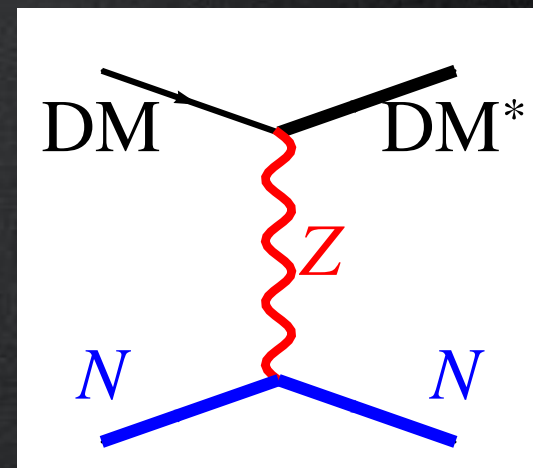
~~tree level,
scalar~~



one loop

Still viable under
which conditions?

- real particle
(Majorana fermion, real scalar)
- hypercharge $Y = 0$
- SD interactions only
- inelastic scattering



Candidates

new physics at
the TeV scale

thermal
freeze-out

WIMPs

LHC

AMS, Fermi, CTA
Antares, Icecube

Direct
Detection

1. even without a larger framework, WIMPs are **still appealing**
2. the frontier is **multi-TeV**
3. searches are **complementary** and still have **ground to cover**

Minimalistic approach

Minimalistic approach

On top of the SM, add **only** one extra multiplet $\mathcal{X} = \begin{pmatrix} \chi_1 \\ \chi_2 \\ \vdots \end{pmatrix}$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \bar{\mathcal{X}}(i\not{D} + M)\mathcal{X} \quad \text{if } \mathcal{X} \text{ is a fermion}$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + |D_\mu \mathcal{X}|^2 - M^2 |\mathcal{X}|^2 \quad \text{if } \mathcal{X} \text{ is a scalar}$$

and systematically search for the ideal DM candidate...

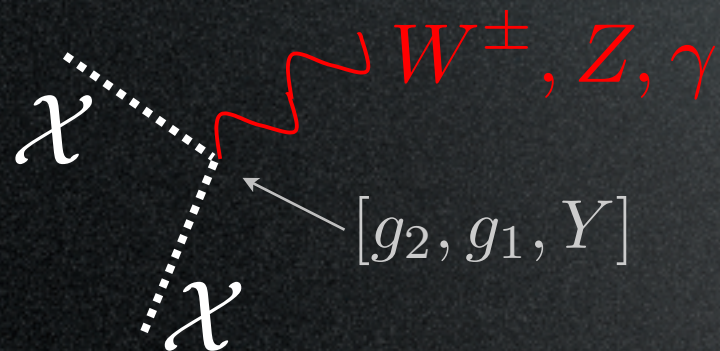
Minimalistic approach

On top of the SM, add **only** one extra multiplet $\chi = \begin{pmatrix} \chi_1 \\ \chi_2 \\ \vdots \end{pmatrix}$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \bar{\chi}(i\not{D} + M)\chi \quad \text{if } \chi \text{ is a fermion}$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + |D_\mu \chi|^2 - M^2 |\chi|^2 \quad \text{if } \chi \text{ is a scalar}$$

gauge interactions



the only parameter,
and will be fixed by Ω_{DM} .

(other terms in the
scalar potential)

(one loop mass splitting)

and systematically search for the ideal DM candidate...

The ideal DM candidate is
weakly int., massive, neutral, stable

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weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin
$\underline{2}$		
$\underline{3}$		
$\underline{4}$		
$\underline{5}$		
$\underline{7}$		

$$\chi = \begin{pmatrix} \chi_1 \\ \chi_2 \\ \vdots \\ \chi_n \end{pmatrix}$$

these are all possible choices:

$n \leq 5$ for fermions

$n \leq 7$ for scalars

to avoid explosion in the running coupling

$$\alpha_2^{-1}(E') = \alpha_2^{-1}(M) - \frac{b_2(n)}{2\pi} \ln \frac{E'}{M}$$

(actually, including 2-loops,

$n \leq 6$ ($n \leq 4$) for real (complex) scalars)

Di Luzio, Nardecchia et al., 1504.00359

← ($\underline{6}$ is similar to $\underline{4}$)

The ideal DM candidate is
weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin
$\underline{2}$	$1/2$	
$\underline{3}$	0	
	1	
$\underline{4}$	$1/2$	
	$3/2$	
$\underline{5}$	0	
	1	
	2	
$\underline{7}$	0	

Each multiplet contains a neutral component with a proper assignment of the hypercharge, according to

$$Q = T_3 + Y \equiv 0$$

e.g. for $n = 2$: $T_3 = \begin{pmatrix} +\frac{1}{2} \\ -\frac{1}{2} \end{pmatrix} \Rightarrow |Y| = \frac{1}{2}$

e.g. for $n = 3$: $T_3 = \begin{pmatrix} +1 \\ 0 \\ -1 \end{pmatrix} \Rightarrow |Y| = 0 \text{ or } 1$

etc.

The ideal DM candidate is
weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin
$\underline{2}$	$1/2$	S
		F
$\underline{3}$	0	S
		F
	1	S
		F
$\underline{4}$	$1/2$	S
		F
	$3/2$	S
		F
$\underline{5}$	0	S
		F
	1	S
		F
	2	S
		F
$\underline{7}$	0	S

Each multiplet contains a neutral component with a proper assignment of the hypercharge, according to

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etc.

The ideal DM candidate is weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin	M (TeV)
$\underline{2}$	$1/2$	S	0.43
		F	1.2
$\underline{3}$	0	S	2.0
		F	2.6
	1	S	1.4
		F	1.8
$\underline{4}$	$1/2$	S	2.4
		F	2.5
	$3/2$	S	2.4
		F	2.5
$\underline{5}$	0	S	5.0
		F	4.2
	1	S	3.5
		F	3.2
	2	S	3.5
		F	3.2
$\underline{7}$	0	S	8.5

The **mass** M is determined by the relic abundance:

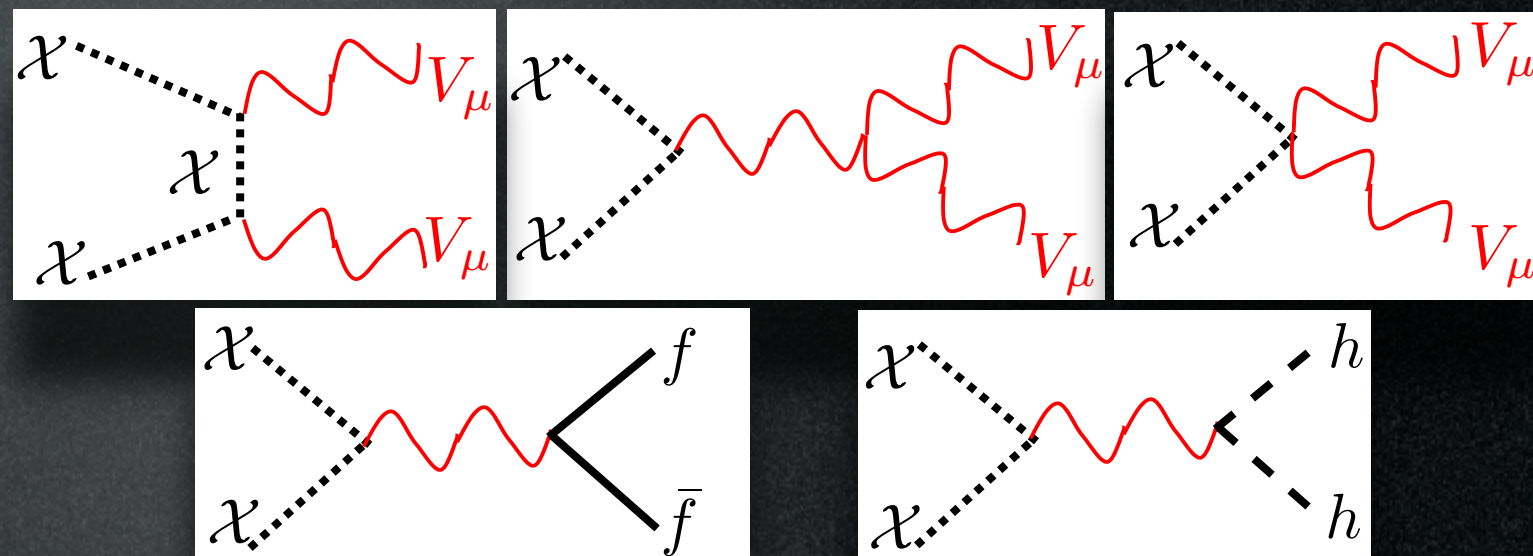
$$\Omega_{\text{DM}} = \frac{6 \cdot 10^{-27} \text{cm}^3 \text{s}^{-1}}{\langle \sigma_{\text{ann}} v \rangle} \cong 0.24$$

for χ scalar

$$\langle \sigma_A v \rangle \simeq \frac{g_2^4 (3 - 4n^2 + n^4) + 16 Y^4 g_Y^4 + 8g_2^2 g_Y^2 Y^2 (n^2 - 1)}{64\pi M^2 g_\chi}$$

for χ fermion

$$\langle \sigma_A v \rangle \simeq \frac{g_2^4 (2n^4 + 17n^2 - 19) + 4Y^2 g_Y^4 (41 + 8Y^2) + 16g_2^2 g_Y^2 Y^2 (n^2 - 1)}{128\pi M^2 g_\chi}$$



(- include co-annihilations)
(- computed for $M \gg M_{Z,W}$)

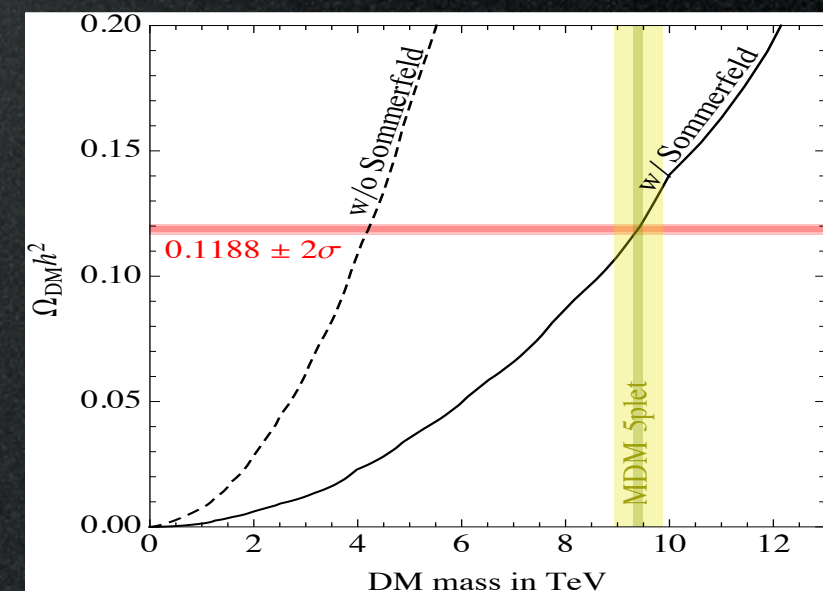
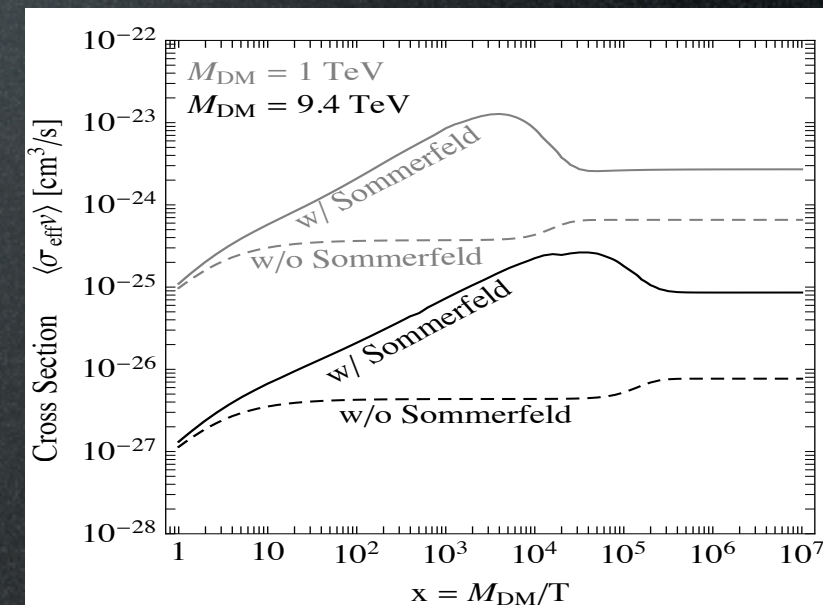
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$SU(2)_L$	$U(1)_Y$	spin	M (TeV)
$\underline{2}$	$1/2$	S	1.0
		F	
$\underline{3}$	0	S	2.5
		F	2.7
	1	S	
		F	
$\underline{4}$	$1/2$	S	
		F	
	$3/2$	S	
		F	
$\underline{5}$	0	S	9.4
		F	
	1	S	
		F	
	2	S	
		F	
$\underline{7}$	0	S	25

Non-perturbative corrections
 (and other smaller corrections) (more later)
 induce modifications:

$$\langle \sigma_{\text{ann}} v \rangle \rightsquigarrow R \cdot \langle \sigma_{\text{ann}} v \rangle + \langle \sigma_{\text{ann}} v \rangle_{p\text{-wave}}$$

$$\text{with } R \sim \mathcal{O}(\text{few}) \rightarrow \mathcal{O}(10^2)$$

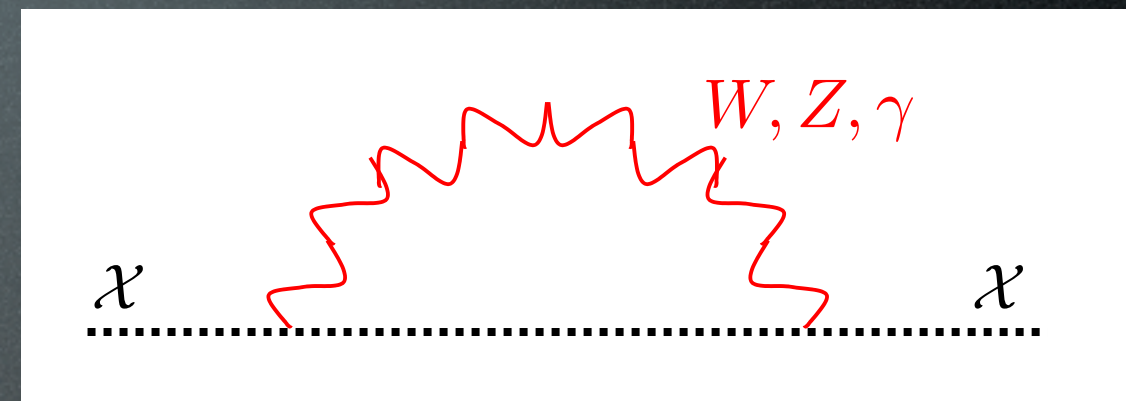


The ideal DM candidate is
weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin	M (TeV)	ΔM (MeV)
$\underline{2}$	$1/2$	S	1.0	348
		F		342
$\underline{3}$	0	S	2.5	166
		F	2.7	166
	1	S		540
		F		526
$\underline{4}$	$1/2$	S		353
		F		347
	$3/2$	S		729
		F		712
$\underline{5}$	0	S	9.4	166
		F		166
	1	S		537
		F		534
	2	S		906
		F		900
$\underline{7}$	0	S	25	166

EW loops induce
a **mass splitting** ΔM
inside the n-uplet:

tree
level



$$M_Q - M_{Q'} = \frac{\alpha_2 M}{4\pi} \left\{ (Q^2 - Q'^2) s_W^2 f\left(\frac{M_Z}{M}\right) + (Q - Q')(Q + Q' - 2Y) \left[f\left(\frac{M_W}{M}\right) - f\left(\frac{M_Z}{M}\right) \right] \right\}$$

with $f(r) \xrightarrow{r \rightarrow 0} -2\pi r$

The neutral component
is the lightest



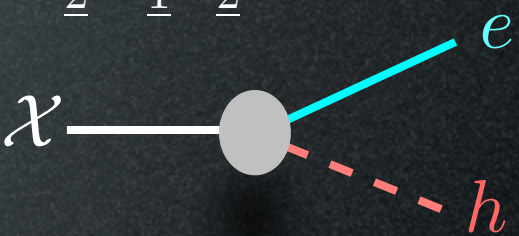
The ideal DM candidate is

weakly int., massive, neutral, stable

$SU(2)_L$	$U(1)_Y$	spin	M (TeV)	ΔM (MeV)	decay ch.
<u>2</u>	1/2	S		348	EL
		F	1.0	342	EH
<u>3</u>	0	S	2.5	166	HH^*
		F	2.7	166	LH
	1	S		540	HH, LH
		F		526	LH
<u>4</u>	1/2	S		353	HHH^*
		F		347	(LHH^*)
	3/2	S		729	HHH
		F		712	(LHH)
<u>5</u>	0	S		166	(HHH^*H^*)
		F	9.4	166	—
	1	S		537	$(HH^*H^*H^*)$
		F		534	—
	2	S		906	$(H^*H^*H^*H^*)$
		F		900	—
<u>7</u>	0	S	25	166	$(\chi\chi H^*H)$

List all **allowed SM couplings**:

e.g. $\chi \begin{matrix} 1/2 & -1 & 1/2 \\ \underline{2} & \underline{1} & \underline{2} \end{matrix} EH$



e.g. $\chi \begin{matrix} 1/2 & -1/2 & 1/2 & -1/2 \\ \underline{4} & \underline{2} & \underline{2} & \underline{2} \end{matrix} LHH^*$

dim=5 operator, induces $\tau \sim \Lambda^2 \text{TeV}^{-3} \ll t_{\text{universe}}$ for $\Lambda \sim M_{\text{Pl}}$

dim=5, loop decay

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dim=5 operator, induces
 $\tau \sim \Lambda^2 \text{TeV}^{-3} \ll t_{\text{universe}}$
 for $\Lambda \sim M_{\text{Pl}}$

No allowed decay!
Automatically stable!

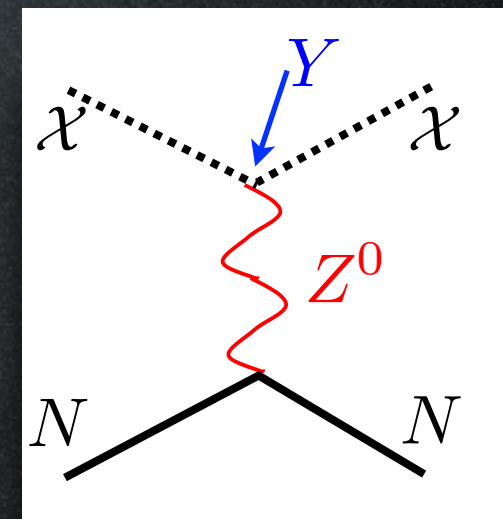
The ideal DM candidate is
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 and
not excluded
 by direct searches!

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Candidates with $Y \neq 0$
 interact as



$$\sigma \simeq G_F^2 M_{\mathcal{N}}^2 Y^2$$

Goodman
Witten
1985

$$\gg \text{present bounds}$$

e.g. LUX, Xenon, PandaX

↓

need $Y = 0$

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		F		900	—
	0	S			
		F			
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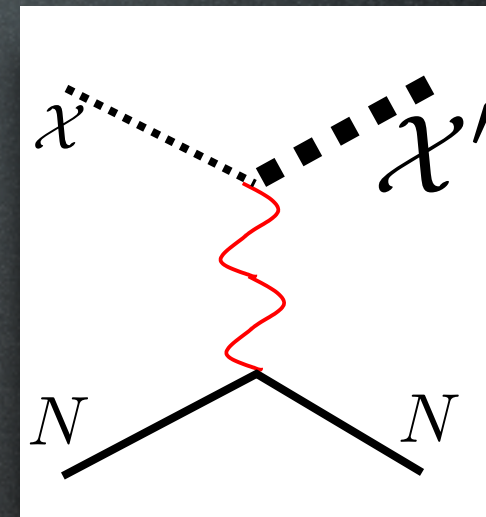
← We have a winner!

(other terms in the
scalar potential)

If you want to cure ill candidates...

$Y \neq 0$: introduce some mechanism to forbid coupling with Z^0 anyway

e.g. mixing with an extra state splits the 2 components of $\mathcal{X}$; if splitting is large enough, NC scattering is kinematically forbidden...



stability: impose some symmetry to forbid decays (e.g. R-parity)...



...the case of SuSy higgsino

mixing is with bino;
even for pure higgsino,
some mixing is 'inevitable'
due to higher dim operators

The ideal DM candidate is
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		F		900	—
<u>7</u>	0	S	25	166	—

← And a second place

← We have a winner!

(other terms in the scalar potential)

Recap:

A fermionic $SU(2)_L$ quintuplet with $Y = 0$
provides a DM candidate with $M = 9.4$ TeV,
which is fully successful:

- neutral

- ***automatically*** stable 

like proton
stability in SM!

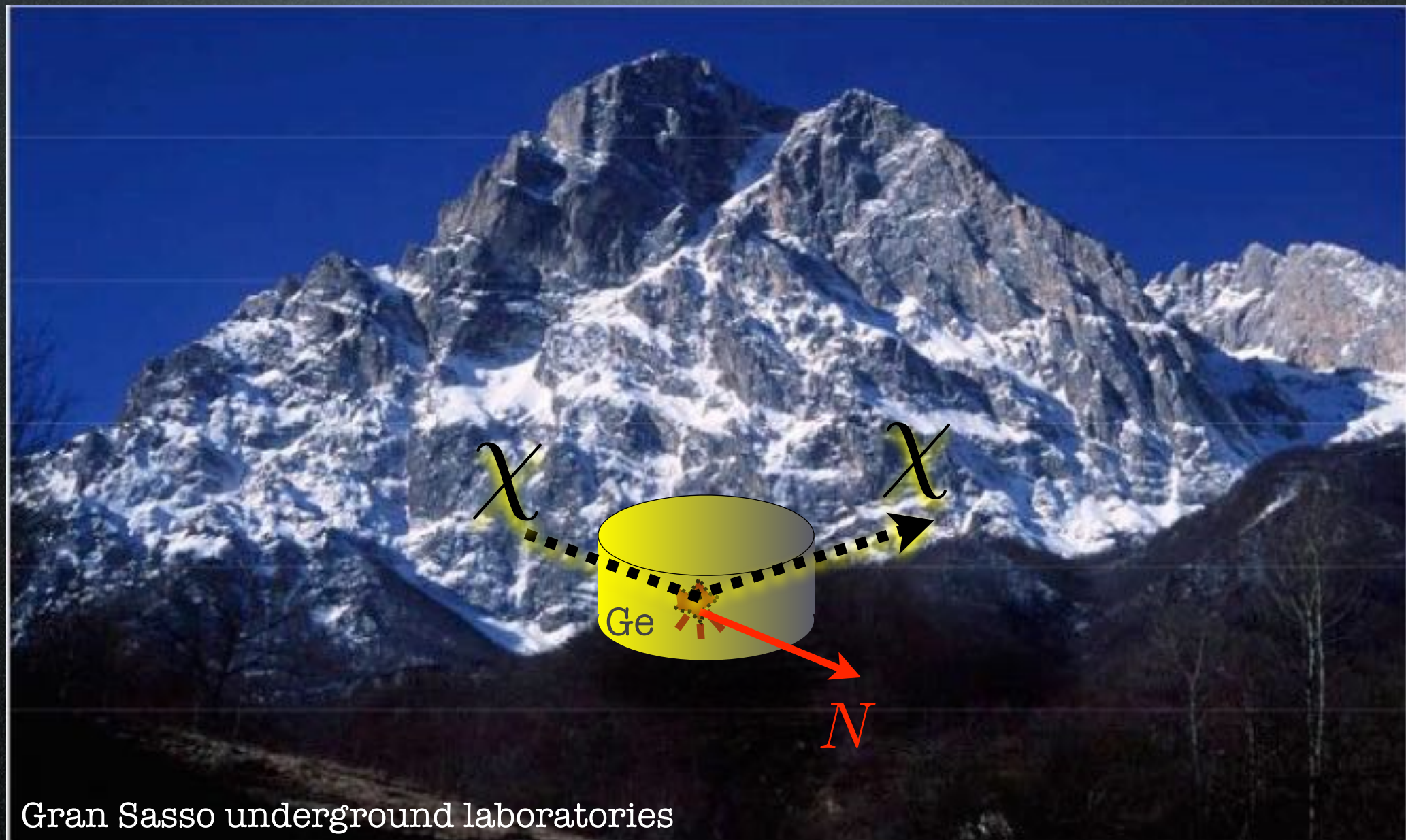
and

not _{yet} discovered by DM searches.

(Other candidates can be cured via non-minimalities.)

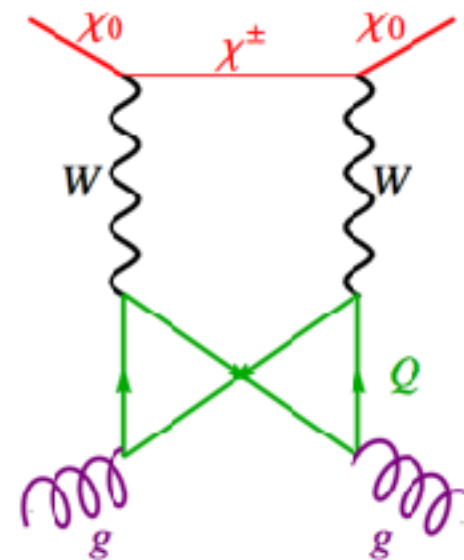
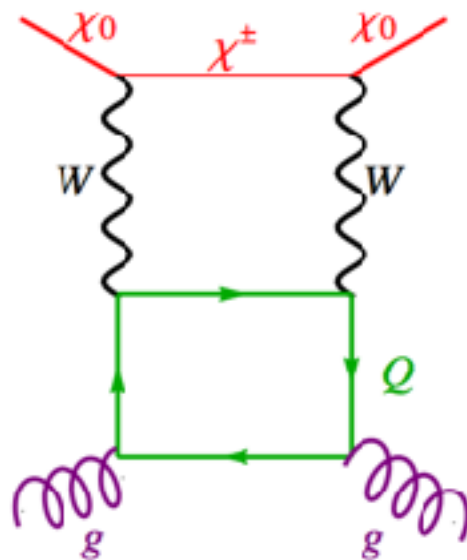
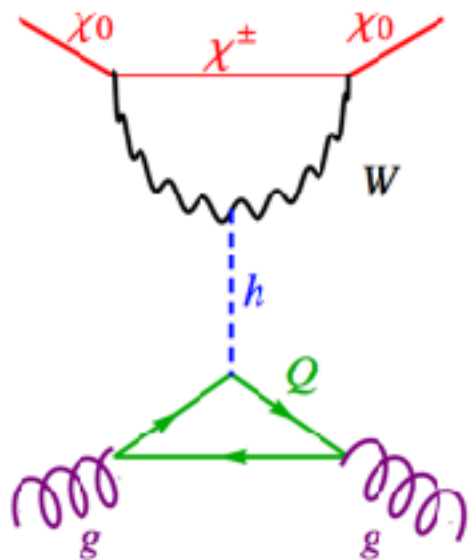
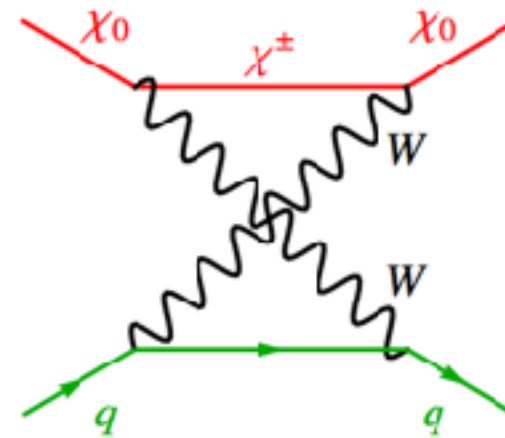
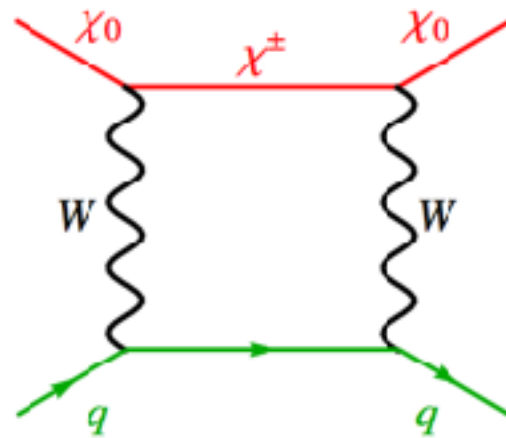
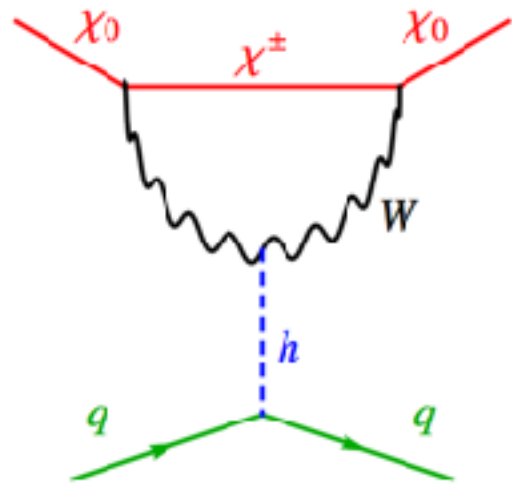
Detection and Phenomenology

Direct Detection



Direct Detection

No tree level scattering.
1-loop and 2-loops:



Cirelli, Fornengo, Strumia
hep-ph/0512090

Essig 0710.1668

Hisano, Ishiwata, Nagata
1007.2601

Hisano, Ishiwata, Nagata,
Takesano 1104.0228

Hill, Solon 1111.0016

Hisano, Ishiwata, Nagata
1010.5985

Farina, Pappadopulo, Strumia
1303.7244

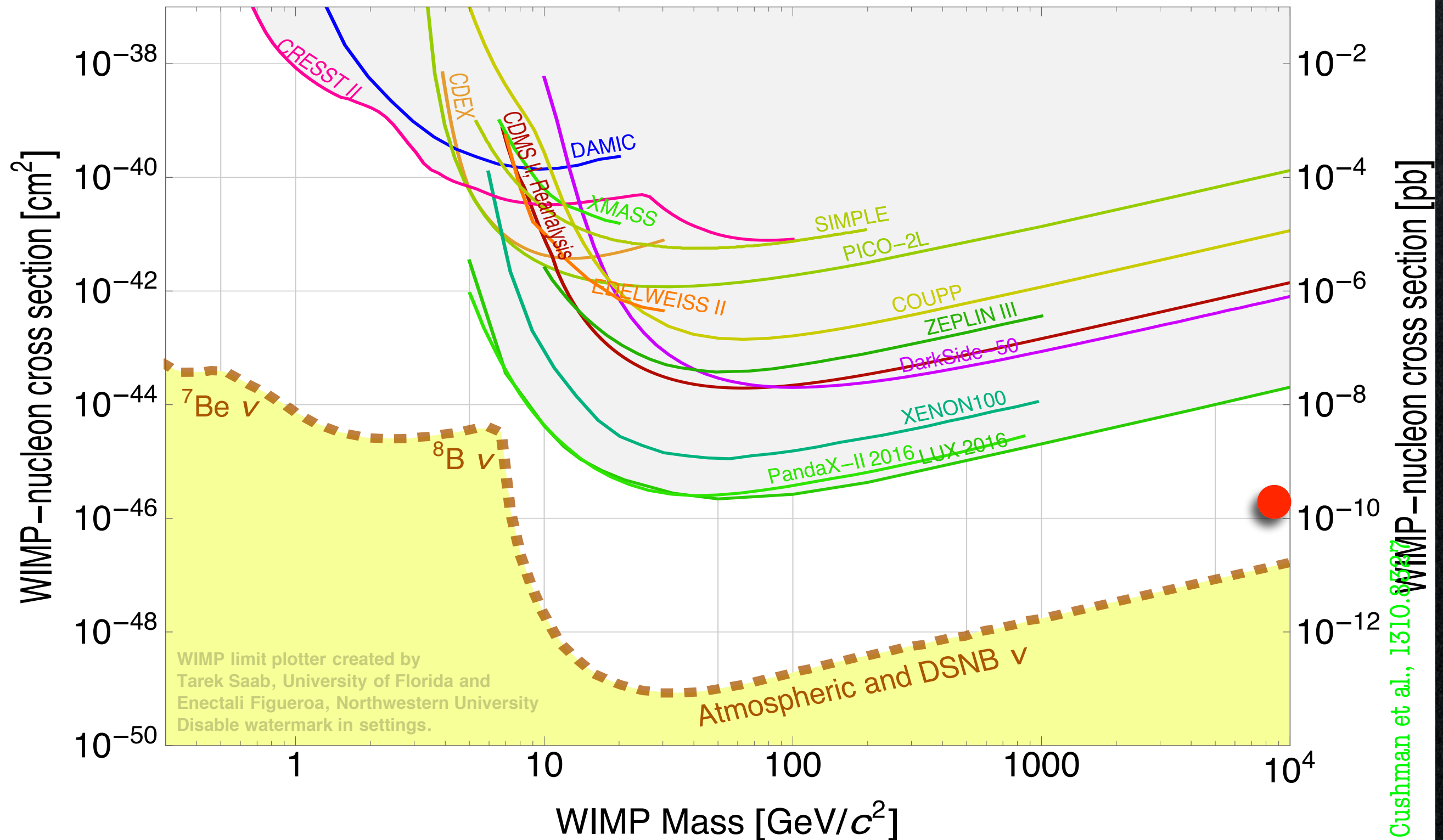
Hill, Solon 1309.4092

Hill, Solon 1401.3339

Hisano, Ishiwata, Nagata
1504.00915

$$\sigma_{\text{SI}}^n = 2 \cdot 10^{-46} \text{ cm}^2$$

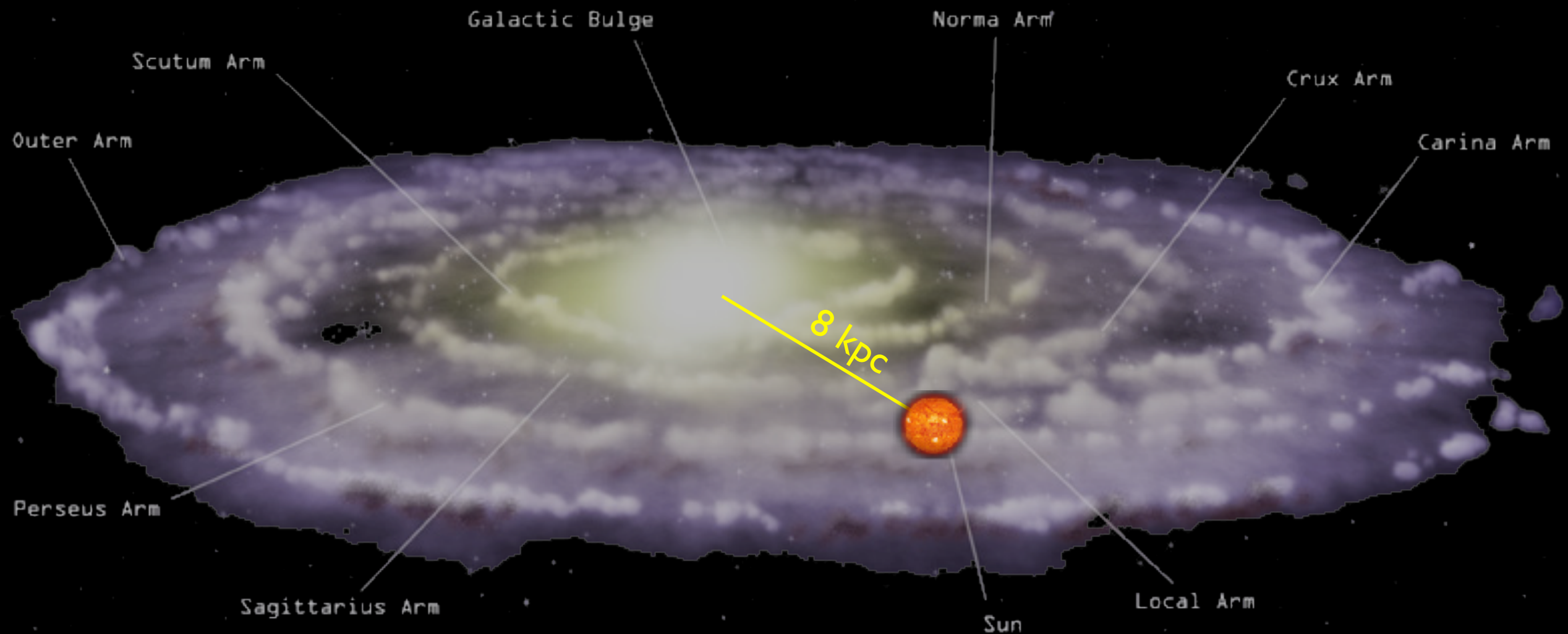
Direct Detection



PS: SD cross section equally challenging [Hisano, Ishiwata, Nagata 1010.5985](#)

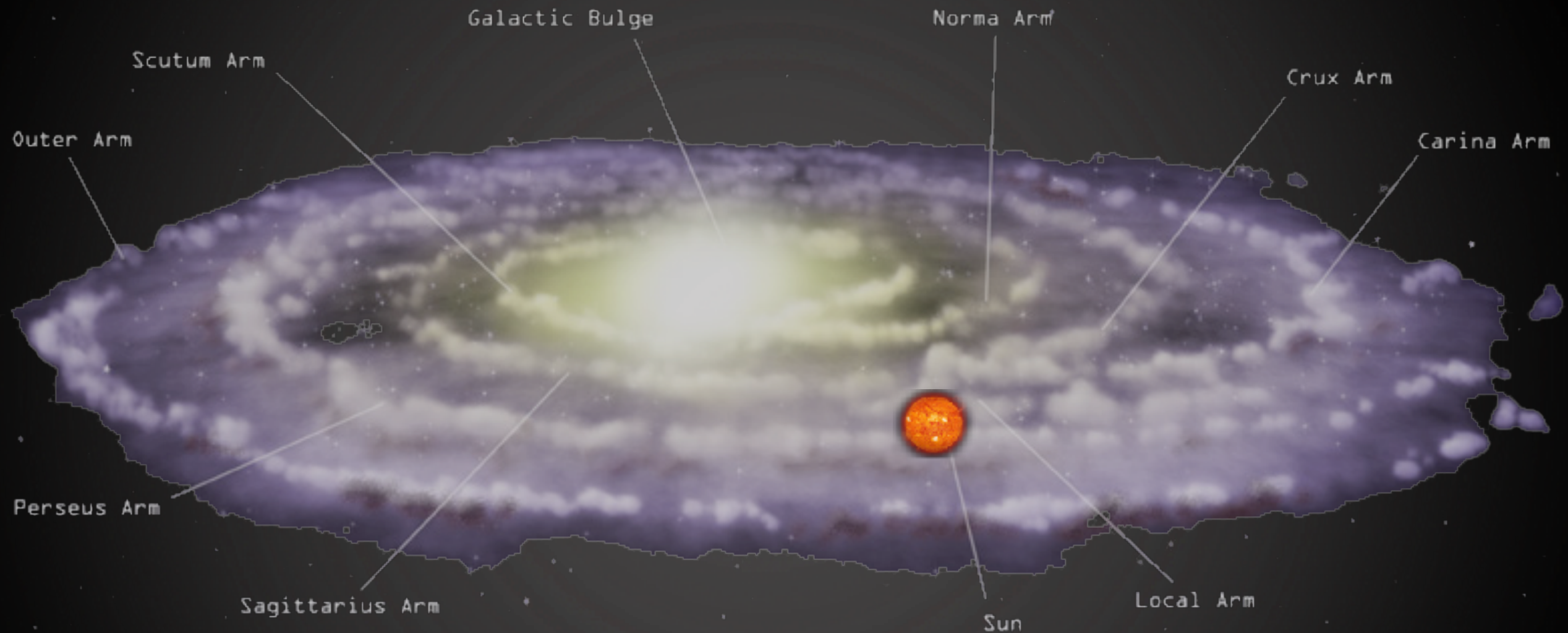
Indirect Detection

i.e. $\nu, \bar{p}, e^+, \gamma, \bar{D}$ from MDM annihilations in MW halo.



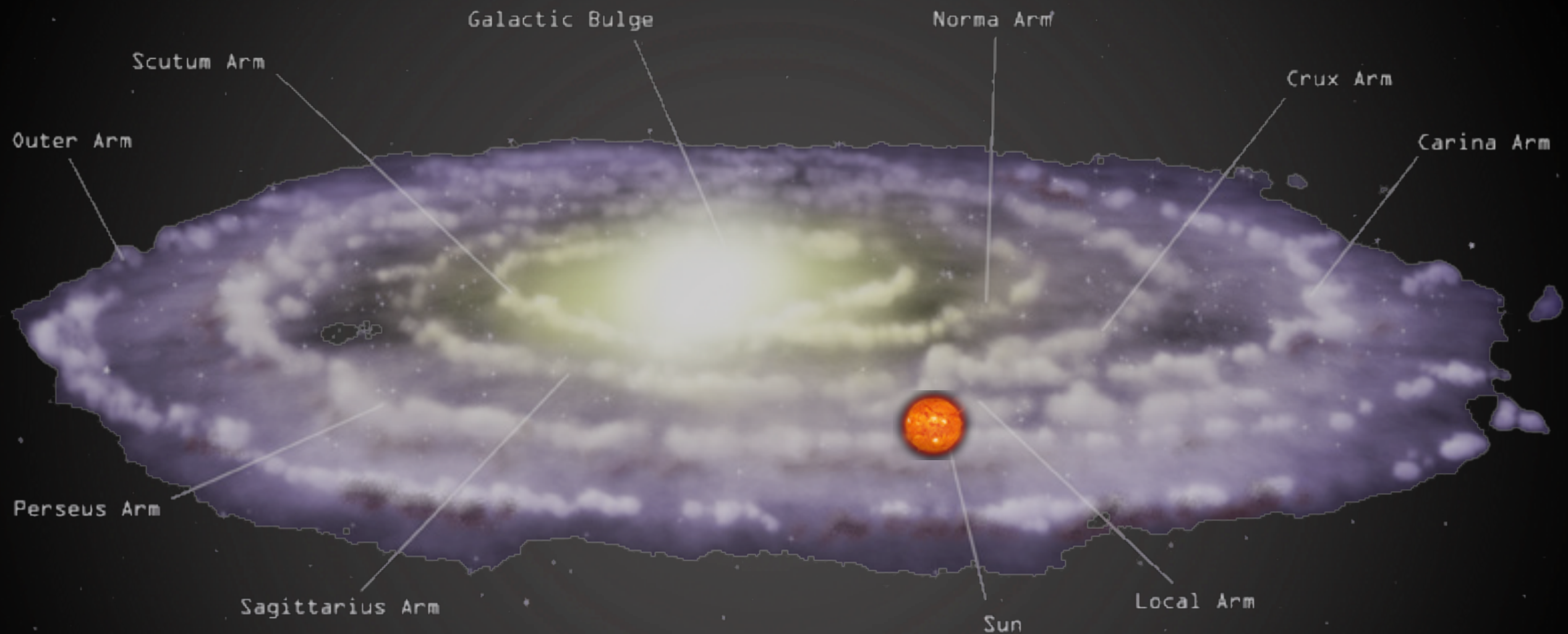
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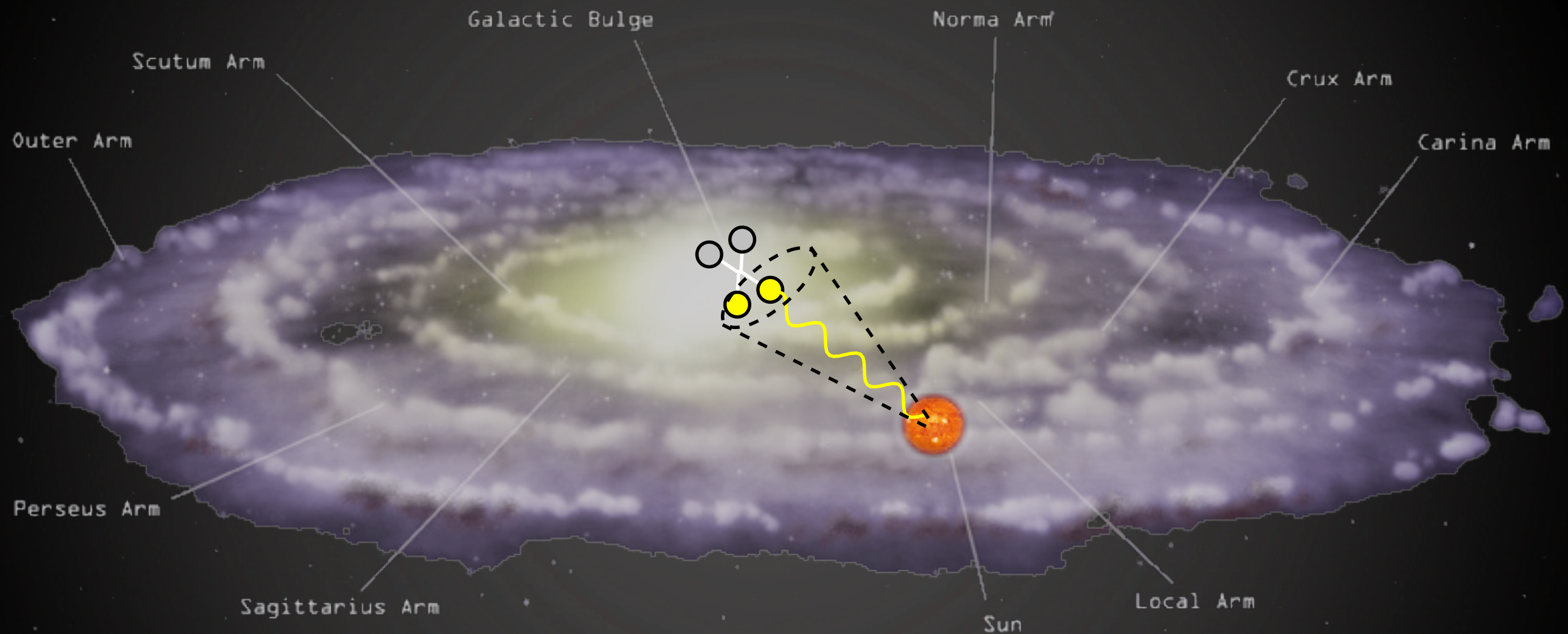
Indirect Detection

γ from MDM annihilations



Indirect Detection

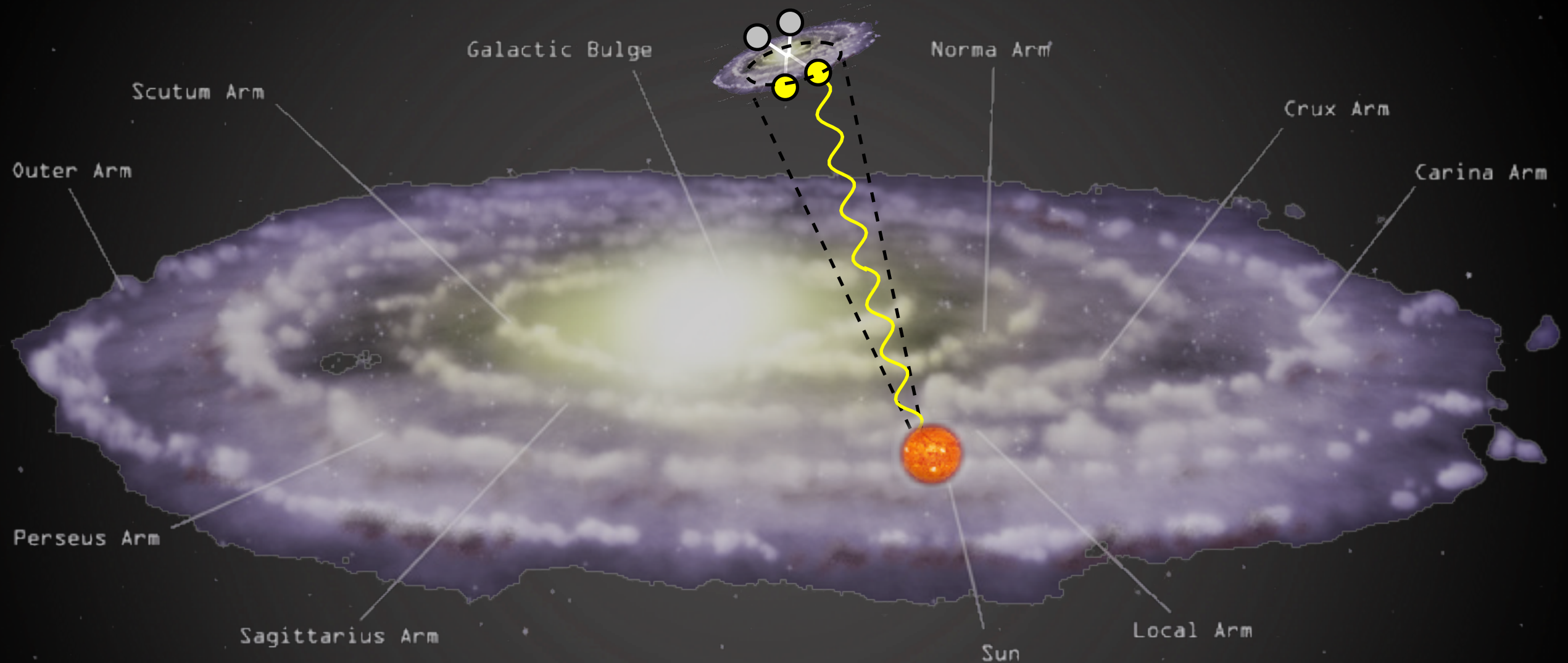
γ from MDM annihilations in galactic center



$$\begin{aligned}
 &DM \rightarrow W^-, Z, b, \tau^-, t, h \dots \rightsquigarrow e^\mp, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma \\
 &DM \rightarrow W^+, Z, \bar{b}, \tau^+, \bar{t}, h \dots \rightsquigarrow e^\pm, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma
 \end{aligned}$$

Indirect Detection

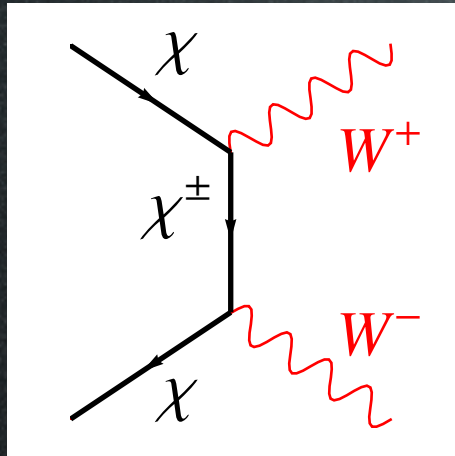
γ from MDM annihilations in dwarf galaxies



$$\begin{aligned}
 &DM \text{---} \bullet \text{---} W^-, Z, b, \tau^-, t, h \dots \rightsquigarrow e^\mp, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma \\
 &DM \text{---} \bullet \text{---} W^+, Z, \bar{b}, \tau^+, \bar{t}, h \dots \rightsquigarrow e^\pm, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma
 \end{aligned}$$

Indirect Detection

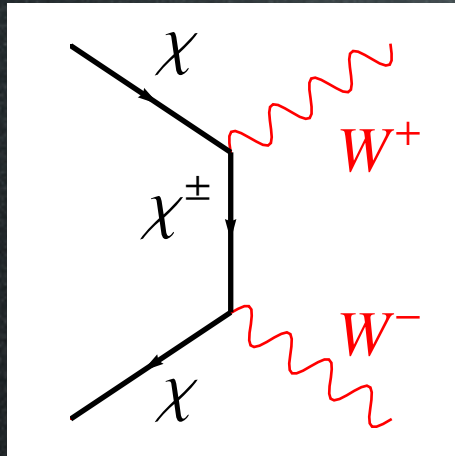
γ from MDM annihilations



$$+ \quad W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$$

Indirect Detection

γ from MDM annihilations

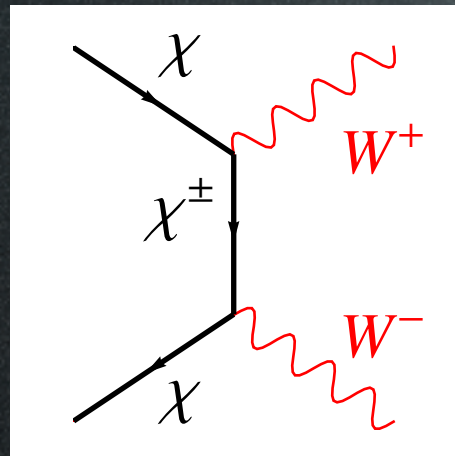


↓
continuum

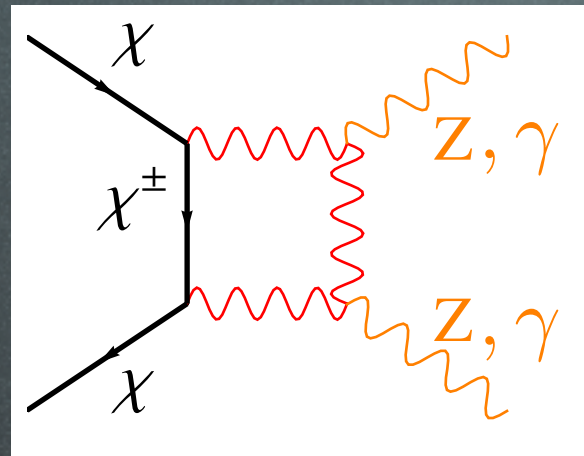
$$+ \quad W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$$

Indirect Detection

γ from MDM annihilations



continuum



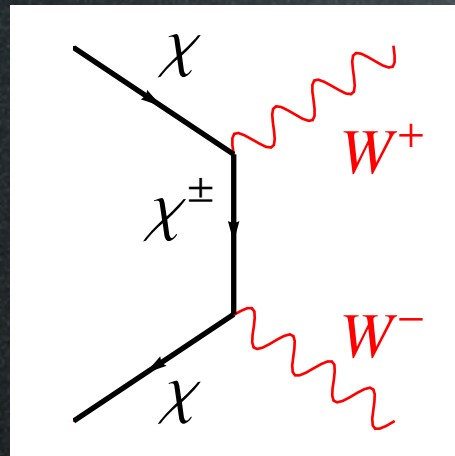
line(s)
(+ continuum)

+ $W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$

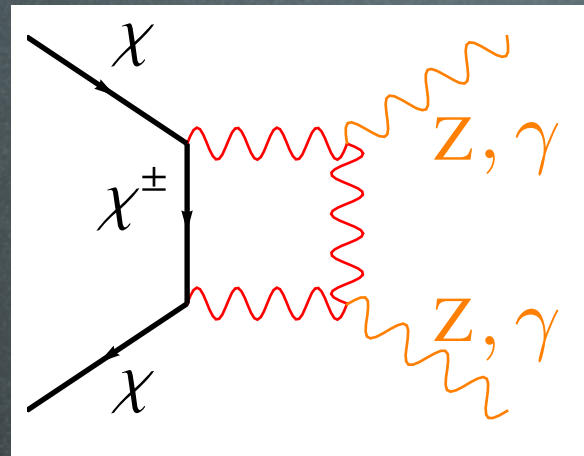
(channels for MDM with $Y=0$)

Indirect Detection

γ from MDM annihilations



continuum



line(s)
(+ continuum)

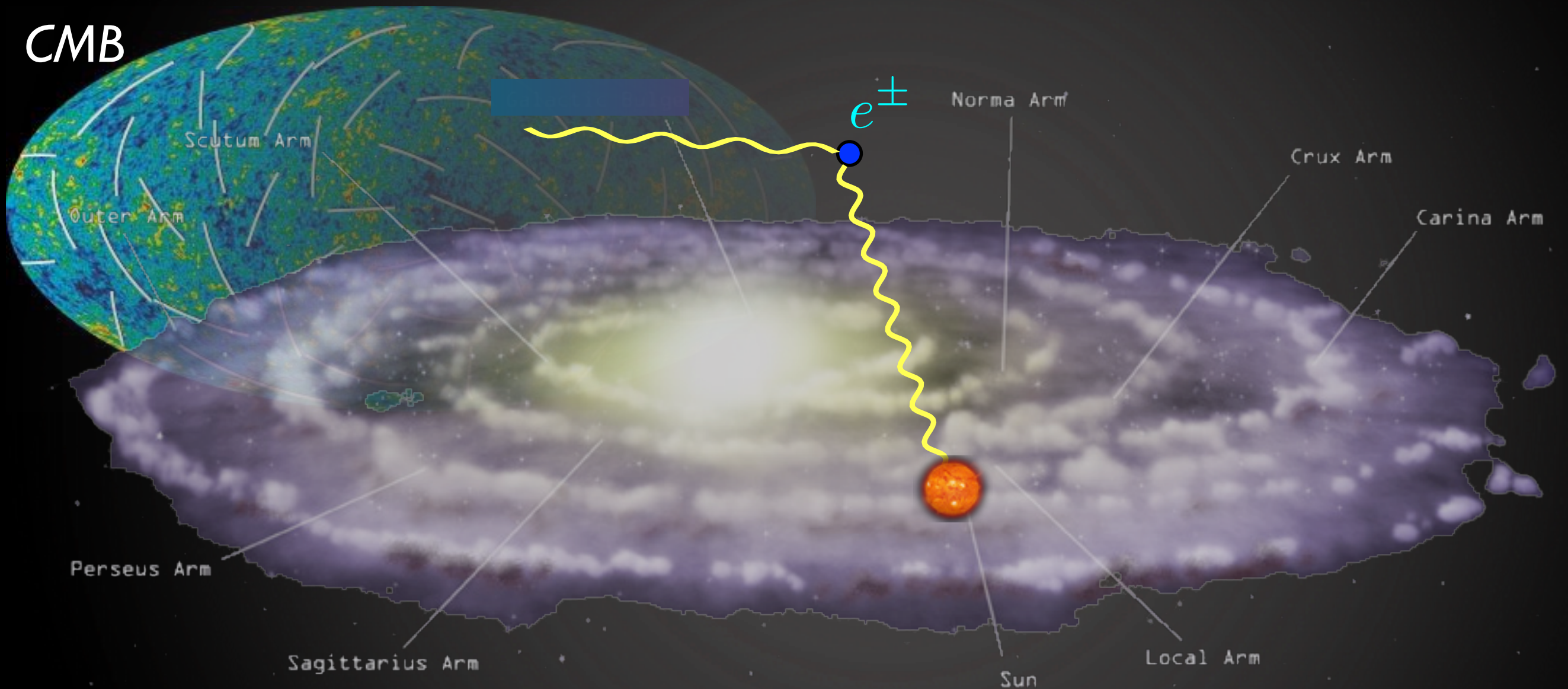
$$+ W^\pm, Z \rightarrow \bar{p} e^+ \gamma \dots$$

(channels for MDM with $Y=0$)

+ ICS

Secondary emission

γ from Inverse Compton on $e^\pm$ in halo

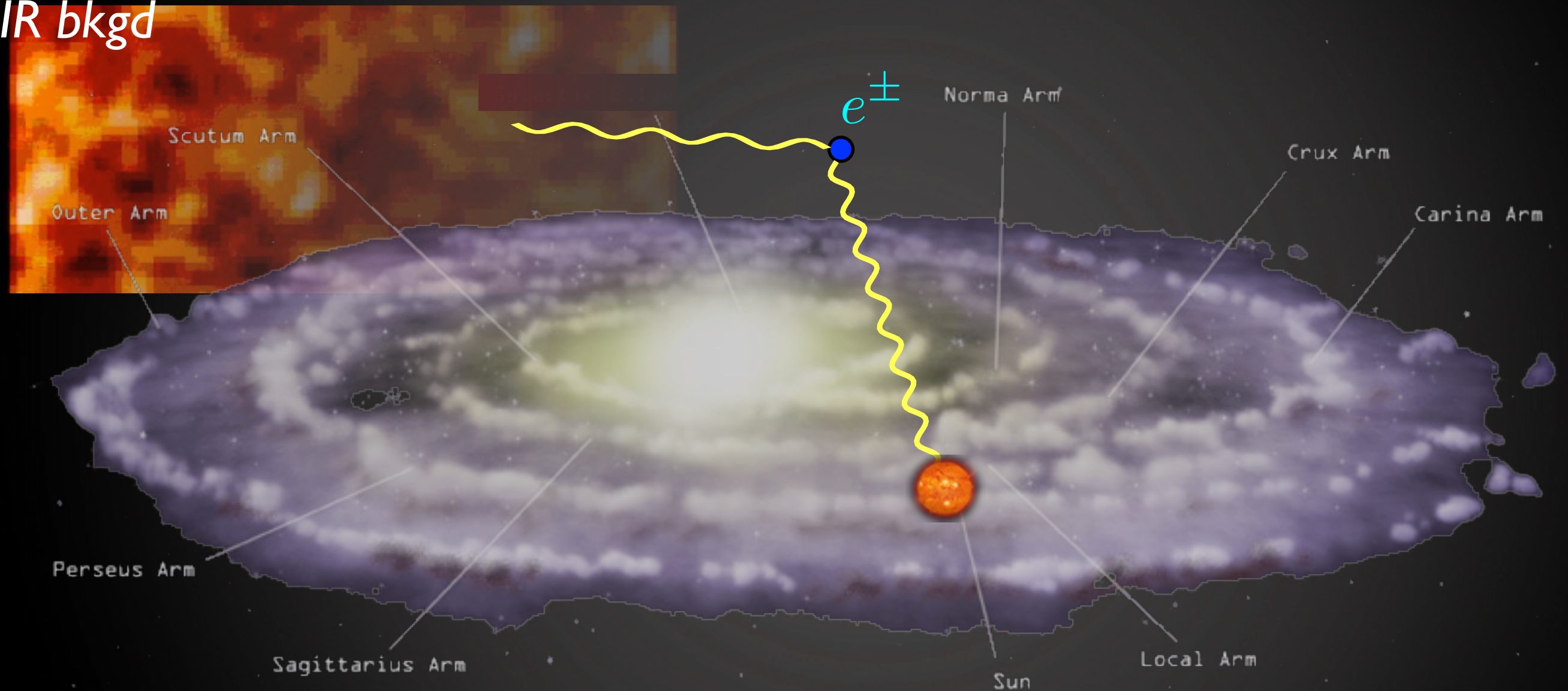


- upscatter of CMB, infrared and starlight photons on energetic $e^\pm$
- probes regions outside of Galactic Center

Secondary emission

γ from Inverse Compton on $e^\pm$ in halo

IR bkgd

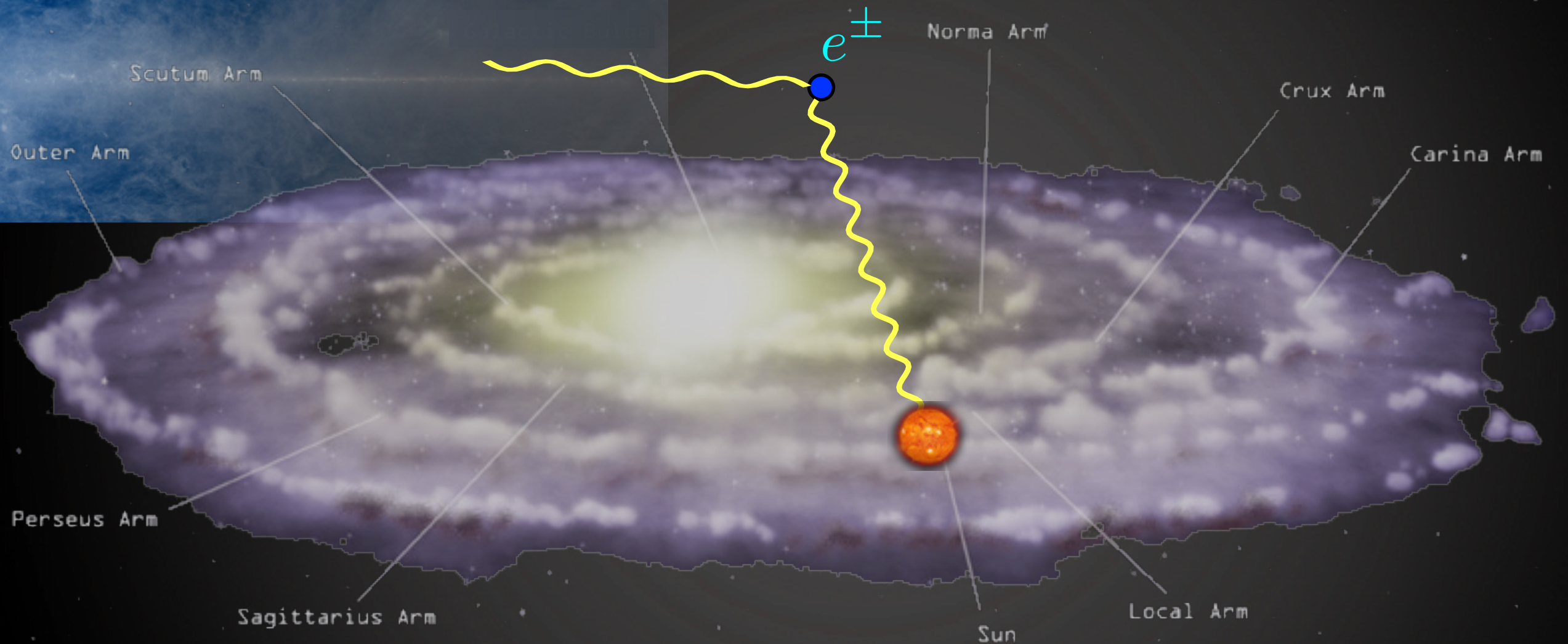


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Star Light

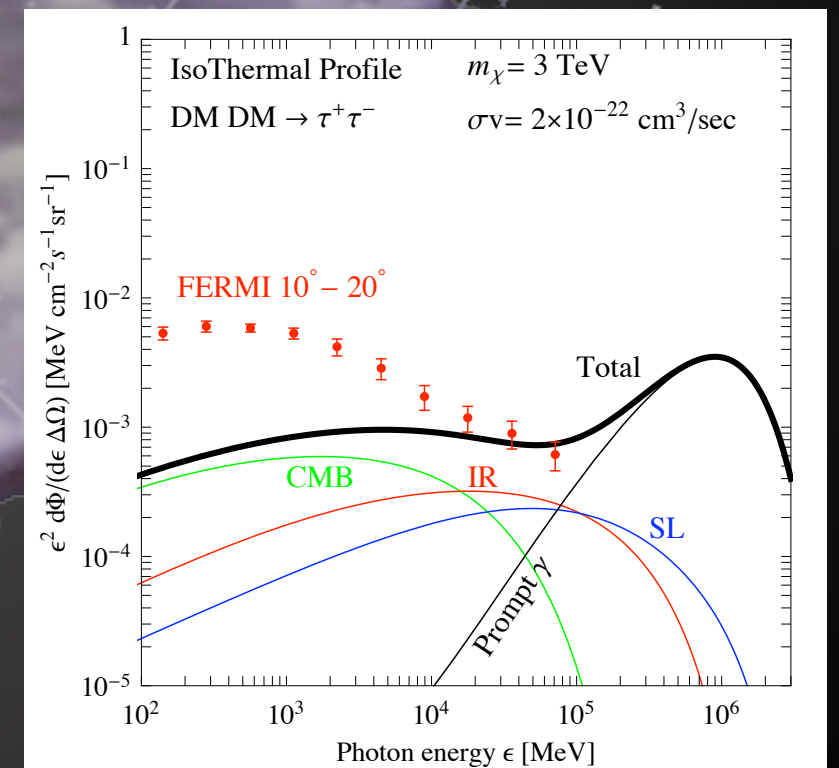
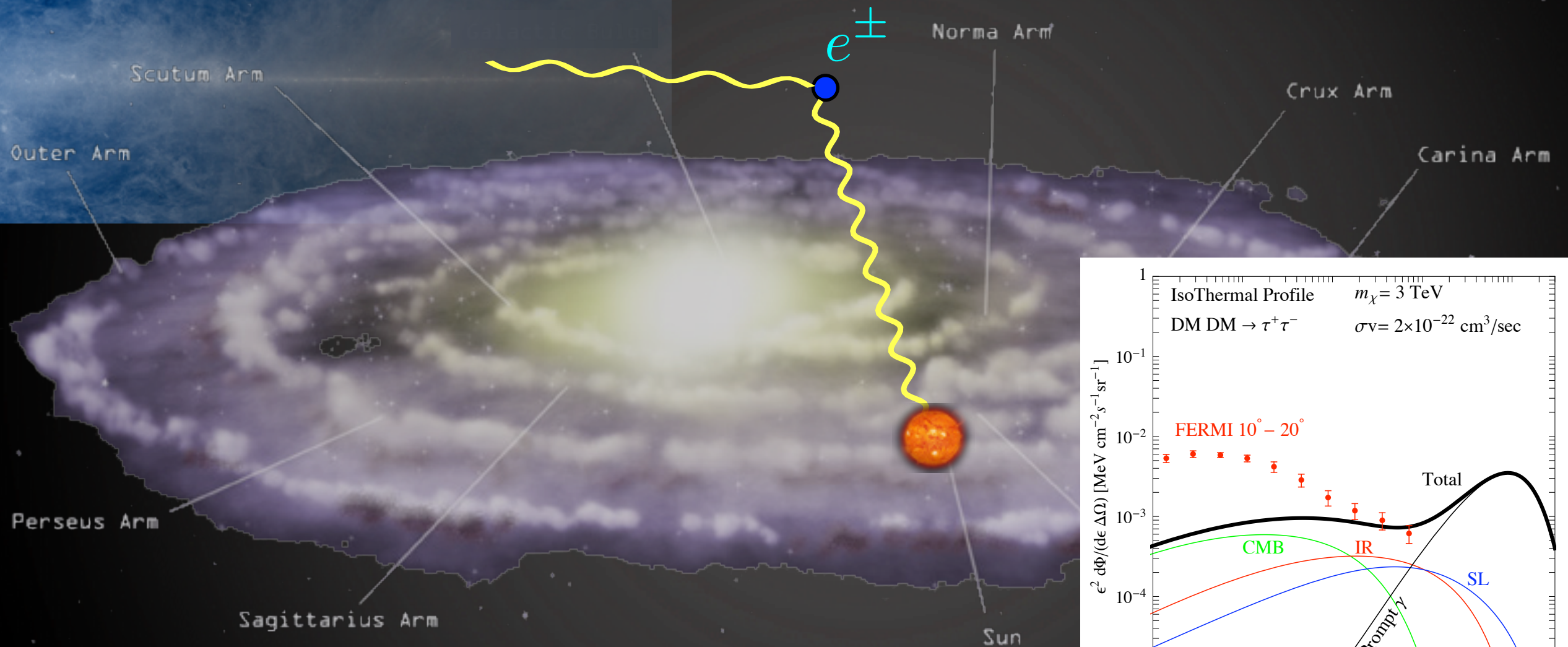


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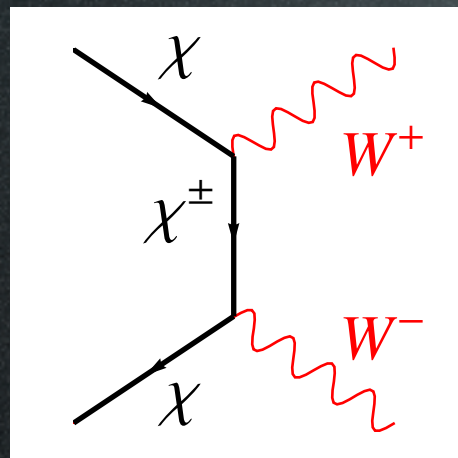
Star Light



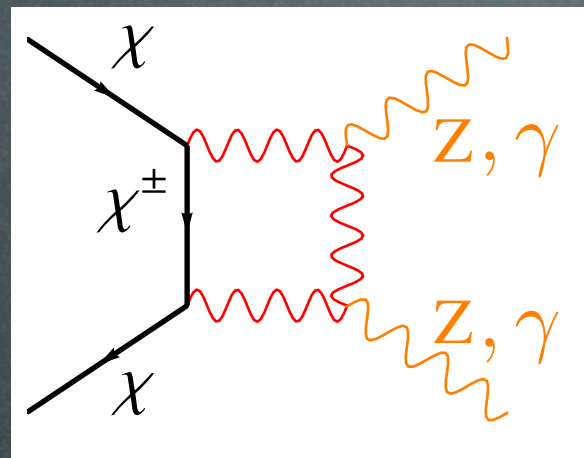
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Indirect Detection

γ from MDM annihilations



continuum



line(s)
(+ continuum)

+ $W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$

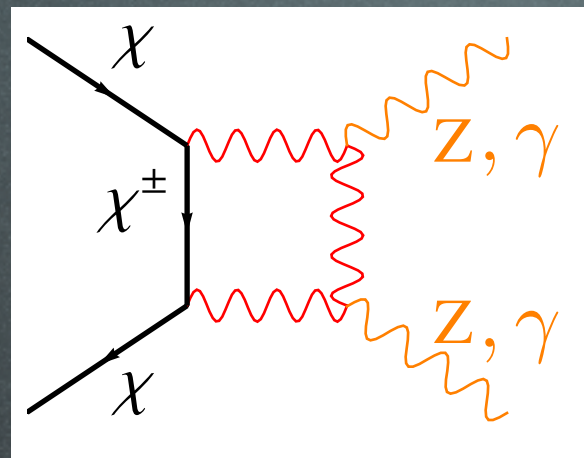
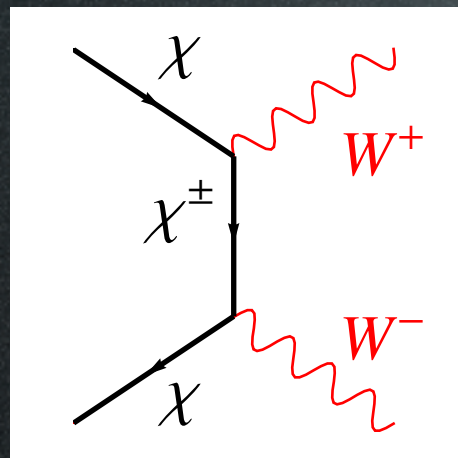
(channels for MDM with $Y=0$)

+ ICS



Indirect Detection

γ from MDM annihilations



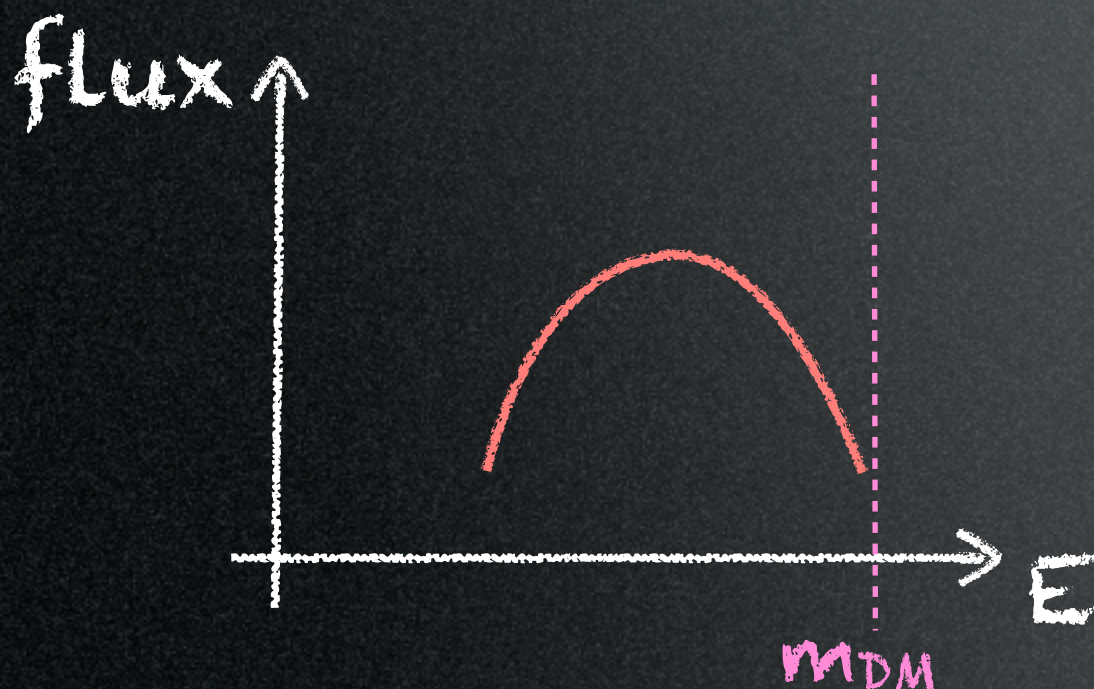
$$+ W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$$

(channels for MDM with $Y=0$)

continuum

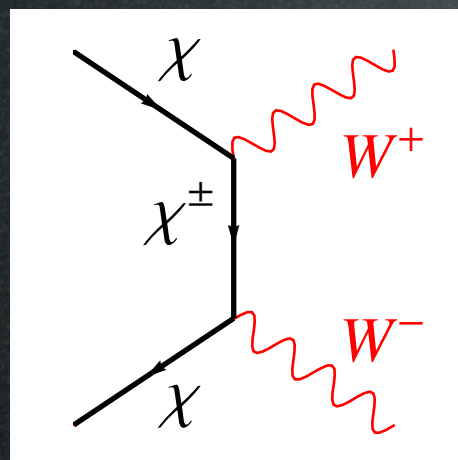
line(s)
(+ continuum)

+ ICS

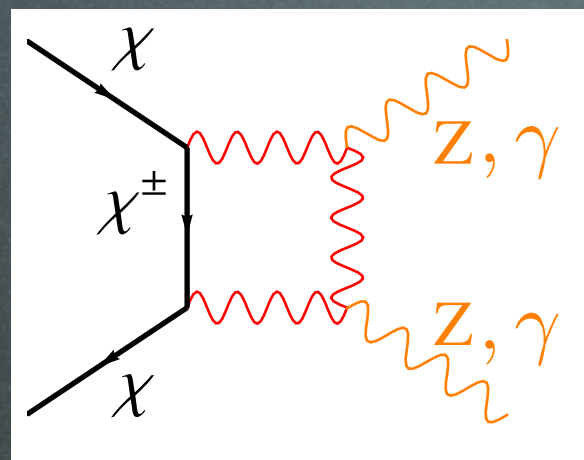


Indirect Detection

γ from MDM annihilations



continuum

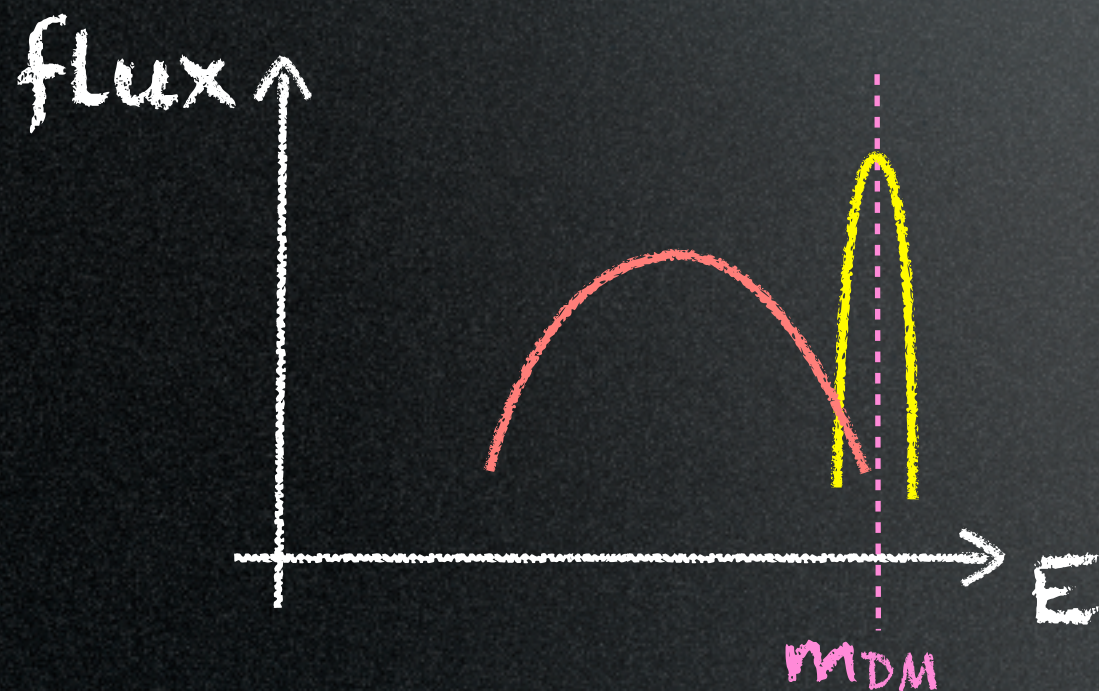


line(s)
(+ continuum)

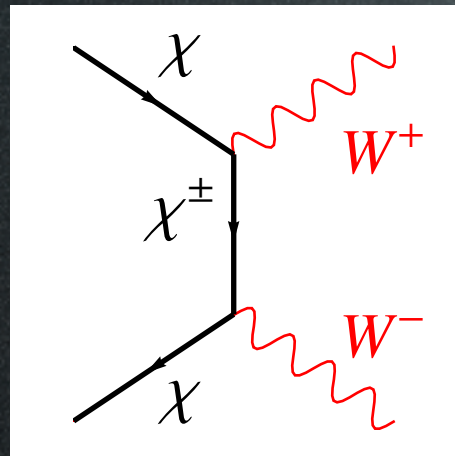
+ $W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$

(channels for MDM with $Y=0$)

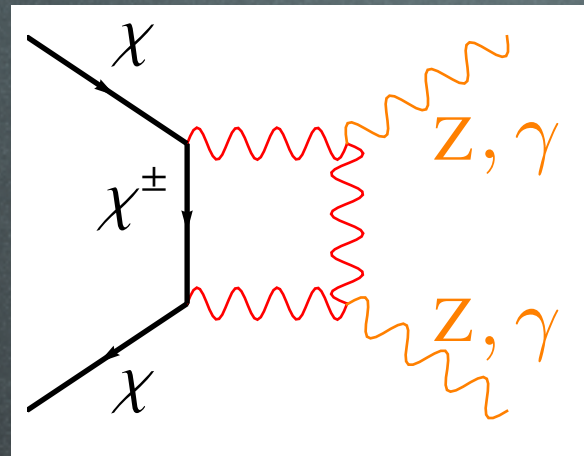
+ ICS



Indirect Detection



continuum

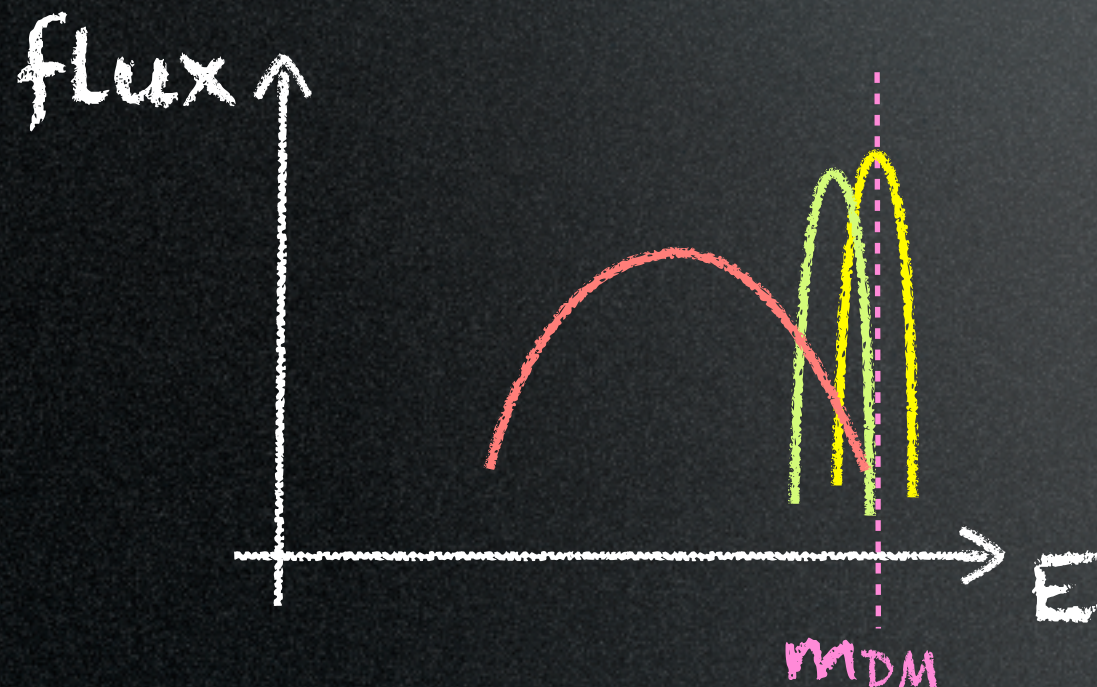


line(s)
(+ continuum)

$$+ W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$$

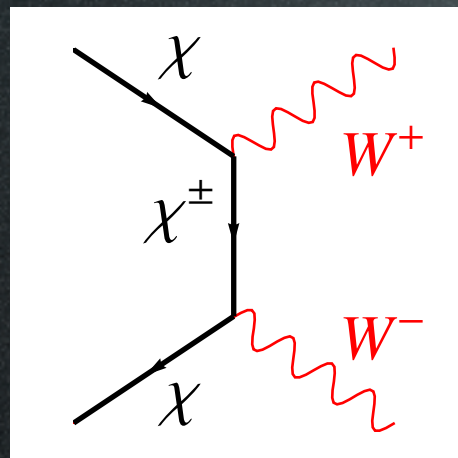
(channels for MDM with $Y=0$)

+ ICS

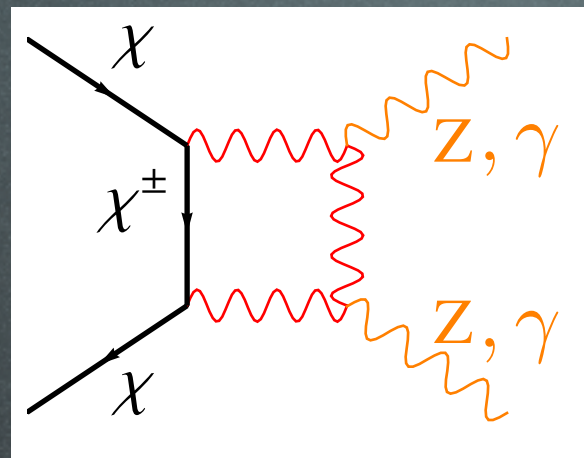


Indirect Detection

γ from MDM annihilations



continuum

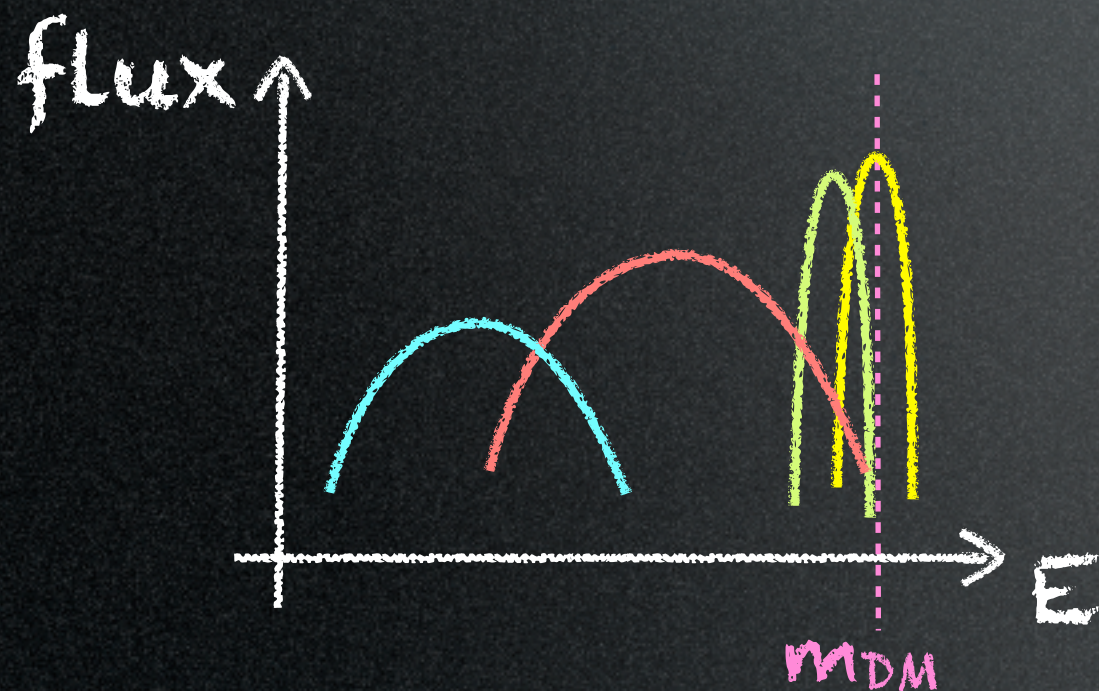


line(s)
(+ continuum)

+ $W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$

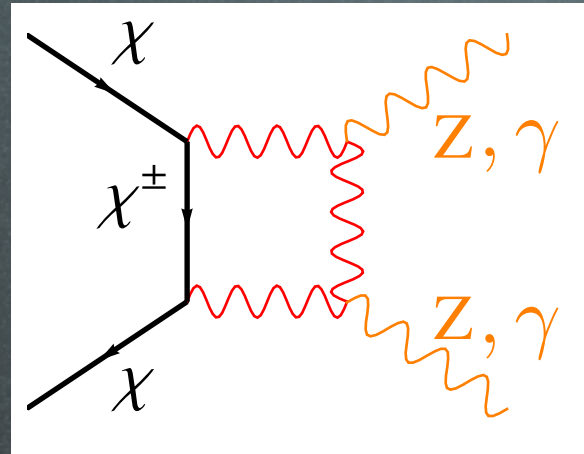
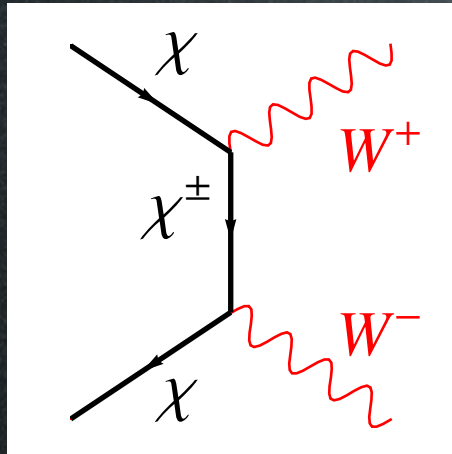
(channels for MDM with $Y=0$)

+ ICS



Indirect Detection

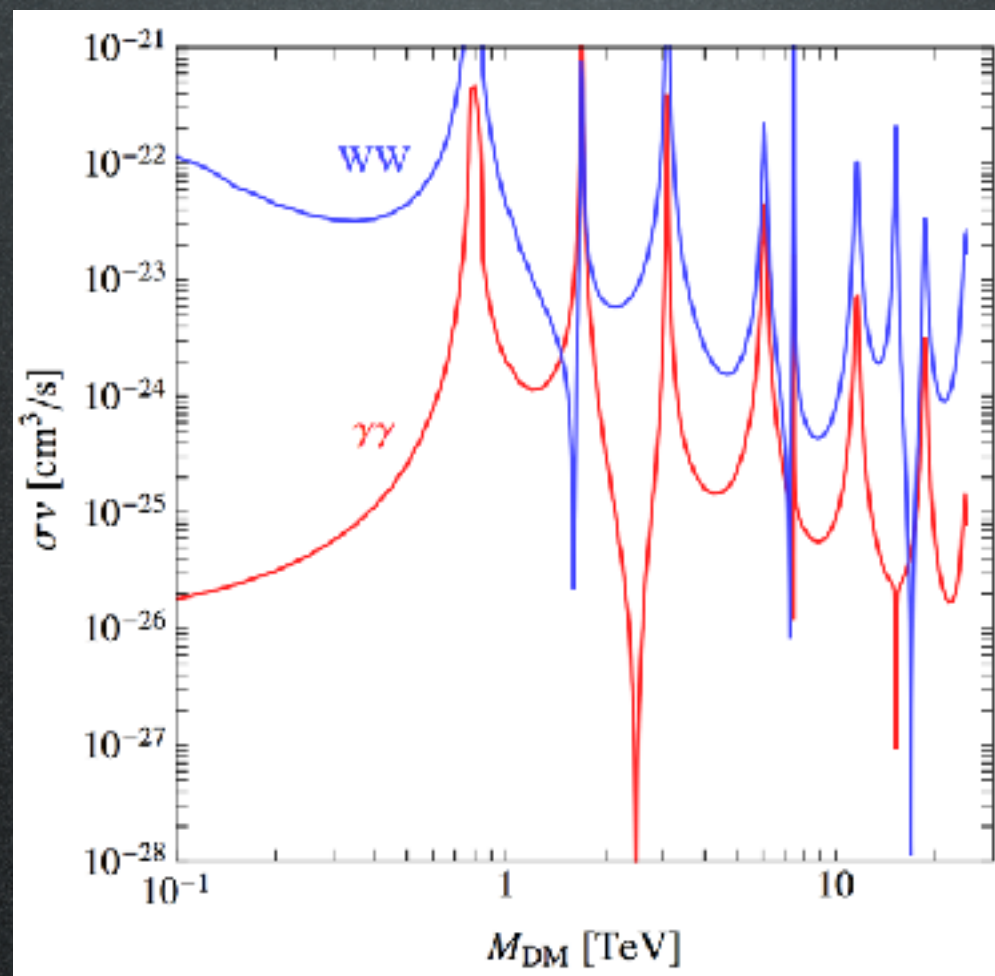
γ from MDM annihilations



$$+ W^\pm, Z \rightarrow \bar{p}, e^+, \gamma \dots$$

(channels for MDM with $Y=0$)

Enhanced cross section due to ‘Sommerfeld corrections’

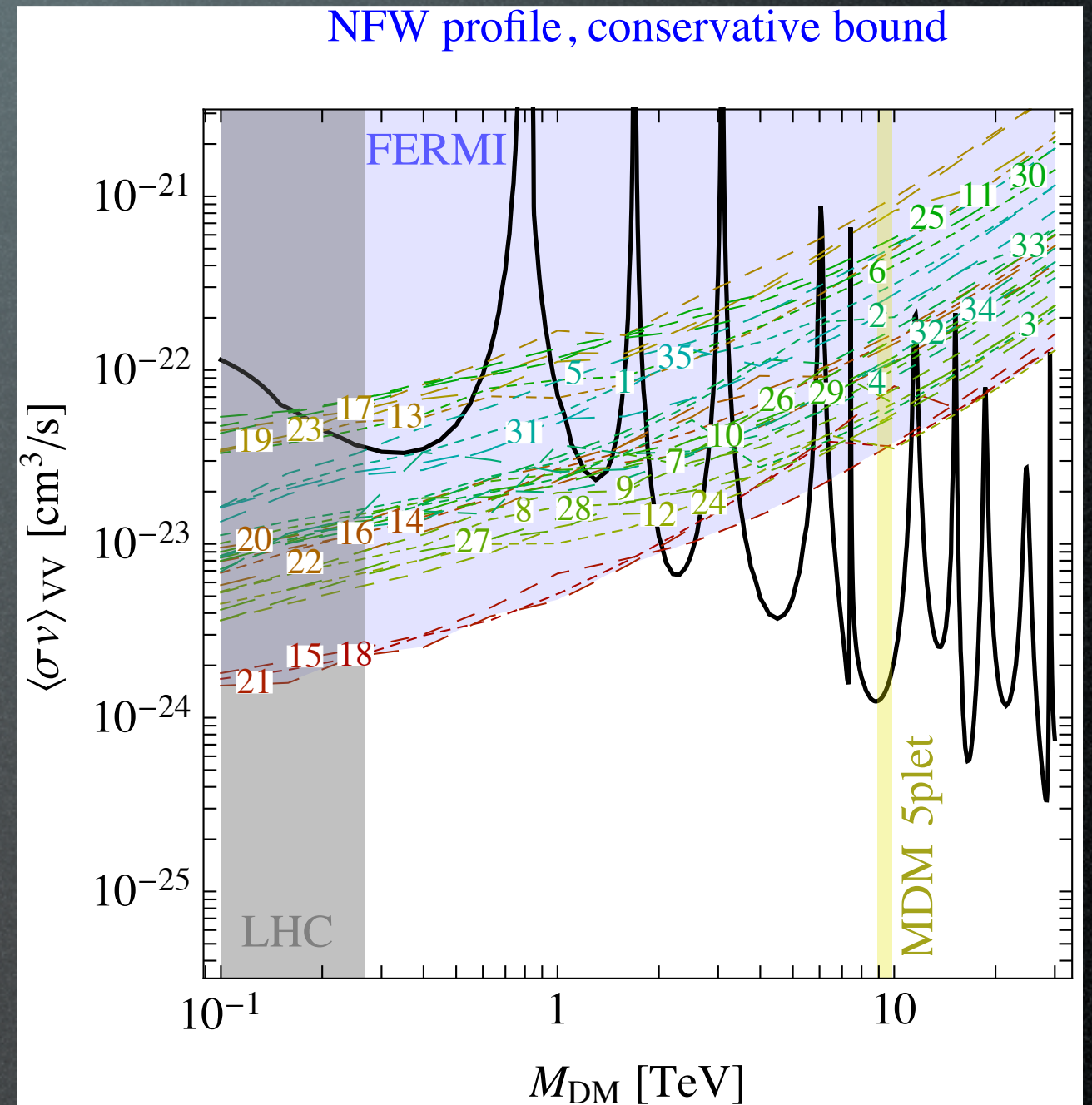
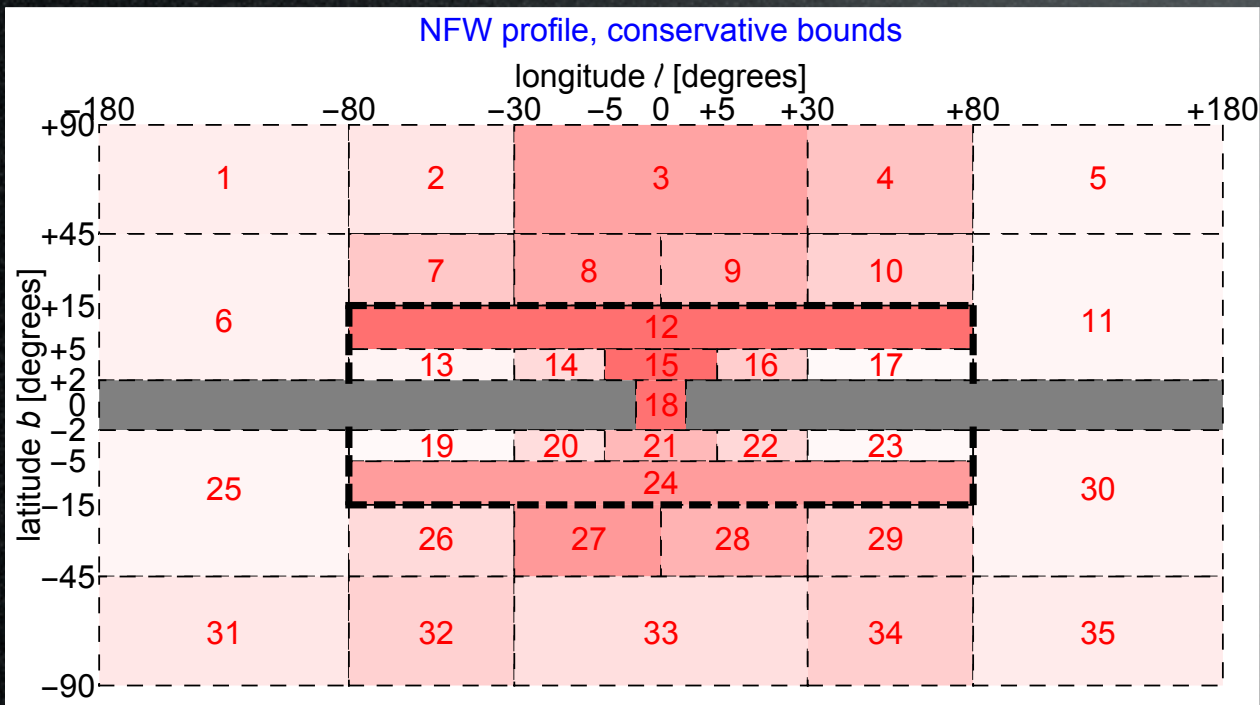
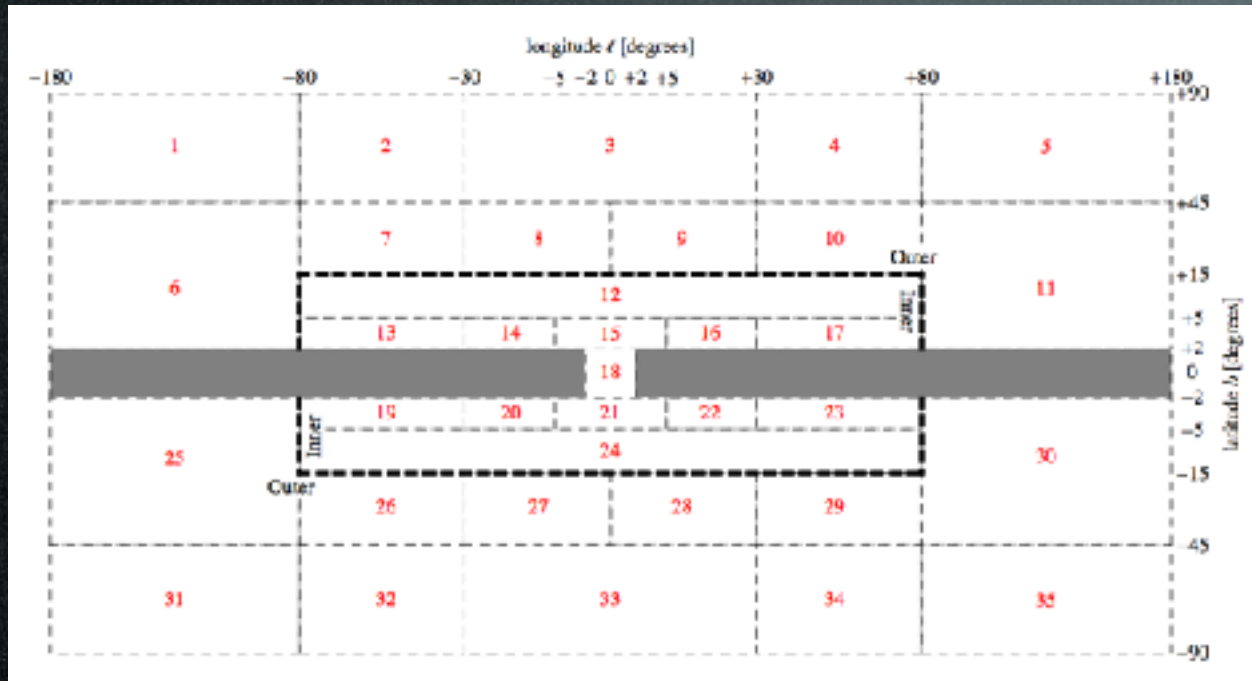


Hisano et al., 2004, 2005
Cirelli, Strumia, Tamburini 2007

Cirelli, Hambye, Panci, Sala, Taoso
1507.05519

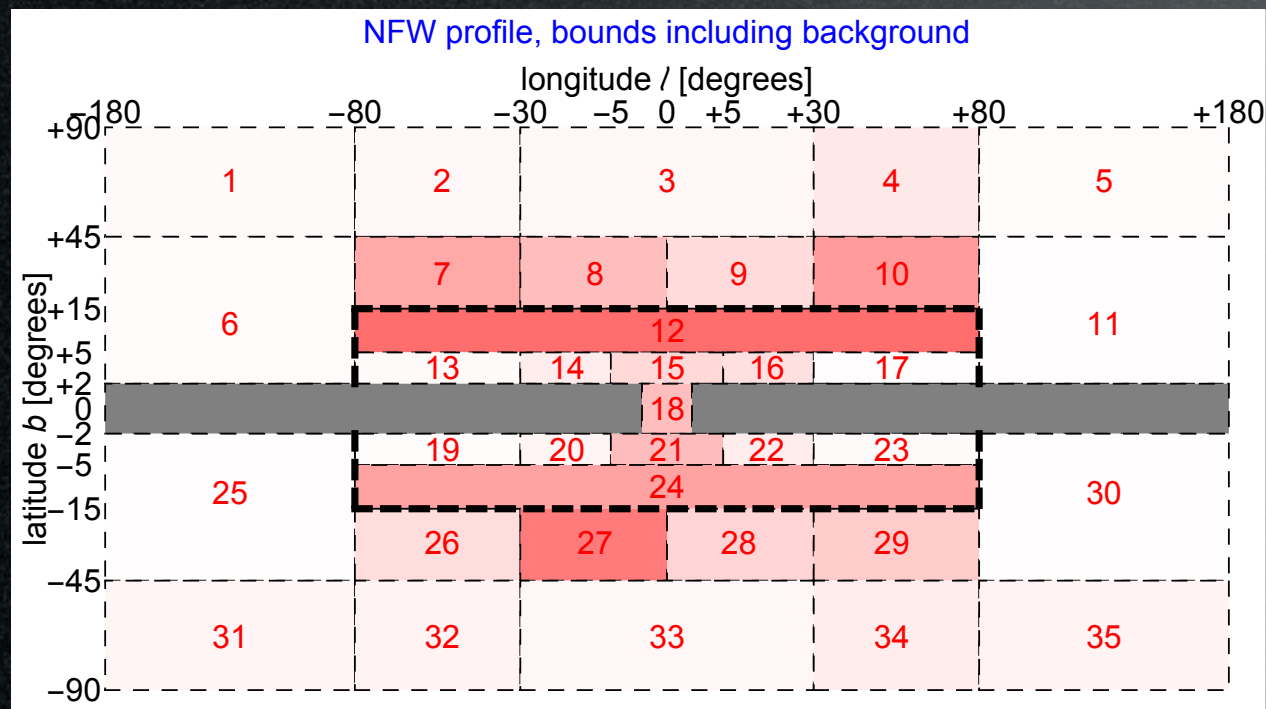
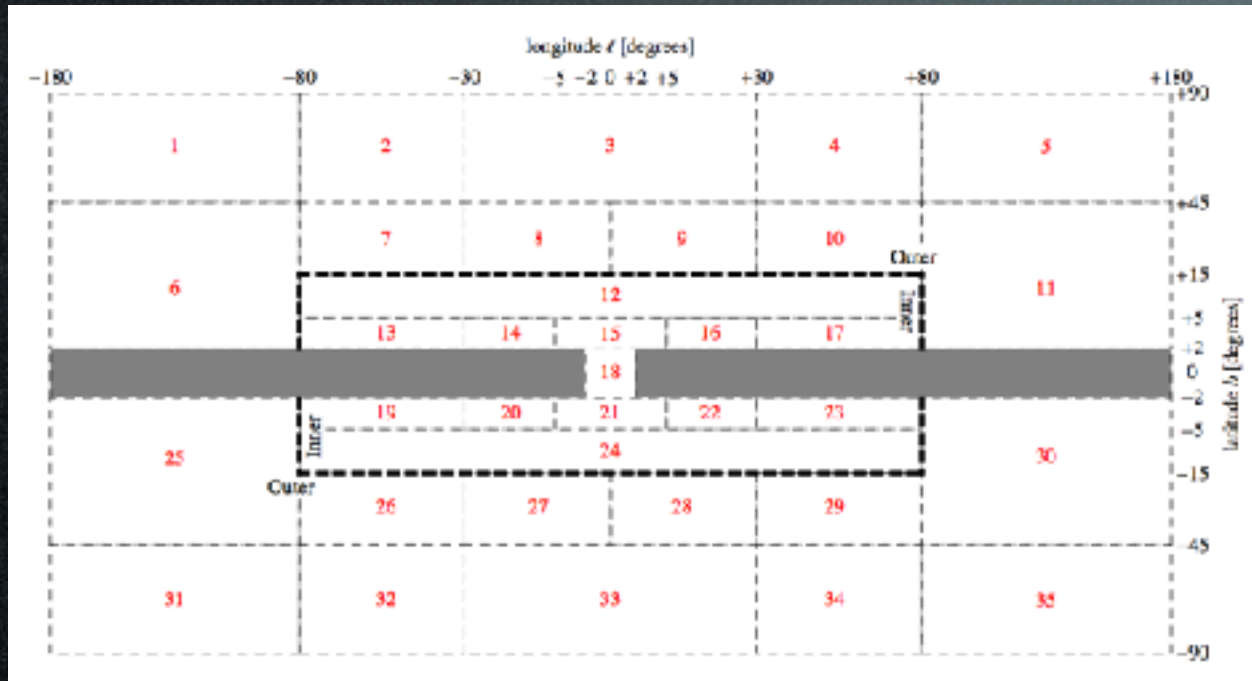
Indirect Detection

FERMI diffuse galactic:

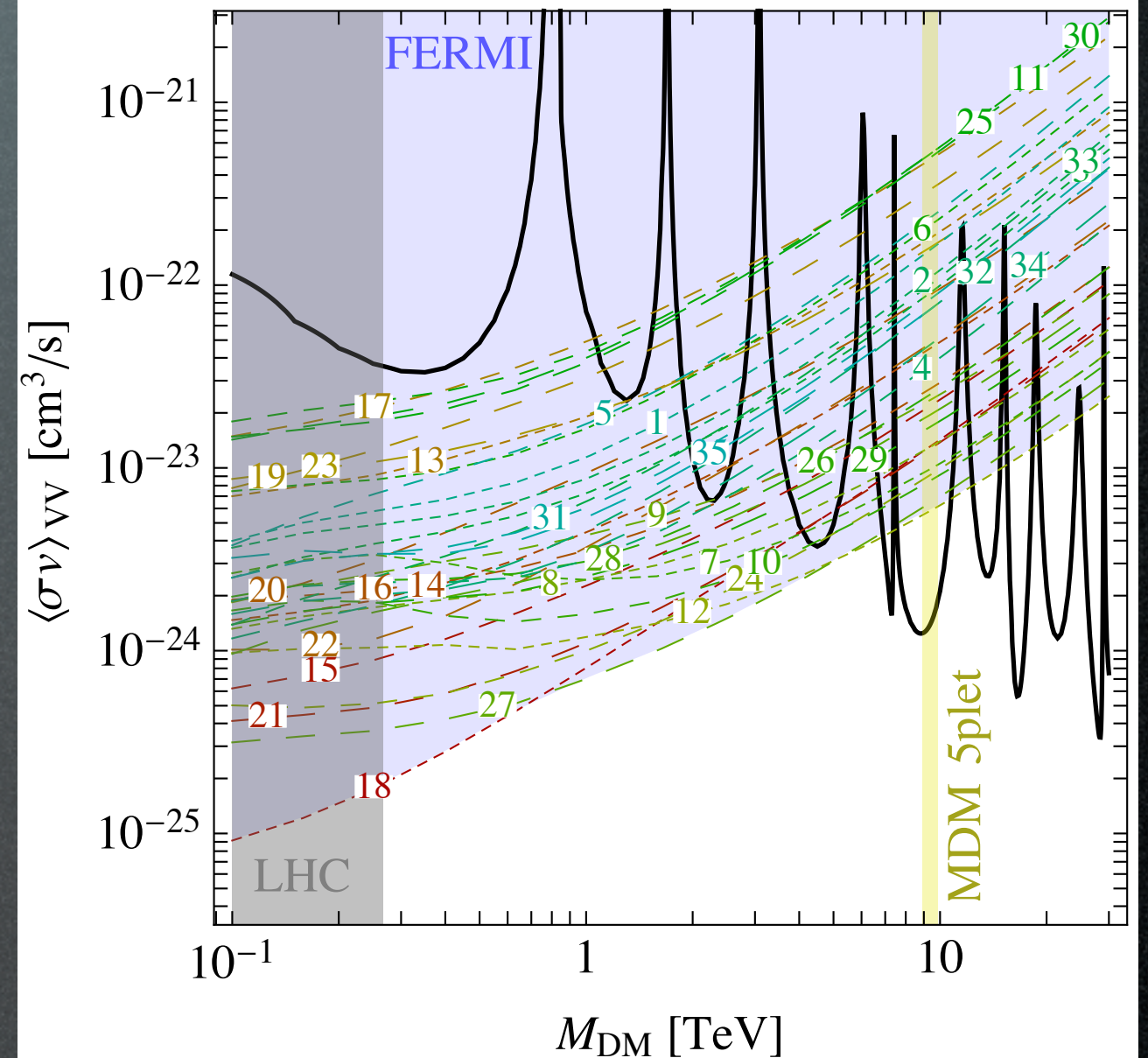


Indirect Detection

FERMI diffuse galactic:

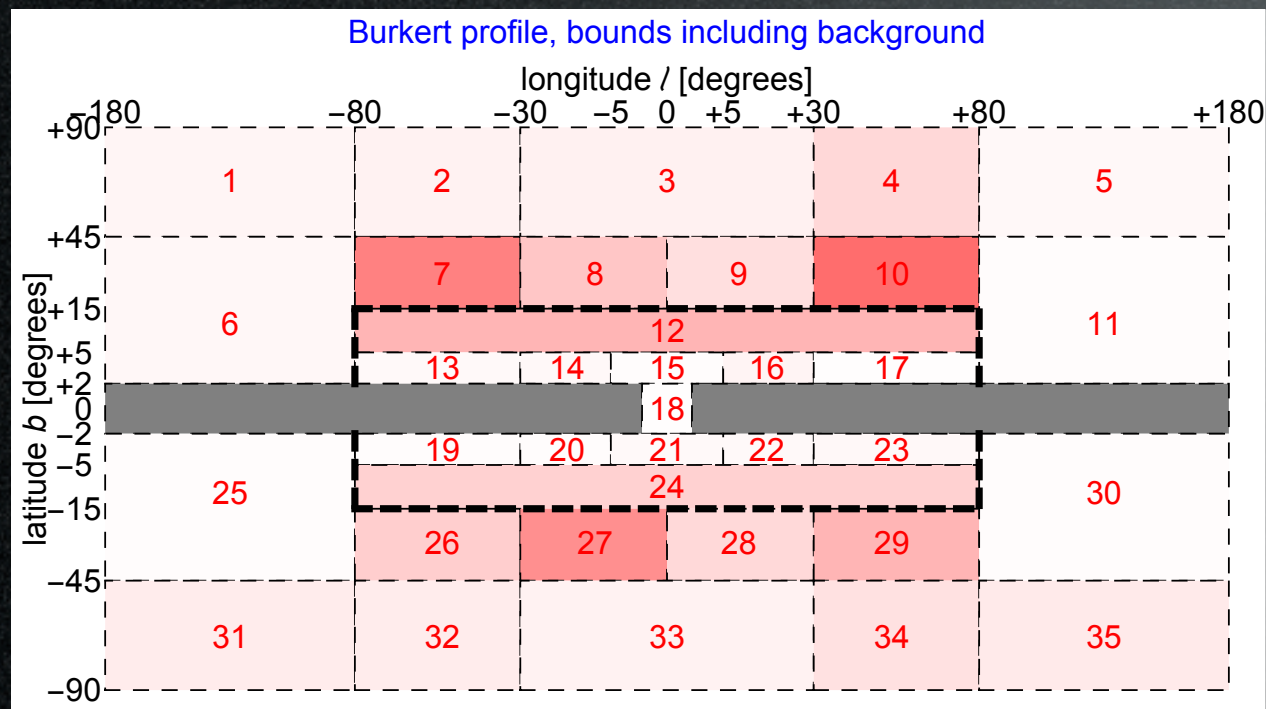
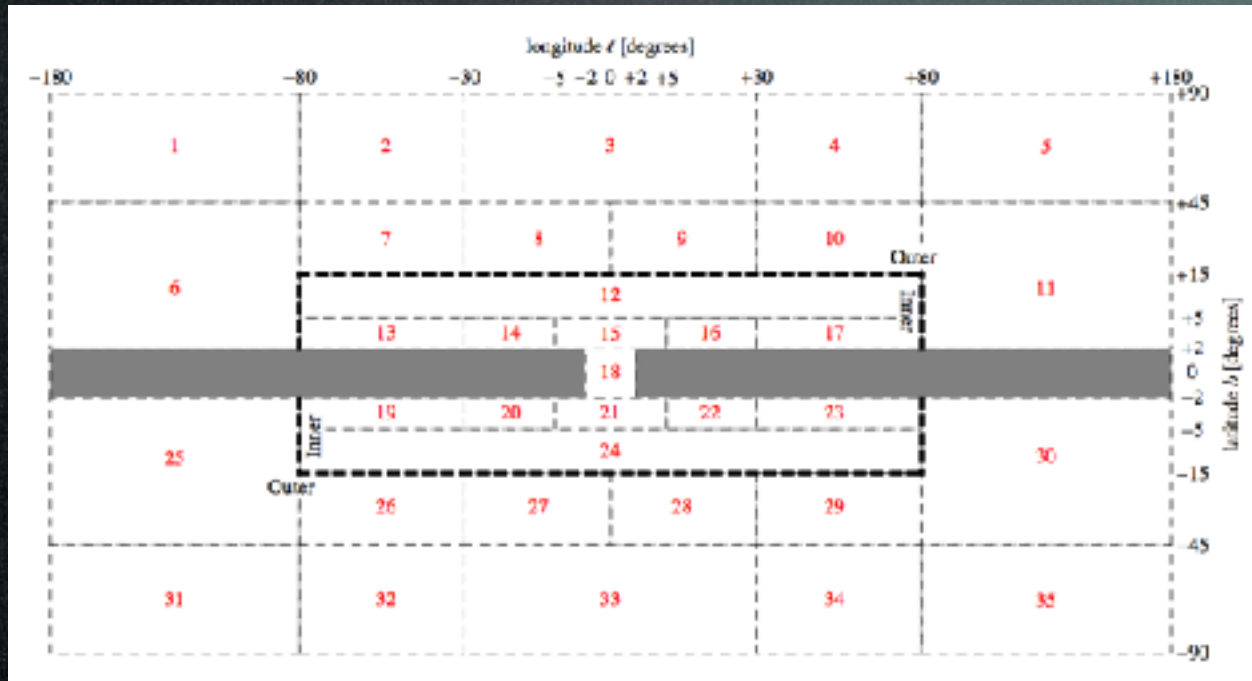


NFW profile, including background

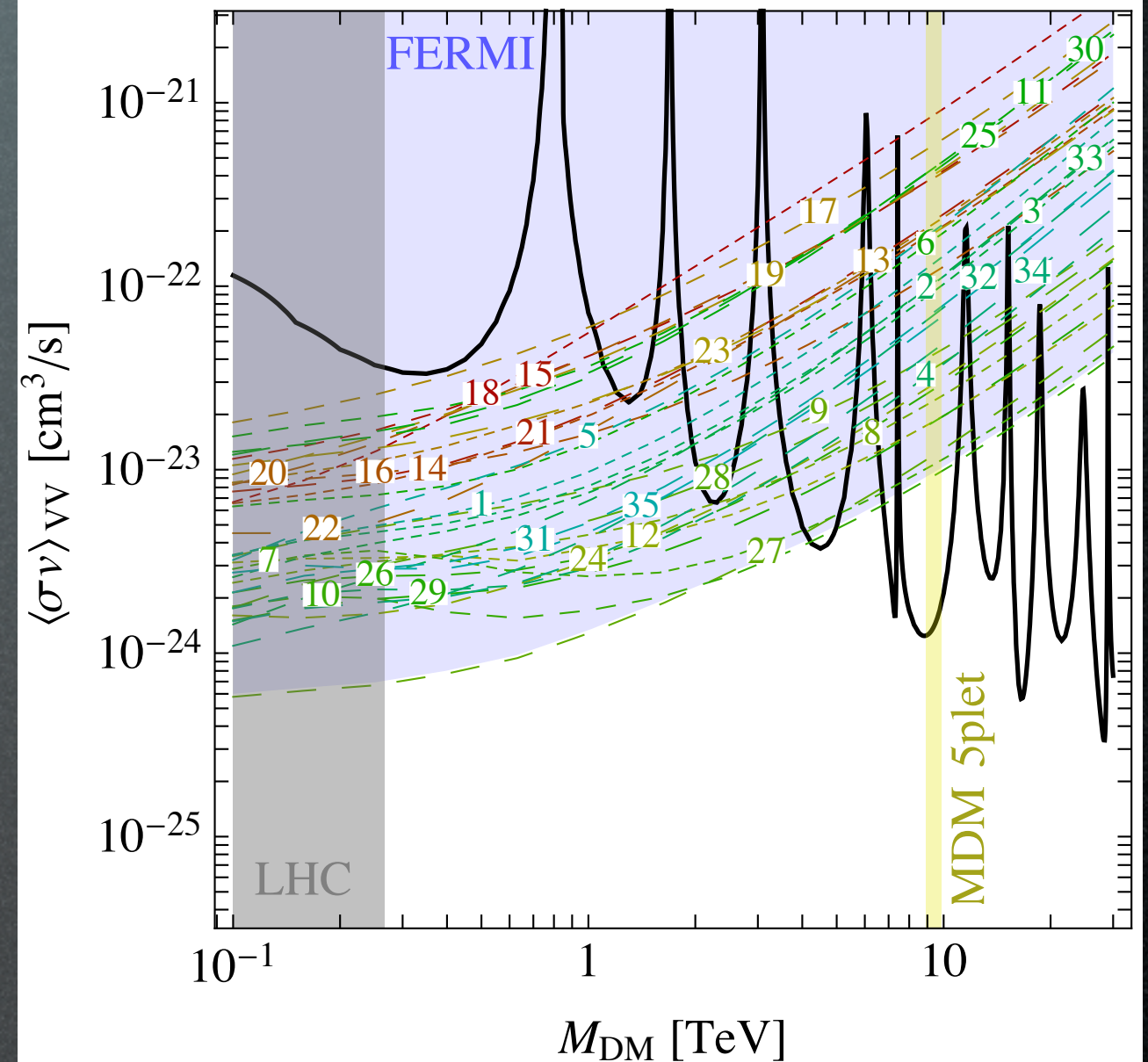


Indirect Detection

FERMI diffuse galactic:

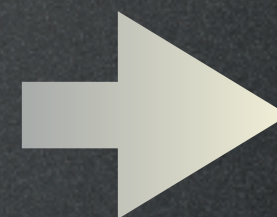
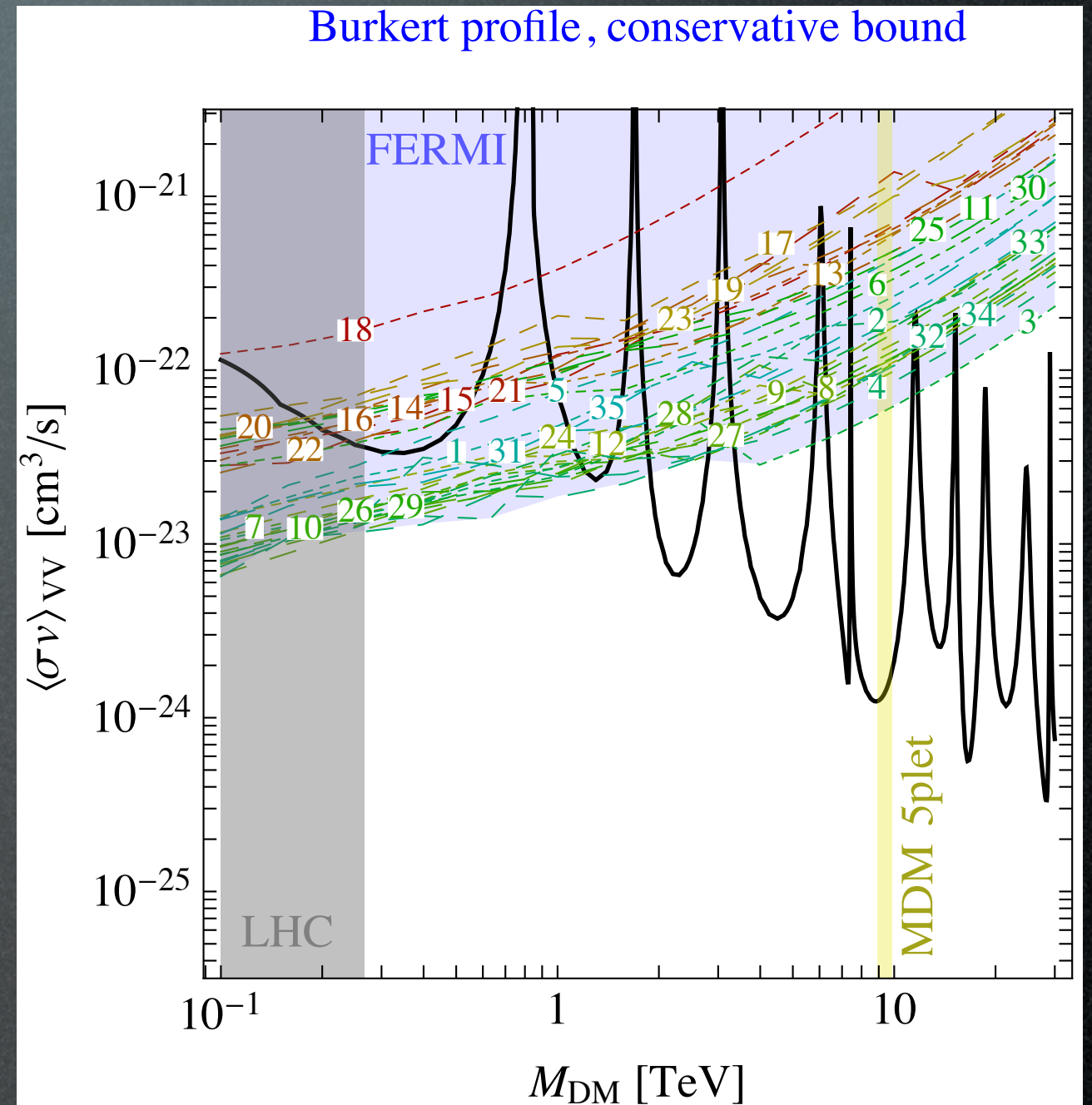
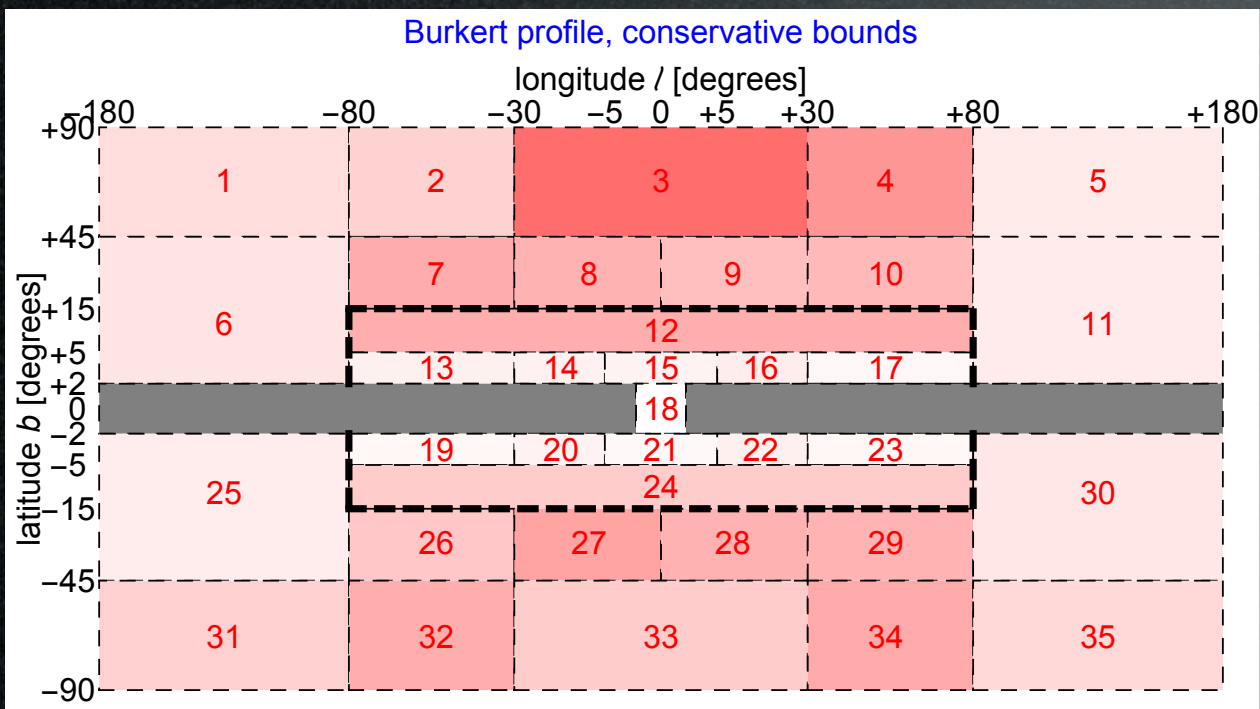
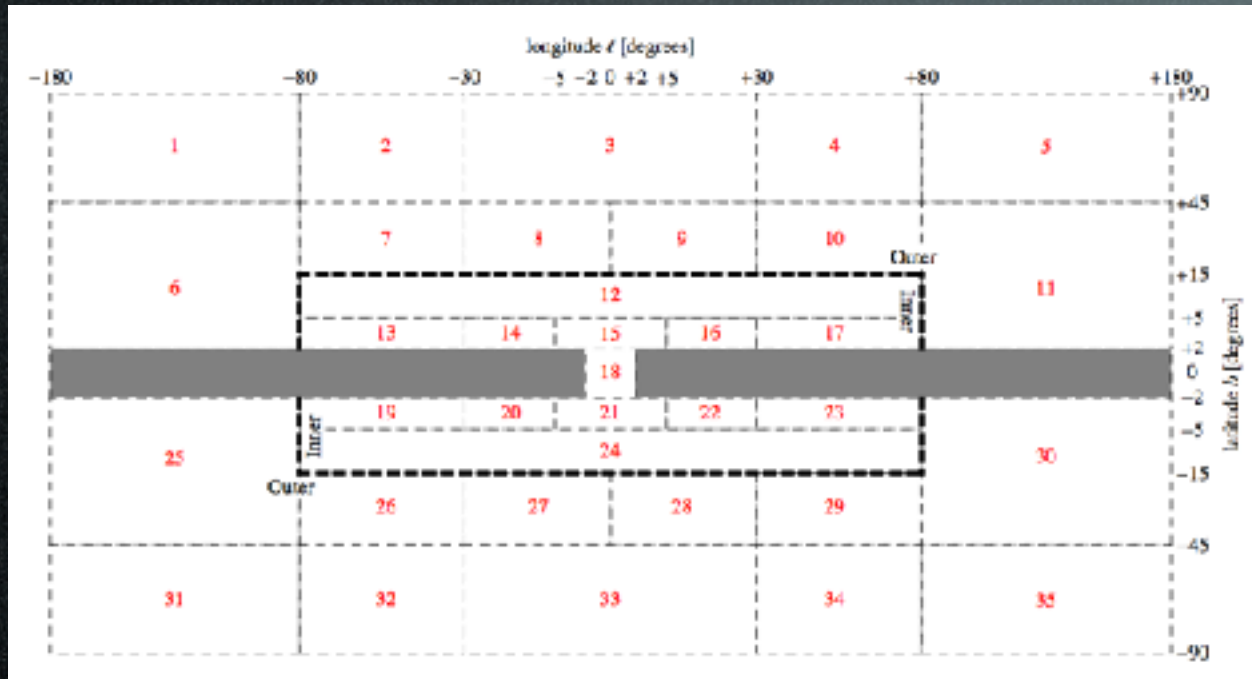


Burkert profile, including background



Indirect Detection

FERMI diffuse galactic:



relevant constraints but
MDM 5plet not probed

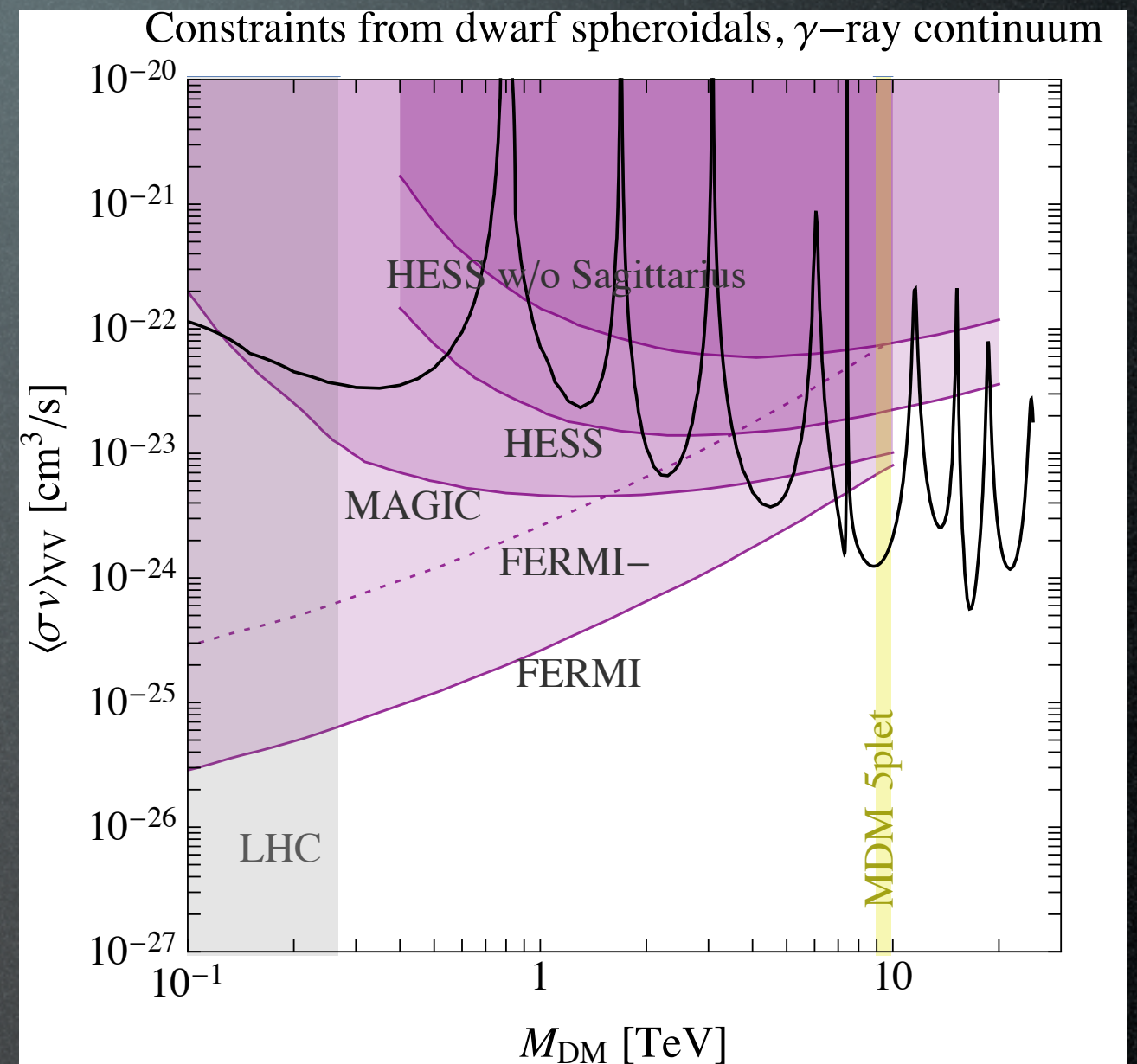
Indirect Detection

dSphs galaxies, search for **continuum** γ -rays:

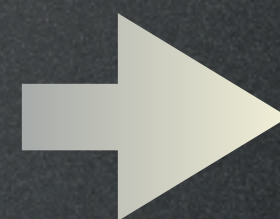
FERMI: 15 dSphs, 6yrs, 'Pass-8' - 1503.02641

HESS: 4 dSphs, incl Sagittarius - 1410.2589

MAGIC: Segue1 - 1312.1535



Cirelli, Hambye, Panci, Sala, Taoso 1507.05519



relevant constraints but
MDM 5plet not probed

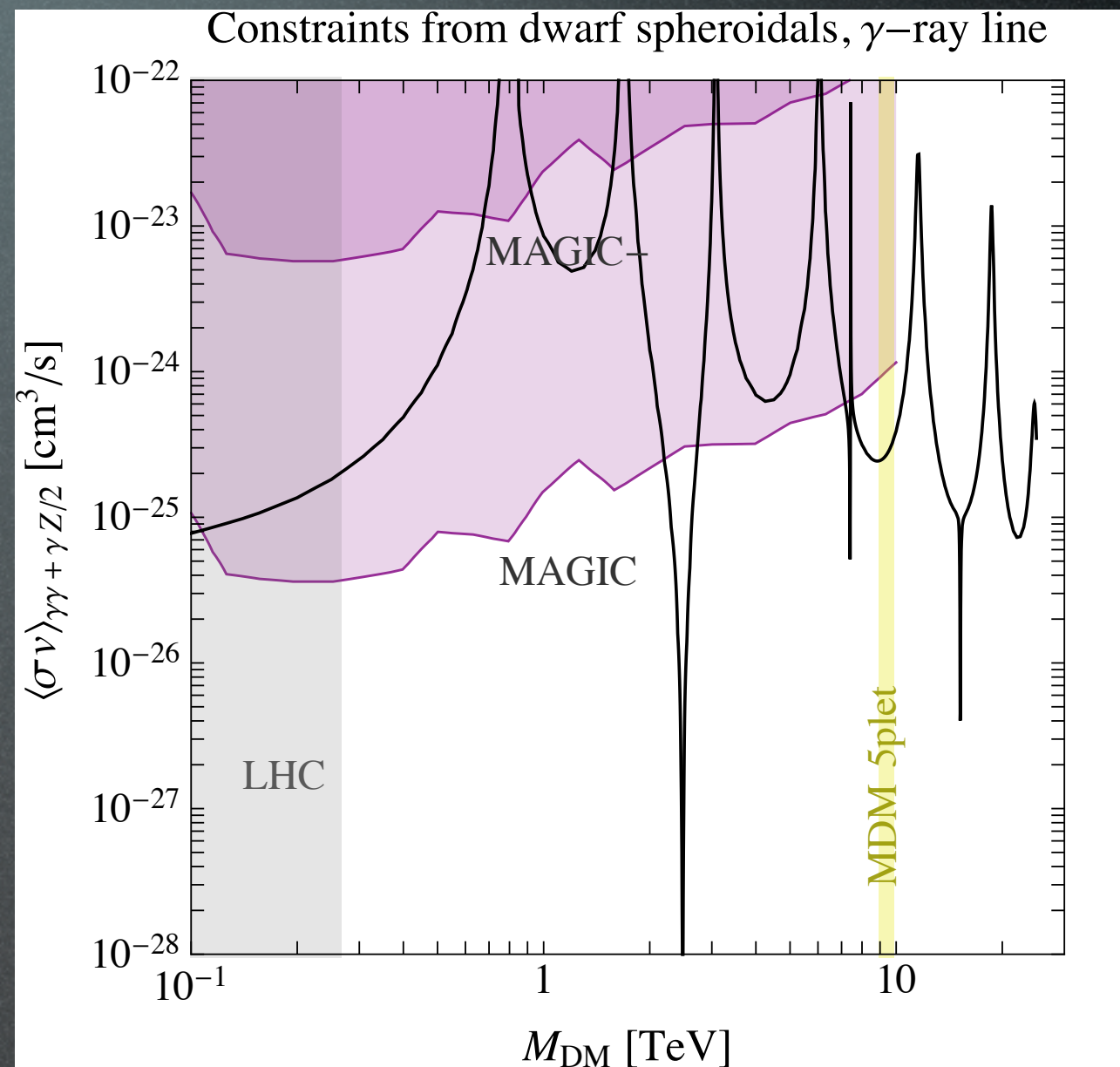
Indirect Detection

dSphs galaxies, search for γ -ray lines:

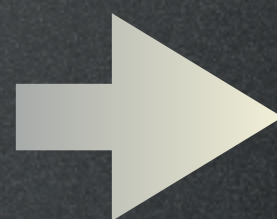
MAGIC: Segue1 - 1312.1535

NB large uncertainties in dSPhs
'J factor', i.e. DM-brightness

e.g. Bonnavard et al., 1504.02048



Cirelli, Hambye, Panci, Sala, Taoso 1507.05519



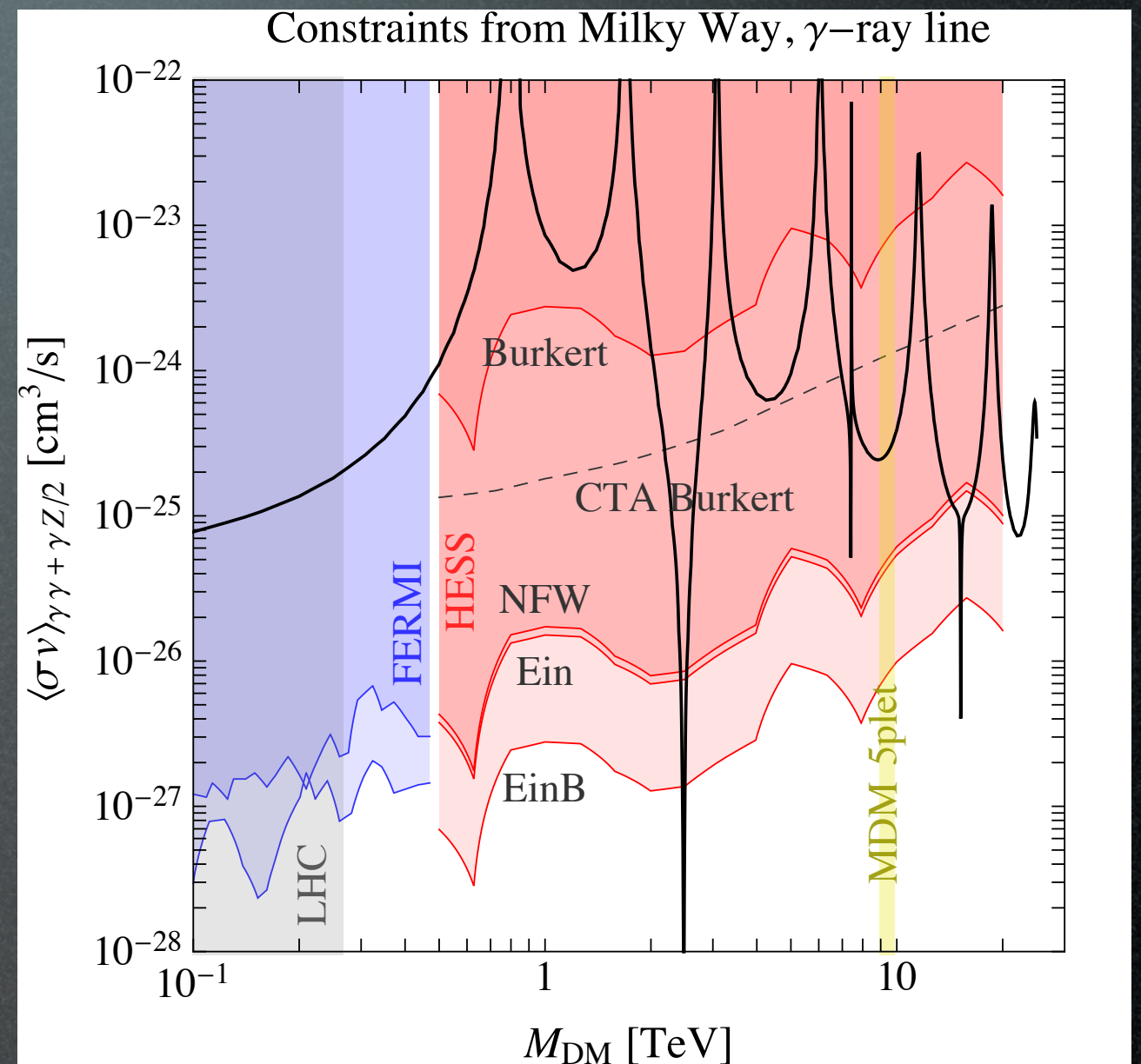
relevant constraints but
MDM 5plet not probed

Indirect Detection

MW center area, search for γ -ray lines:

FERMI: 1506.00013

HESS: 1301.1173



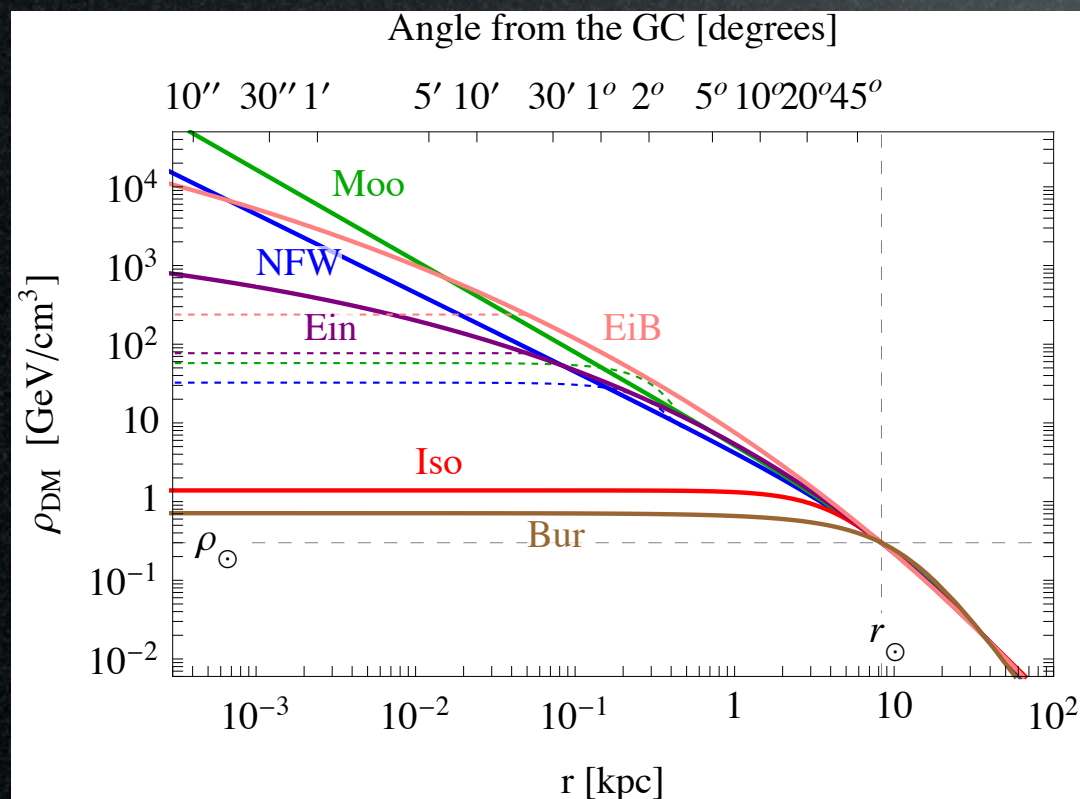
Indirect Detection

MW center area, search for γ -ray lines:

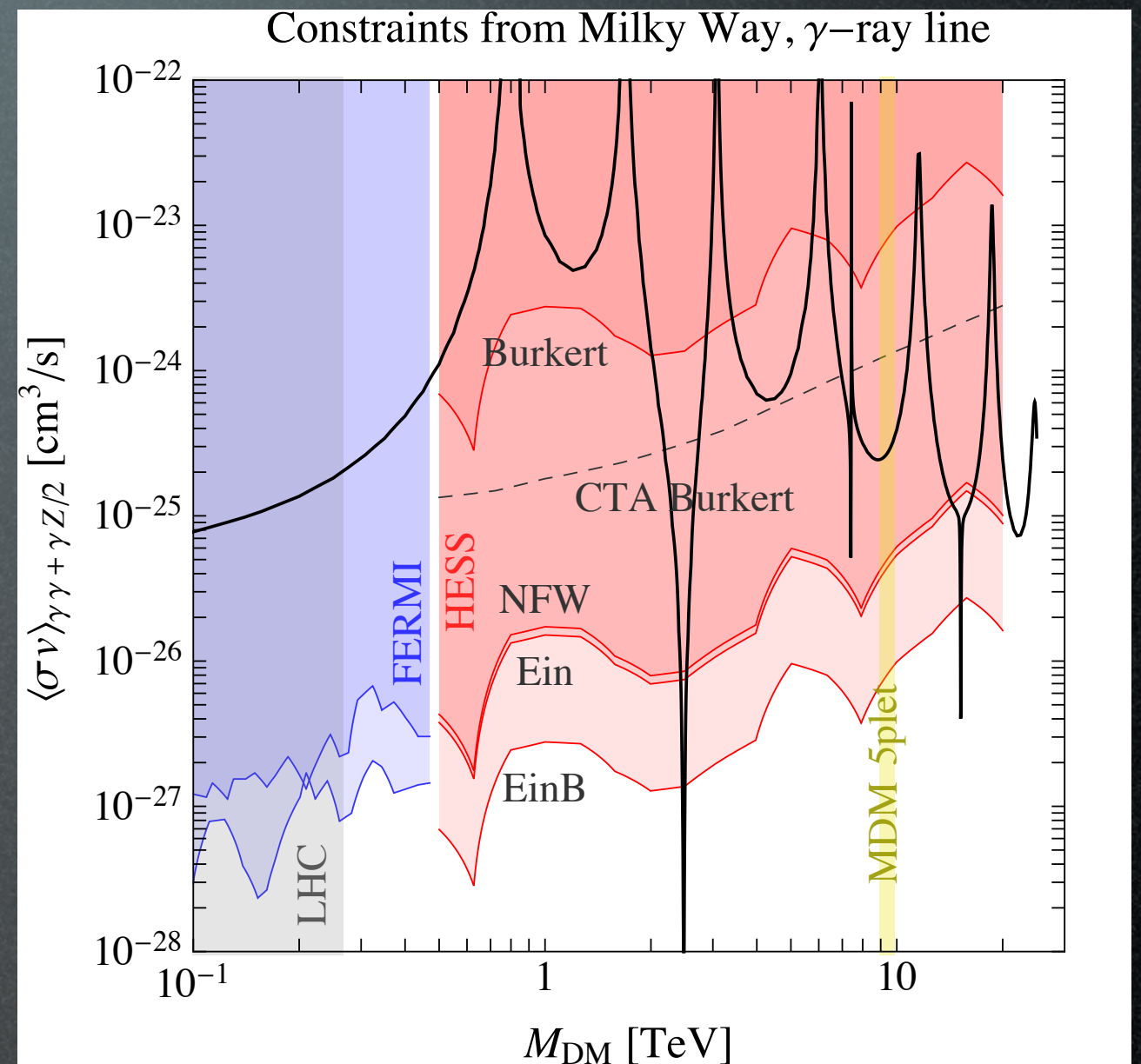
FERMI: 1506.00013

HESS: 1301.1173

Uncertainties in DM profile:



e.g. Cirelli et al., 1012.4515



Cirelli, Hambye, Panci, Sala, Taoso 1507.05519

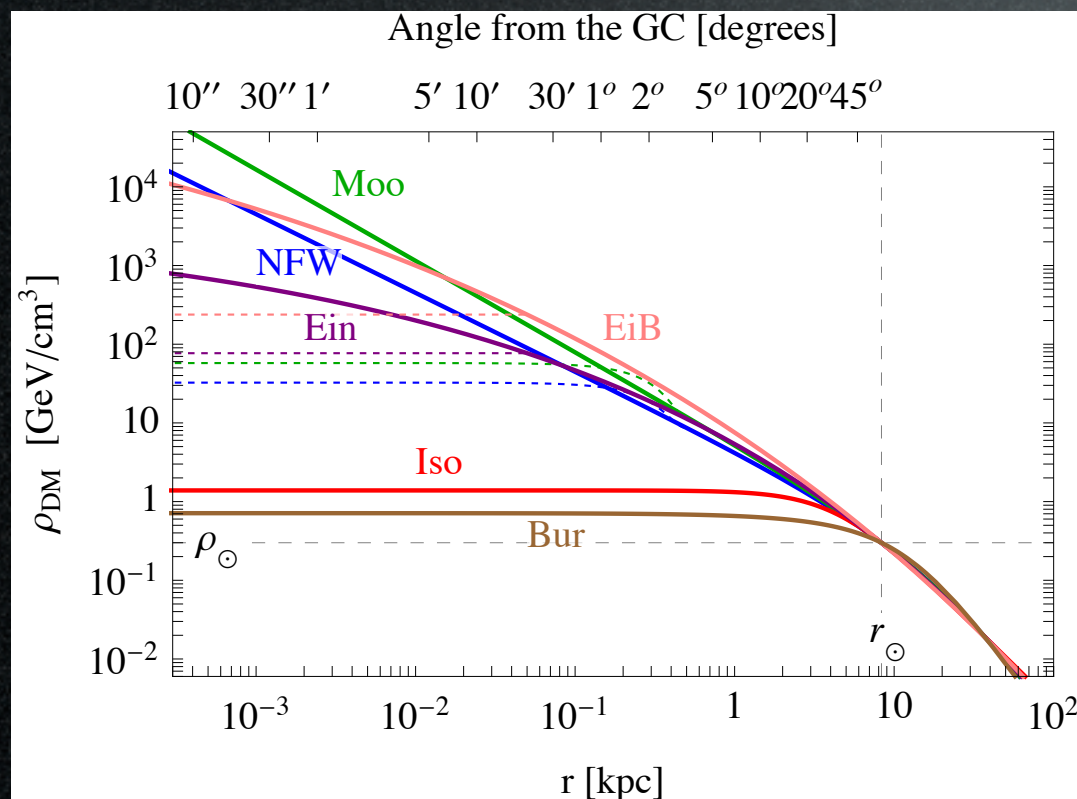
Indirect Detection

MW center area, search for γ -ray lines:

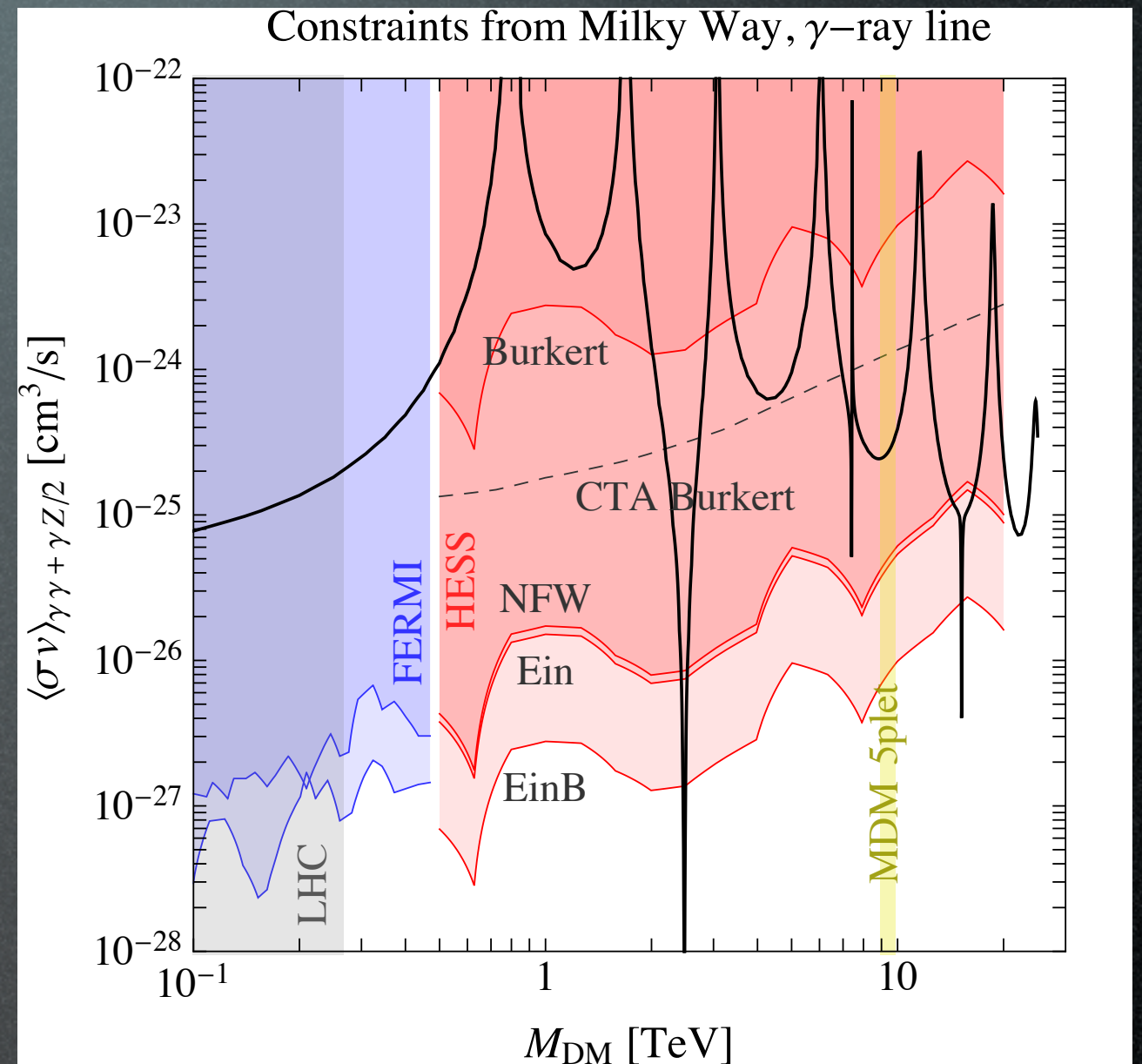
FERMI: 1506.00013

HESS: 1301.1173

Uncertainties in DM profile:



e.g. Cirelli et al., 1012.4515



Cirelli, Hambye, Panci, Sala, Taoso 1507.05519

MDM excluded if cuspy
MDM not probed if cored

Consistent conclusions in: Garcia-Cely et al. 1507.05536

Post scriptum

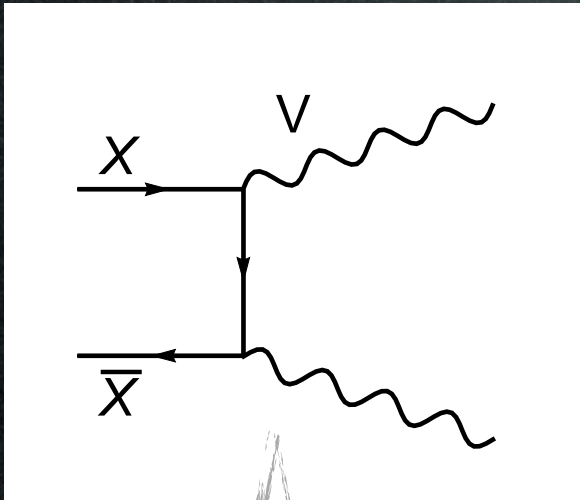
Bound state formation is relevant

Pospelov, Ritz 2009
March-Russell, West 2009
Shepherd, Tait, Zaharijas 2009
K.Petraki+, 2014+

Post scriptum

Bound state formation is relevant

Pospelov, Ritz 2009
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K.Petraki+, 2014+

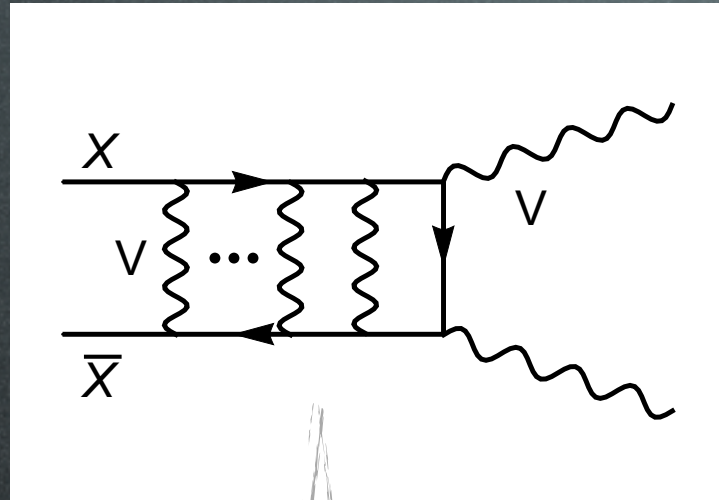
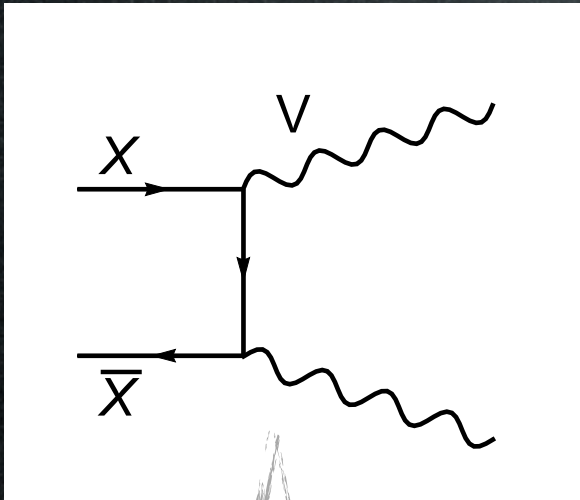


tree level annihilation

Post scriptum

Bound state formation is relevant

Pospelov, Ritz 2009
March-Russell, West 2009
Shepherd, Tait, Zaharijas 2009
K.Petraki+, 2014+



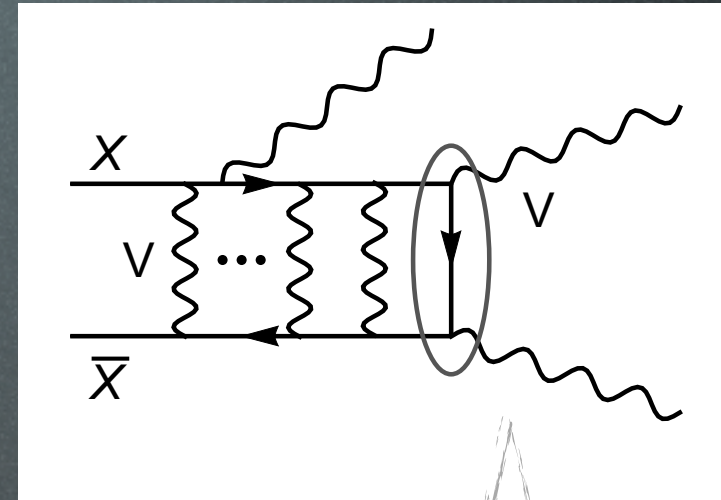
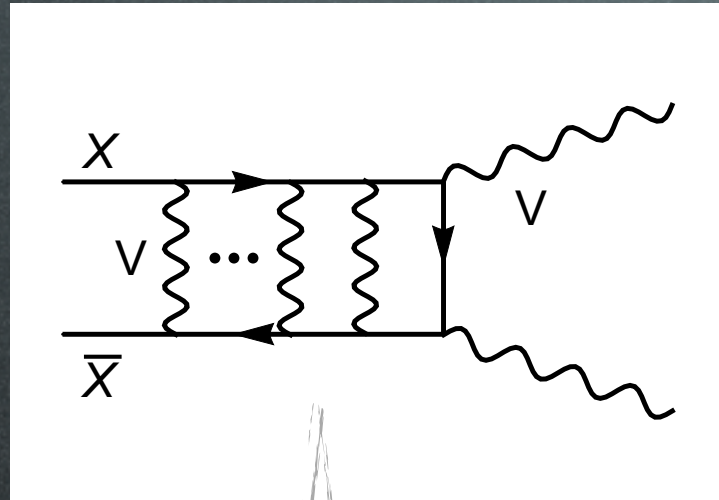
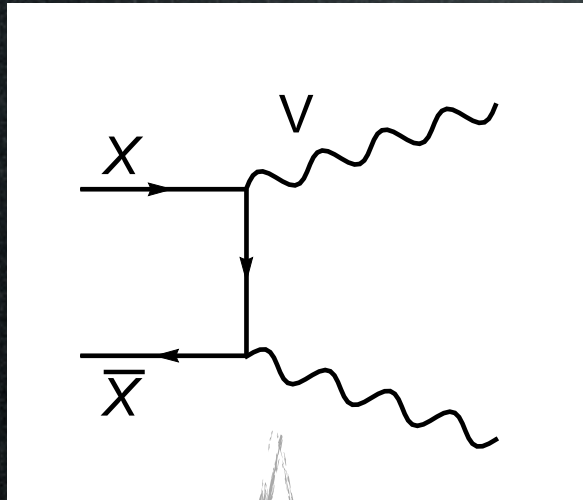
tree level annihilation

size of the XX system
If $\alpha M/m_V \gtrsim 1$, the force is long range:
Sommerfeld enhanced
range

Post scriptum

Bound state formation is relevant

Pospelov, Ritz 2009
March-Russell, West 2009
Shepherd, Tait, Zaharijas 2009
K.Petraki+, 2014+



tree level annihilation

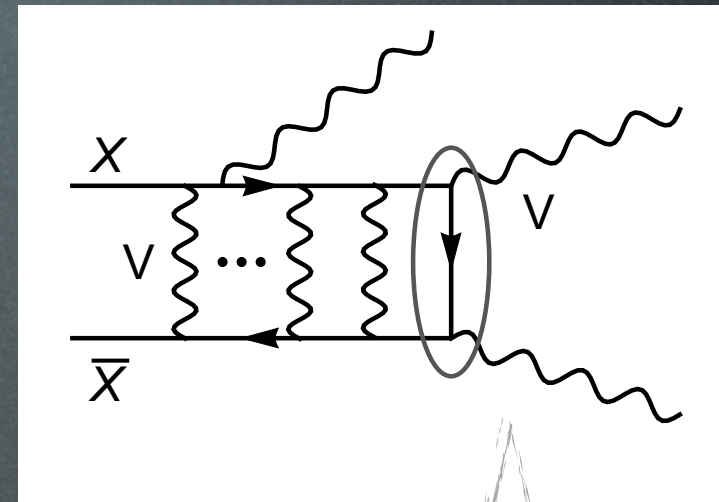
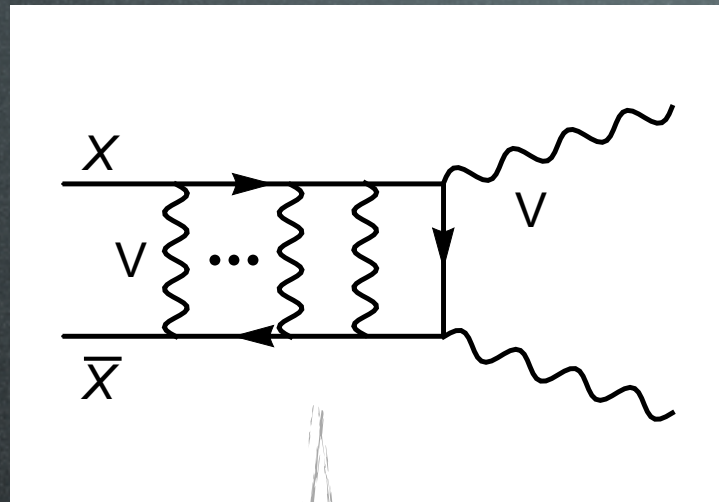
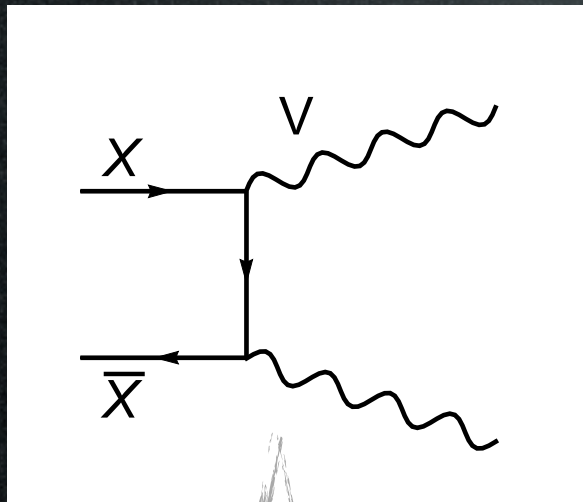
binding energy of the XX system
If $\alpha^2 \overset{\text{binding energy of the } XX \text{ system}}{M}/2m_V \gtrsim 1$, bound states form
 $\overset{\text{emitted mediator}}{\nearrow}$

size of the XX system
If $\alpha \overset{\text{size of the } XX \text{ system}}{M}/m_V \gtrsim 1$, the force is long range:
 $\nearrow$ range Sommerfeld enhanced

Post scriptum

Bound state formation is relevant

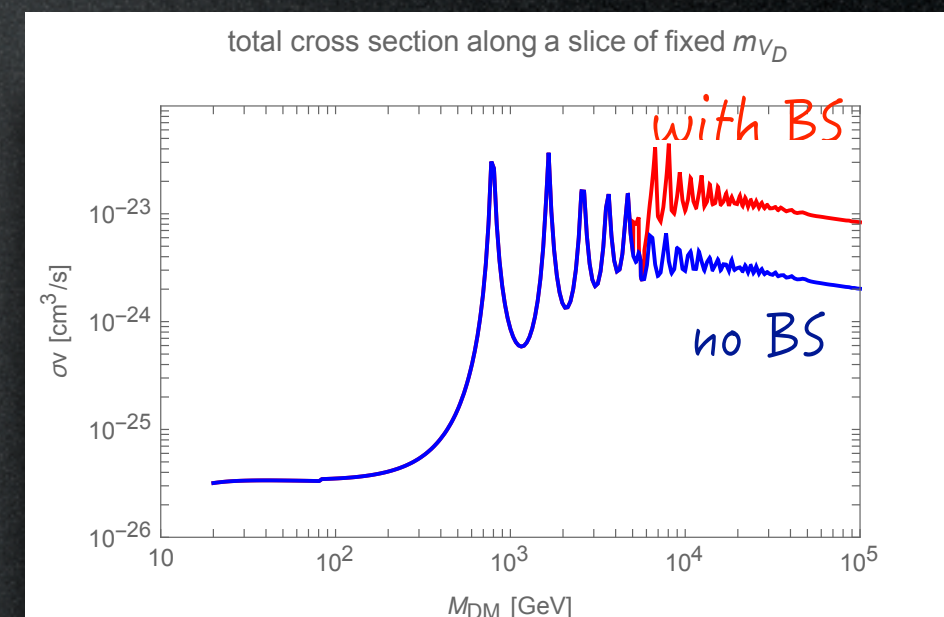
Pospelov, Ritz 2009
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Shepherd, Tait, Zaharijas 2009
K.Petraki+, 2014+



tree level annihilation

binding energy of the XX system
If $\alpha^2 M/2m_V \gtrsim 1$, bound states form
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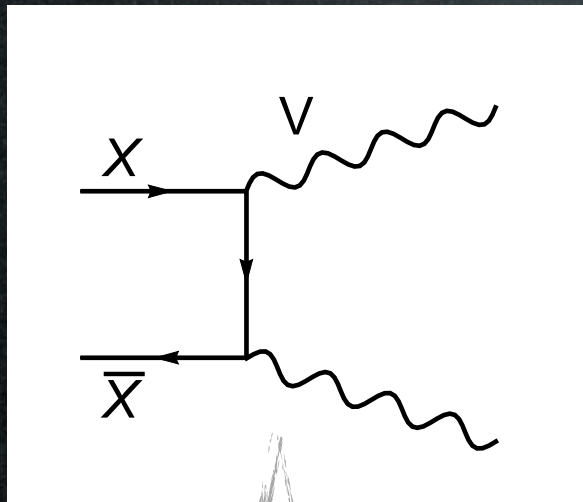
size of the XX system
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range
Sommerfeld enhanced



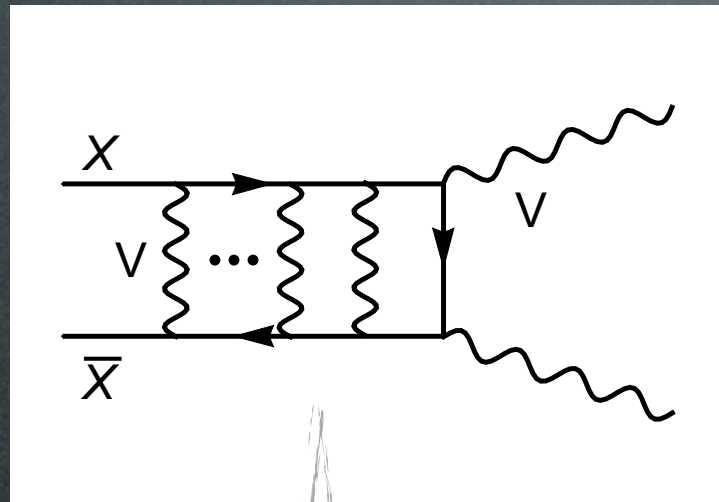
Post scriptum

Bound state formation is relevant

Pospelov, Ritz 2009
March-Russell, West 2009
Shepherd, Tait, Zaharijas 2009
K.Petraki+, 2014+



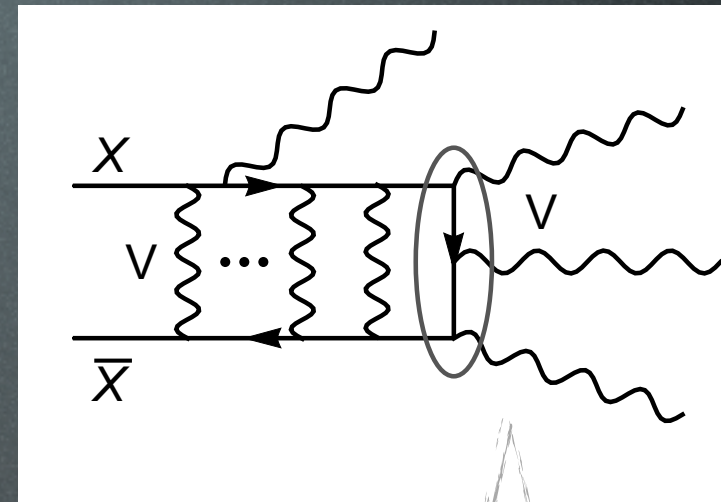
tree level annihilation



size of the XX system

If $\alpha M/m_V \gtrsim 1$, the force is **long range**:
Sommerfeld enhanced

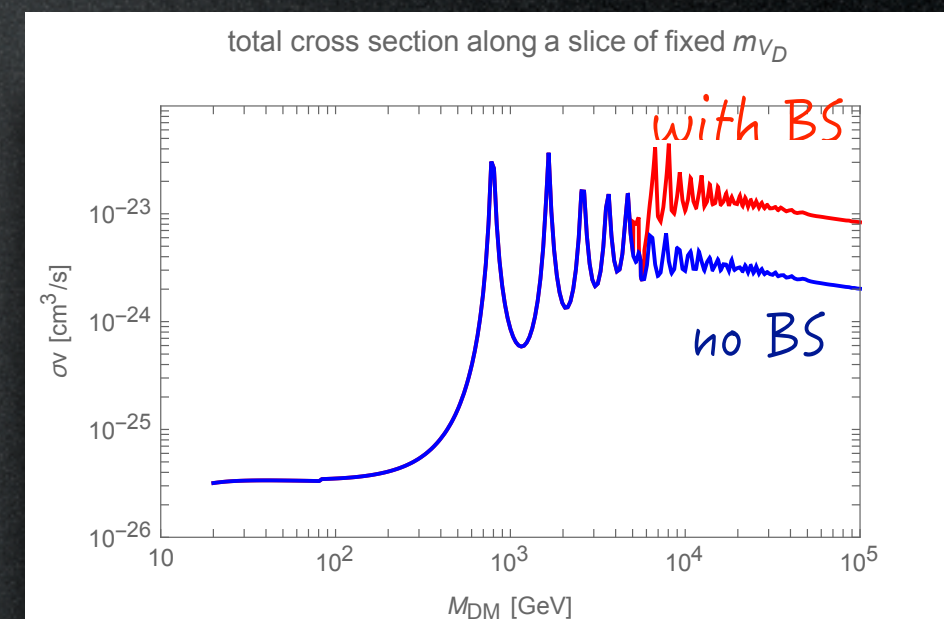
range



binding energy of the XX system

If $\alpha^2 M/2m_V \gtrsim 1$, **bound states** form

emitted mediator

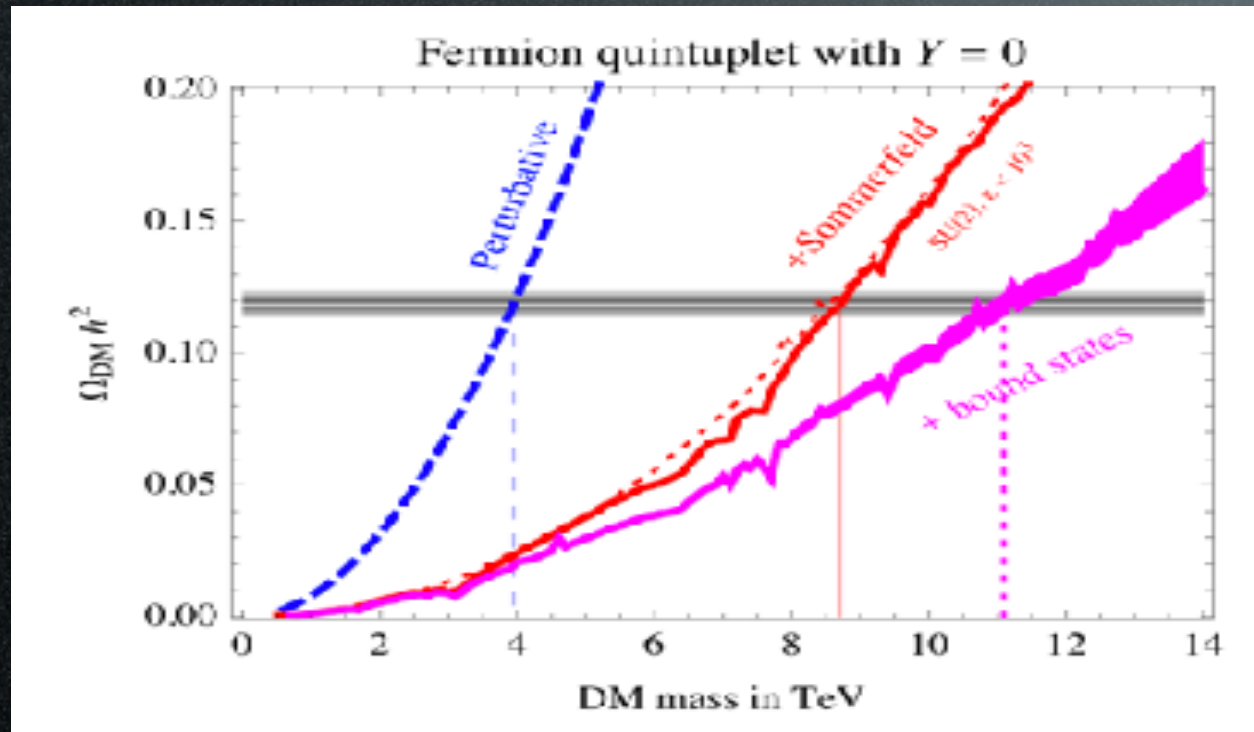


Post scriptum

Bound state formation

Mitridate, Redi, Smirnov, Strumia 1702.01141

impact on thermal mass and indirect detection

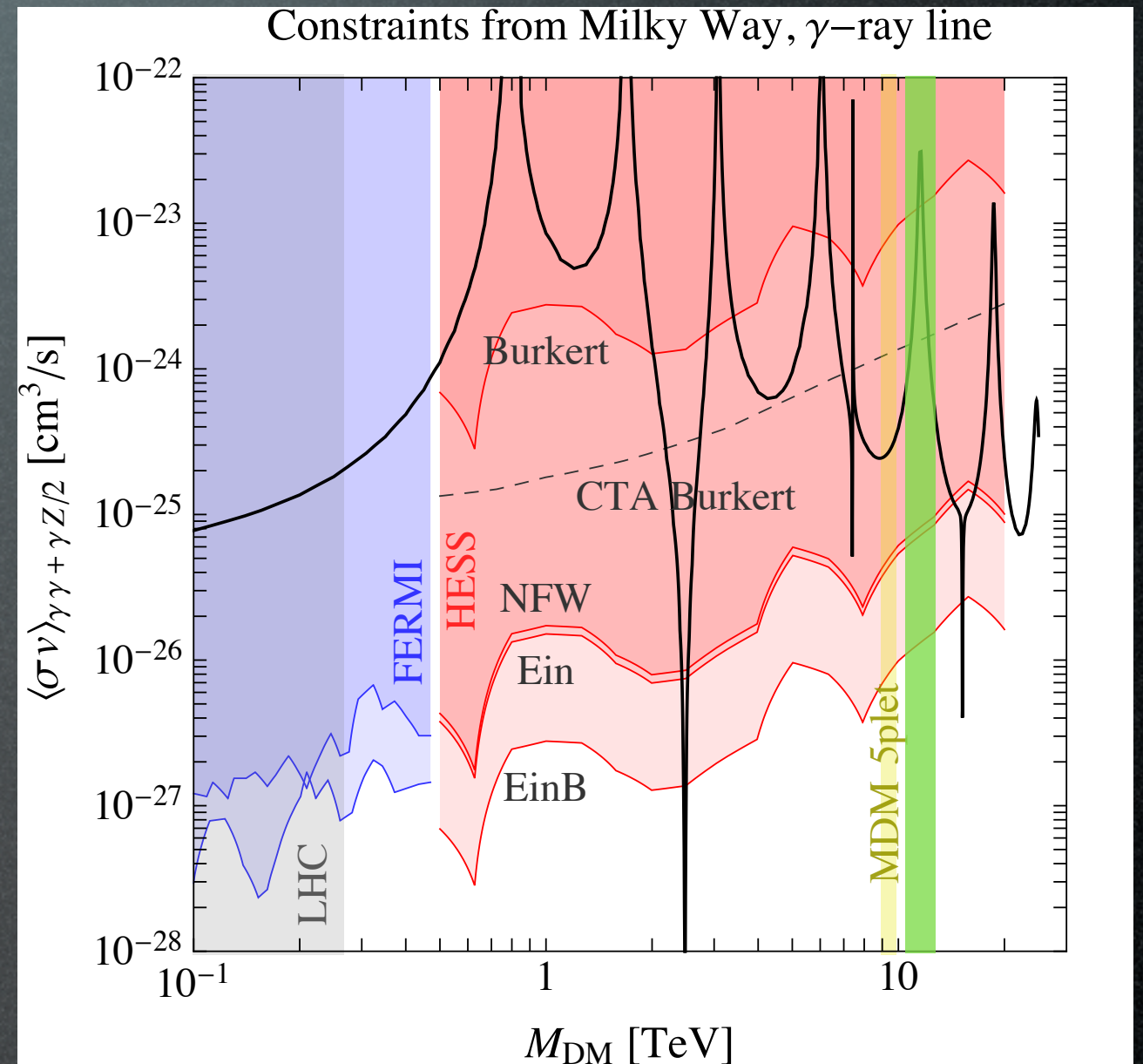
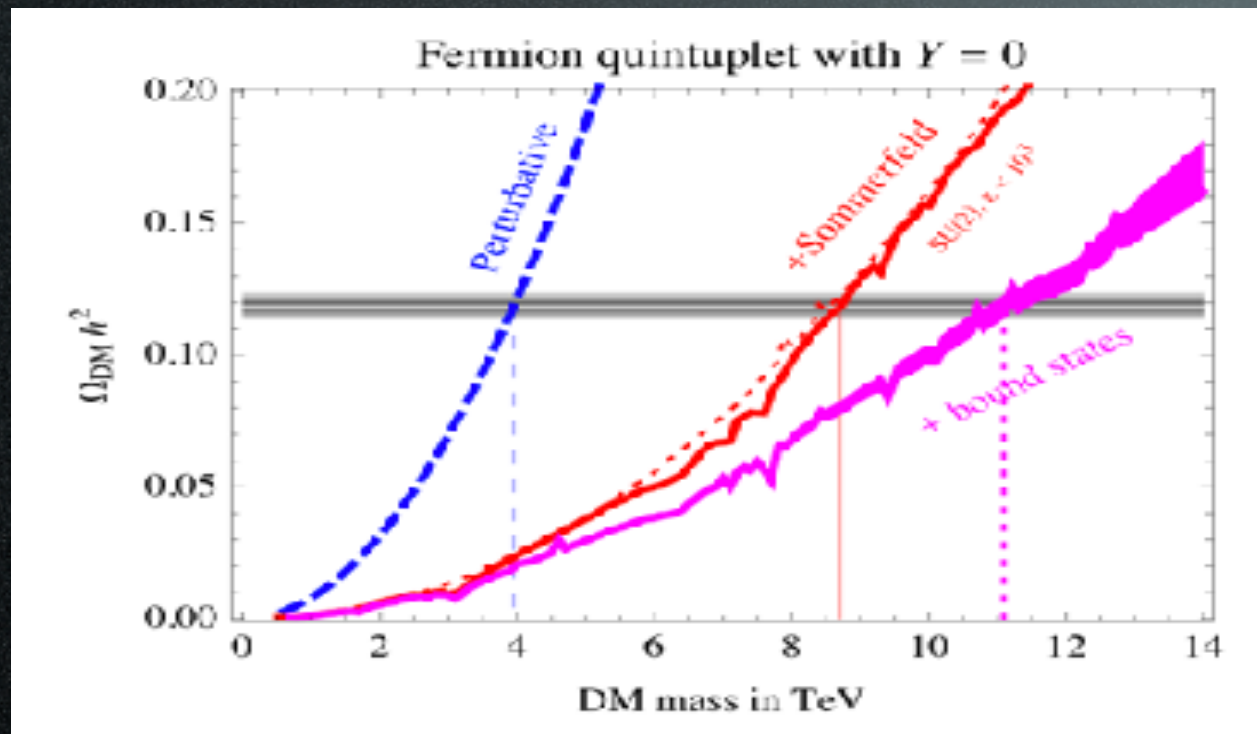


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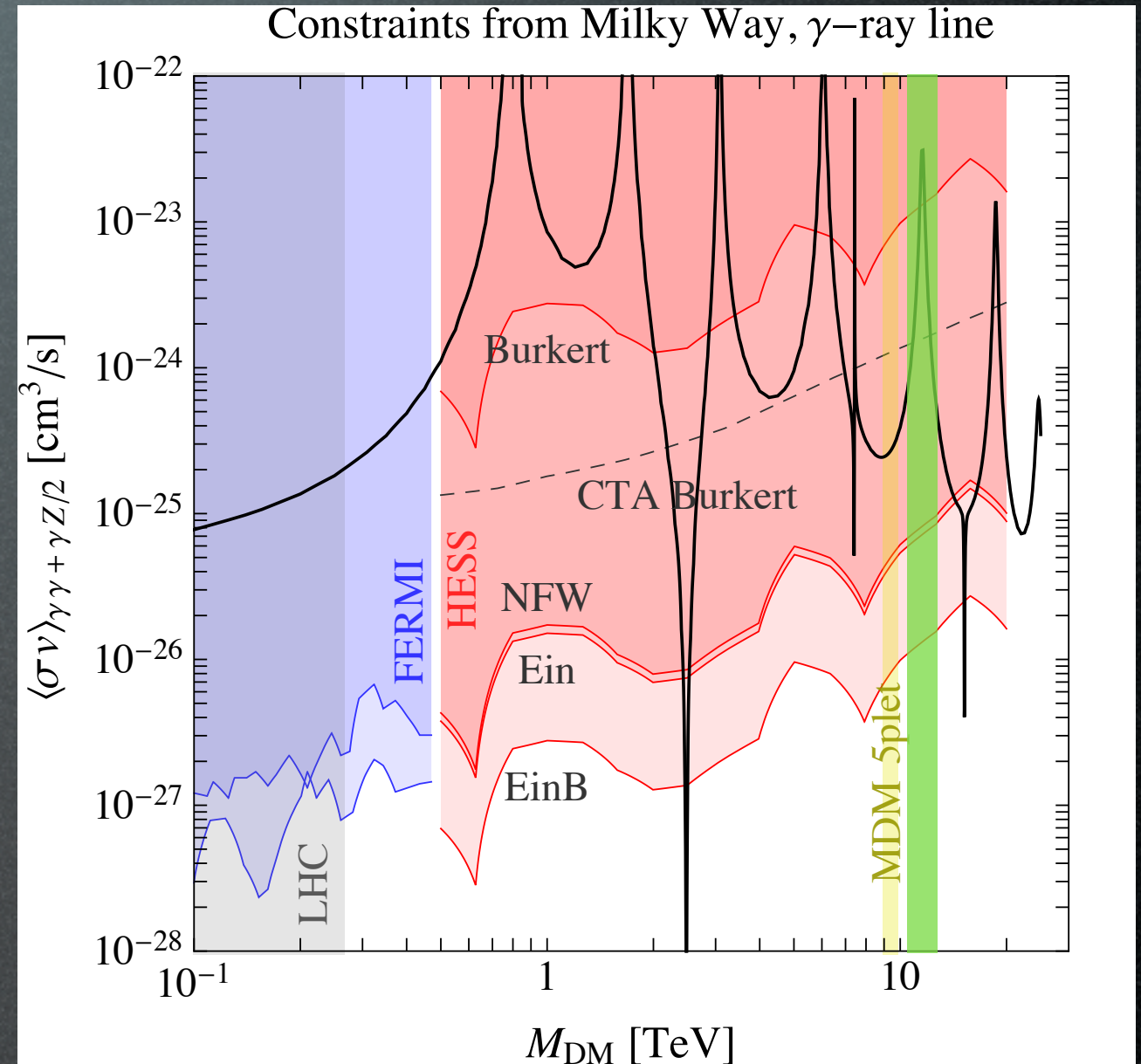
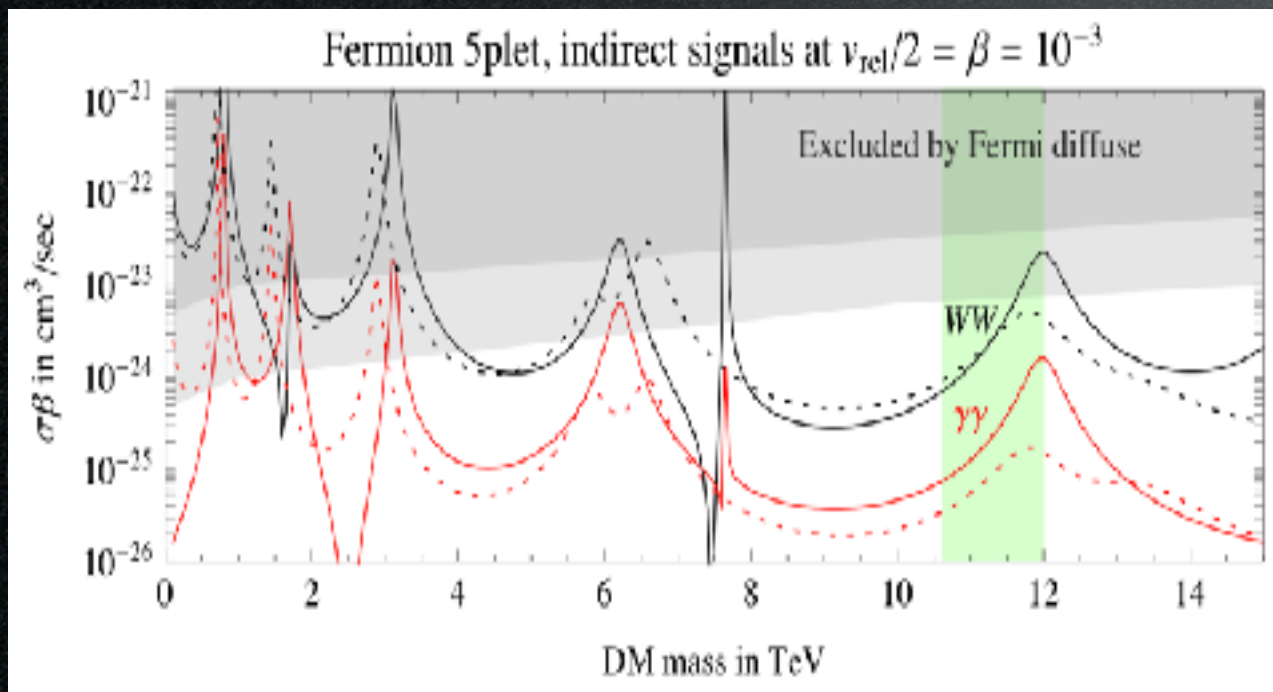
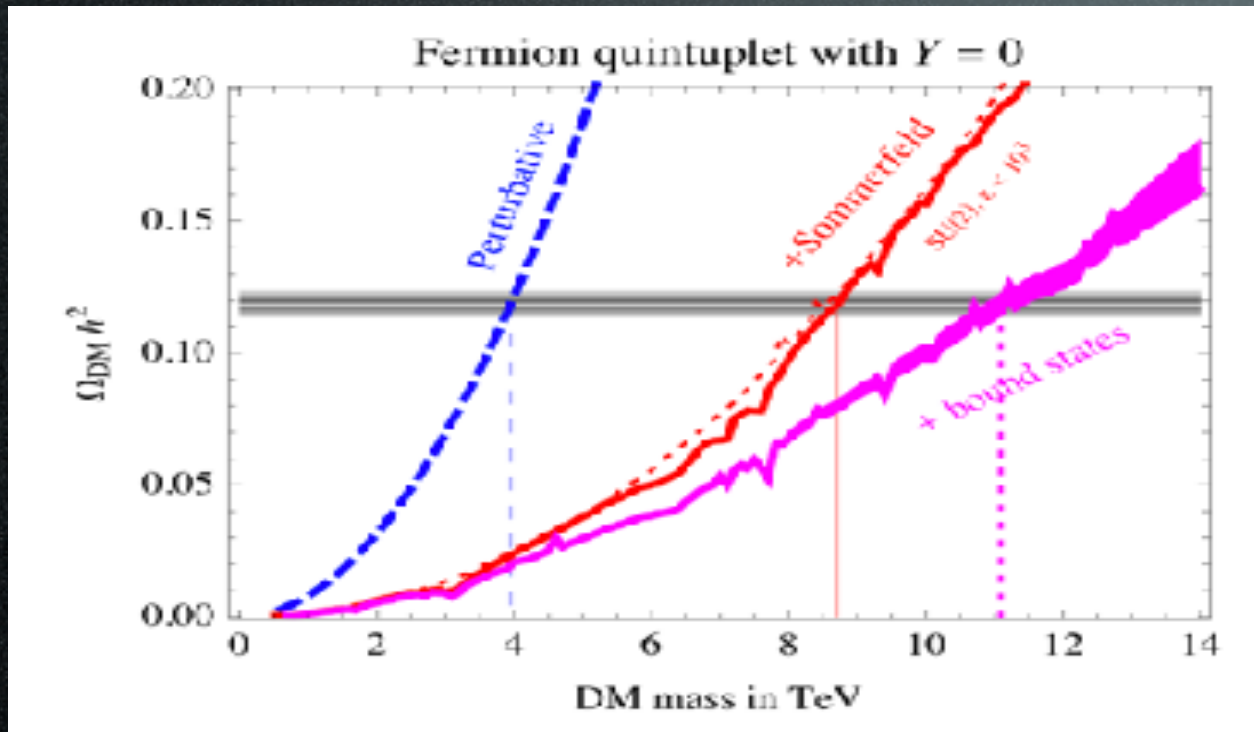


Post scriptum

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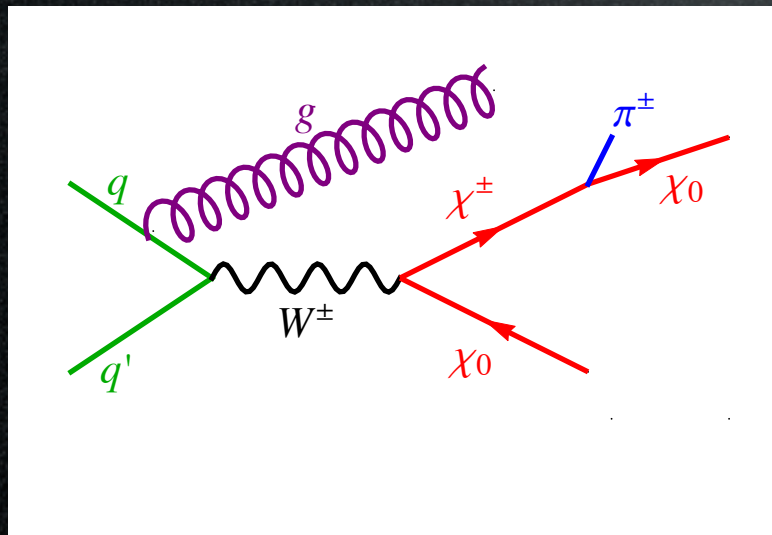
←
Mitridate, Redi,
Smirnov, Strumia
1702.01141v2

Collider searches

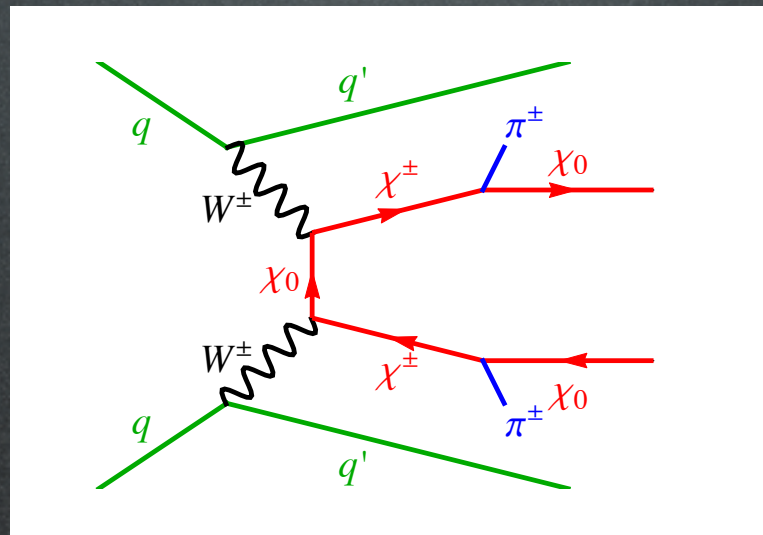
At 9.4 TeV, no hope at the LHC.

- relax the mass constraint
- consider next-to-minimal cases (e.g. the triplet = pure wino)
- explore reach of 100 TeV future collider

Mono-X

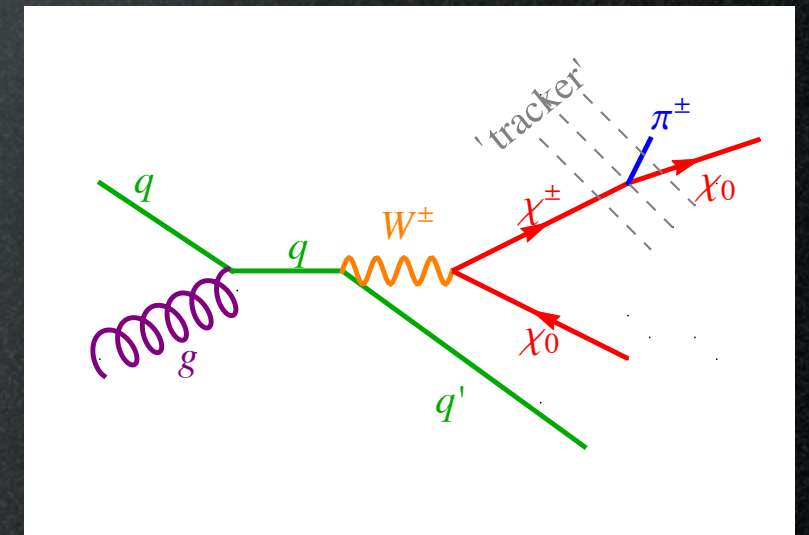


VBF



di-jets + MET

Disappearing tracks



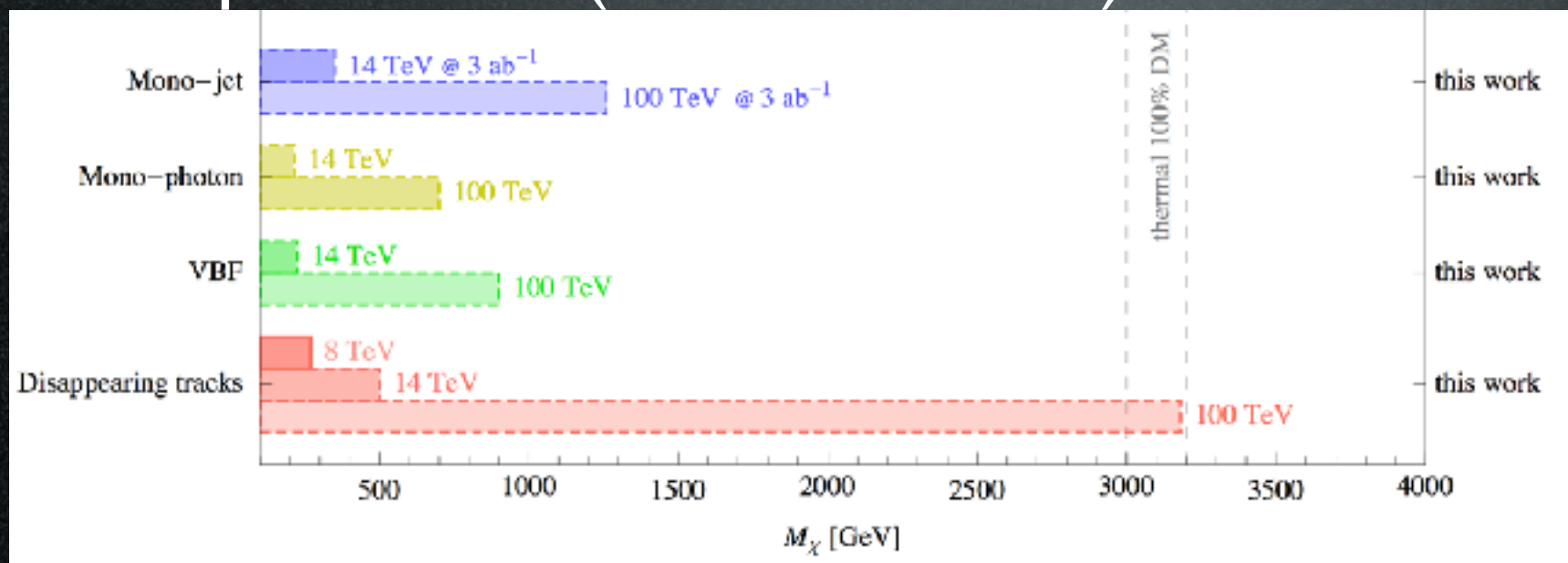
or: indirect searches

Collider searches

At 9.4 TeV, no hope at the LHC.

- relax the mass constraint
- consider next-to-minimal cases (e.g. the triplet = pure wino)
- explore reach of 100 TeV future collider

For triplet MDM (a.k.a. wino DM)



Cirelli, Sala, Taoso 1407.7058

For 5plet MDM

Model	$\sqrt{s} = 8 \text{ TeV}$				$\sqrt{s} = 14 \text{ TeV}$					
	ATLAS		CMS		Exclude			Discover		
	Expected	Observed	Expected	Observed	500%	100%	20%	500%	100%	20%
Wino	224	238	203	195	354	483	635	287	394	514
Majorana Fiveplet	256	267	234	226	410	524	668	340	448	576
Dirac Fiveplet	283	293	259	251	465	599	743	381	503	639

Ostdiek, 1506.03445

Collider searches

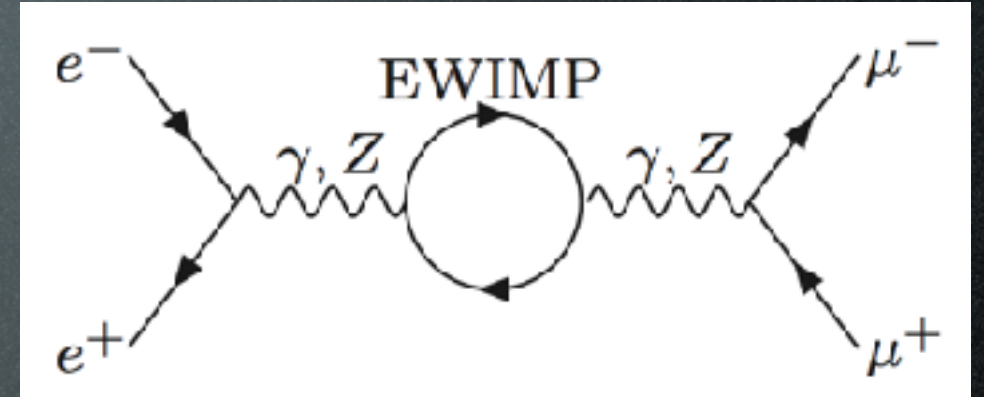
Indirect searches at a future Linear Collider:

Collider searches

Indirect searches at a future Linear Collider:

Harigaya, Ichikawa, Matsumoto..., 1504.03402

Even if $\sqrt{s} < M_{\text{DM}}$, one can see the effects
in precision measurements

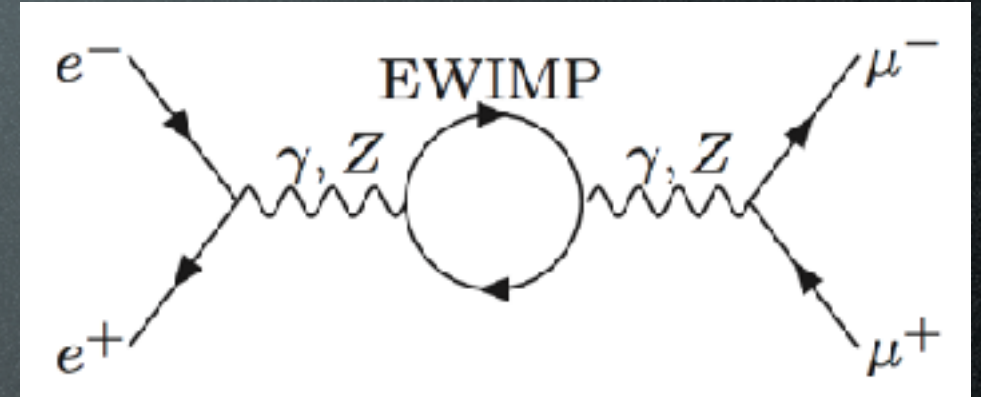


Collider searches

Indirect searches at a future Linear Collider:

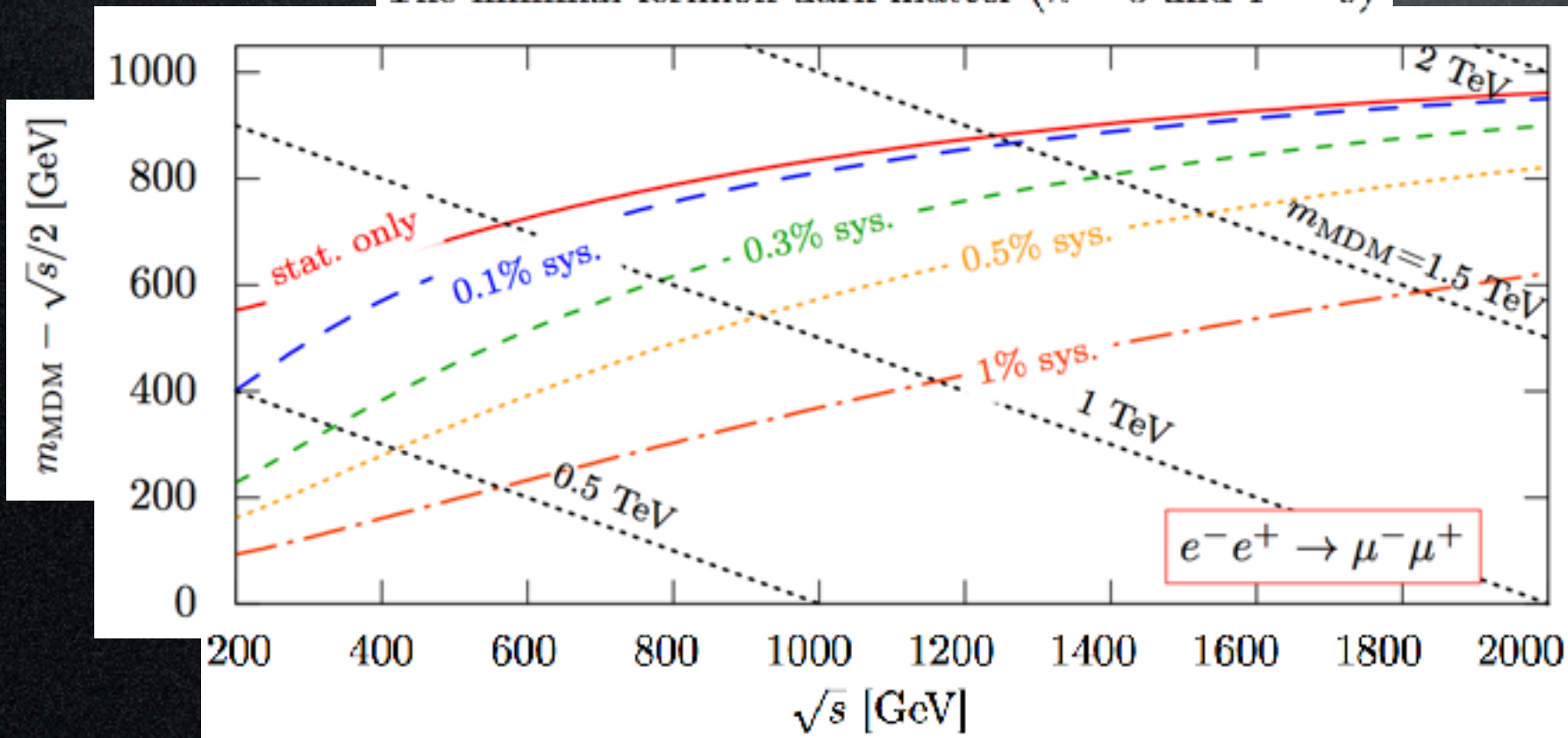
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Result:

The minimal fermion dark matter ($n = 5$ and $Y = 0$)

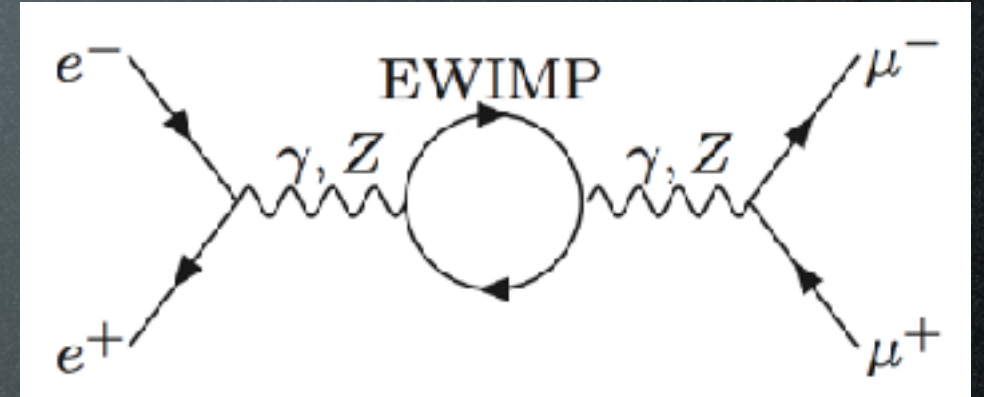


Collider searches

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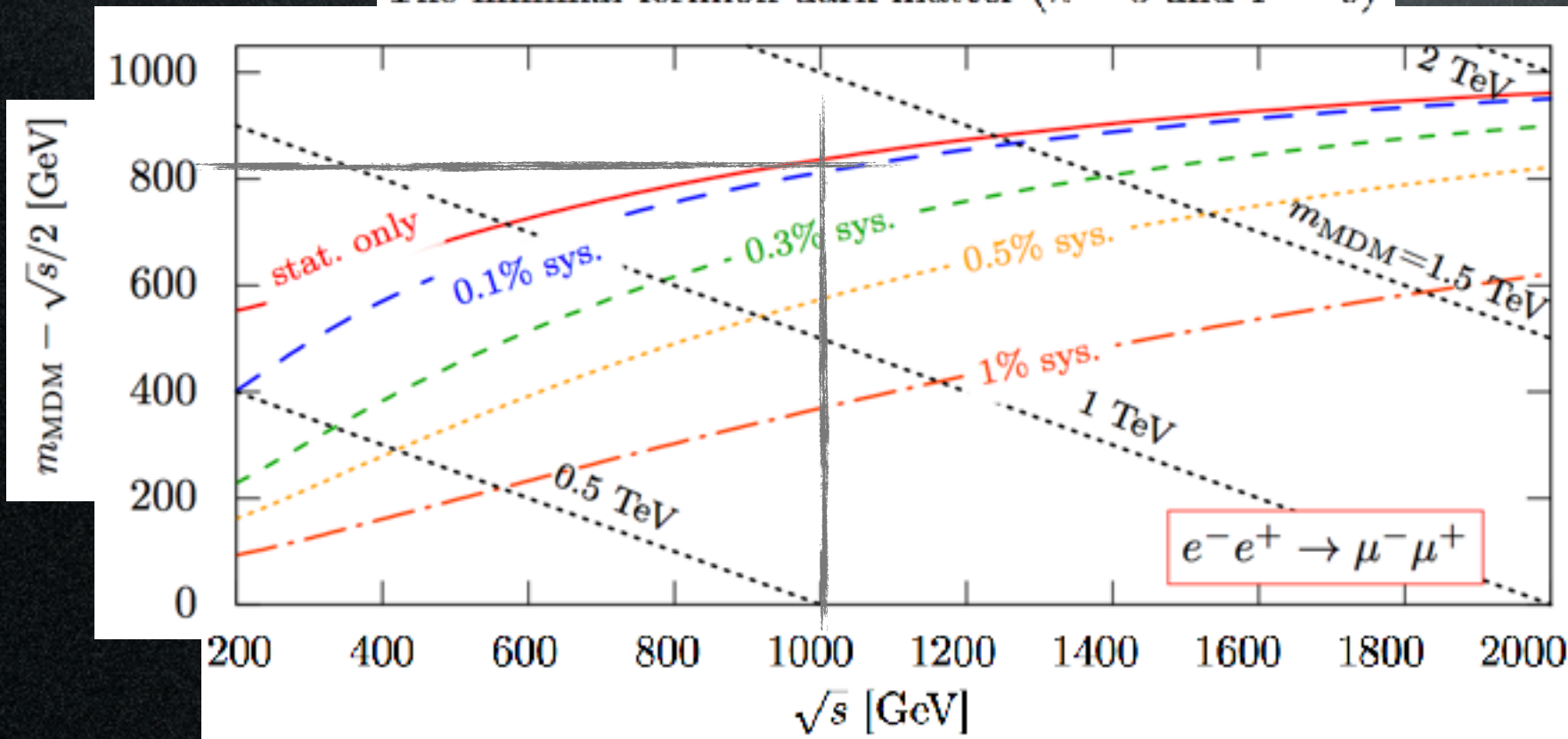
Harigaya, Ichikawa, Matsumoto..., 1504.03402

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Result:

The minimal fermion dark matter ($n = 5$ and $Y = 0$)



How to read the plot:
a LC with $\sqrt{s} = 1 \text{ TeV}$,
assuming only stat uncertainties,
will be sensitive to $m_{\text{DM}} - \sqrt{s}/2 \sim 800 \text{ GeV}$
i.e. $m_{\text{DM}} \sim 1.3 \text{ TeV}$ (indeed see the dotted
isocontours of the DM mass)

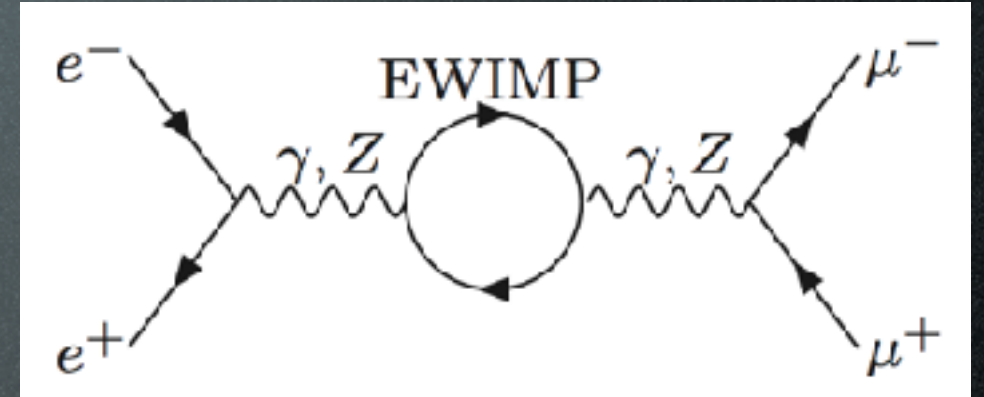
similar plots for other channels
& for other candidates

Collider searches

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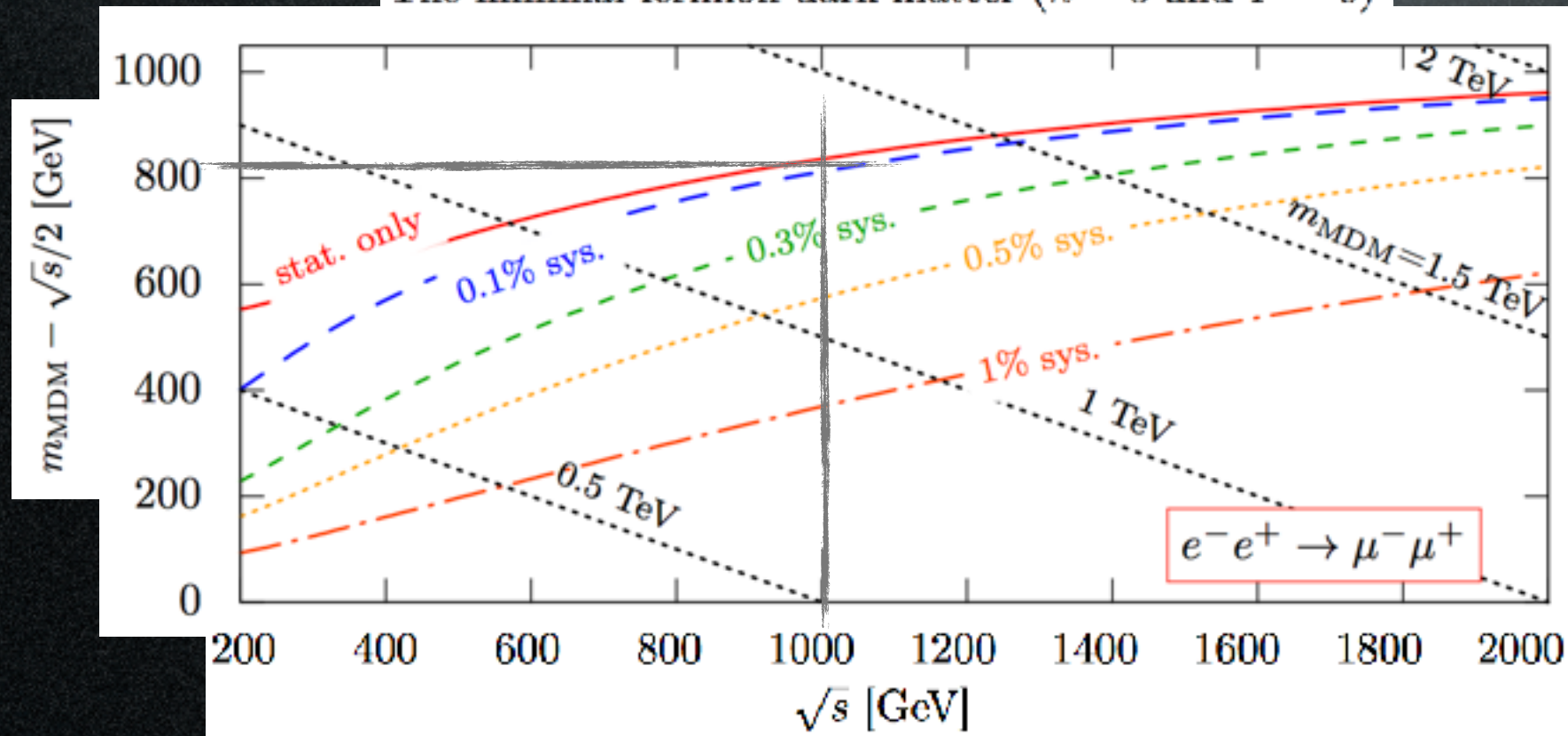
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isocontours of the DM mass)

similar plots for other channels
& for other candidates

one can go beyond the collider energy, but not by much

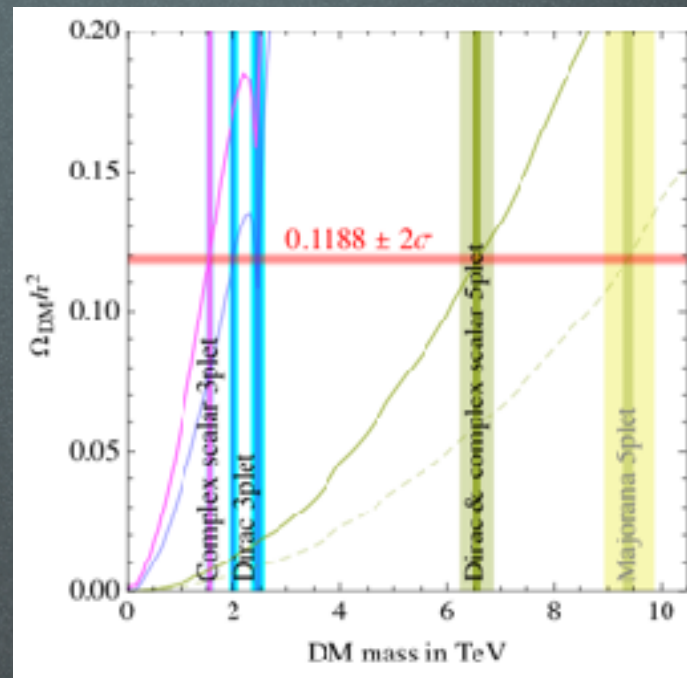
Bonus track

Some interesting recent extensions:

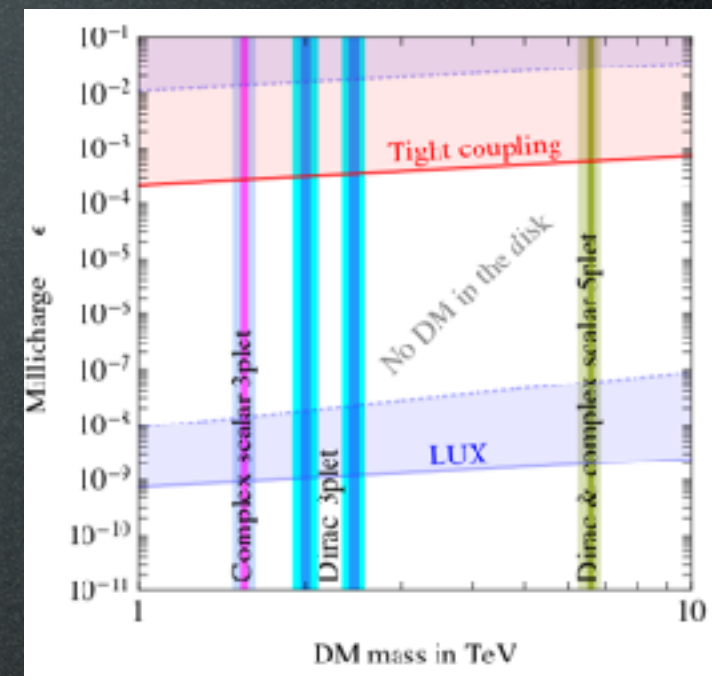
- millicharged MDM

Del Nobile, Nardecchia, Panci 1512.05353

assume $Y = \varepsilon \neq 0$,
-> implies stability
-> for suitable ε , no DD



relic abundance



constraints

- decaying MDM, if $\Lambda < M_{\text{Planck}}$

Del Nobile, Nardecchia, Panci 1512.05353

-> observable consequences in gamma rays

- 'natural' MDM

Fabbrichesi,
Urbano
1510.03861

MDM induces (at 2-loops) m_h corrections => small hierarchy prob

-> supersymmetrize it!:

- fermion/boson cancellations restore naturalness
- stability preserved by SuSy

Bonus track

Some interesting recent extensions:

- asymmetric MDM

Boucenna, Krauss, Nardi 1503.01119

- MDM and vacuum stability

Cai, Ramsey-Musolf et al., 1108.0969

Cai et al., 1508.04034

- non-thermally produced MDM

Aoki, Toma, Vicente 1507.01591

- incorporating neutrino masses

Cai, Schmidt 1603.00255

Ahriche, McDonald, Nasri, Picek 1603.01247

Conclusions 1/2

The DM problem requires physics beyond the SM.

(*) unless it's Primordial Black Holes
generated by higgs perturbations!
Espinosa, Racco, Riotto 1710.11196

Conclusions 1/2

The DM problem requires physics beyond the SM.

WIMPs:

Conclusions 1/2

The DM problem requires physics beyond the SM.

WIMPs:

1. even without a larger framework, they are still appealing
2. the frontier is multi-TeV
3. searches are complementary and still have ground to cover

Conclusions 2/2

The DM problem requires physics beyond the SM.

Focussing on pure WIMPs,
we find one fully successful DM candidate:
massive, neutral, *automatically* stable.

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mass = 9.4 TeV
(or 11 TeV with bound states)

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The phenomenology is precisely computable:

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fermionic $SU(2)_L$ quintuplet with $Y = 0$,
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(or 11 TeV with bound states)

The phenomenology is precisely computable:

- tested by indirect detection (γ ray) exp's:
 - excluded if cuspy
 - not probed if cored

**Back-up
slides**

Production at colliders

OK, given what we (do not) see, how do we interpret?

effectively

simplified models

a simple model

a full theory

Production at colliders

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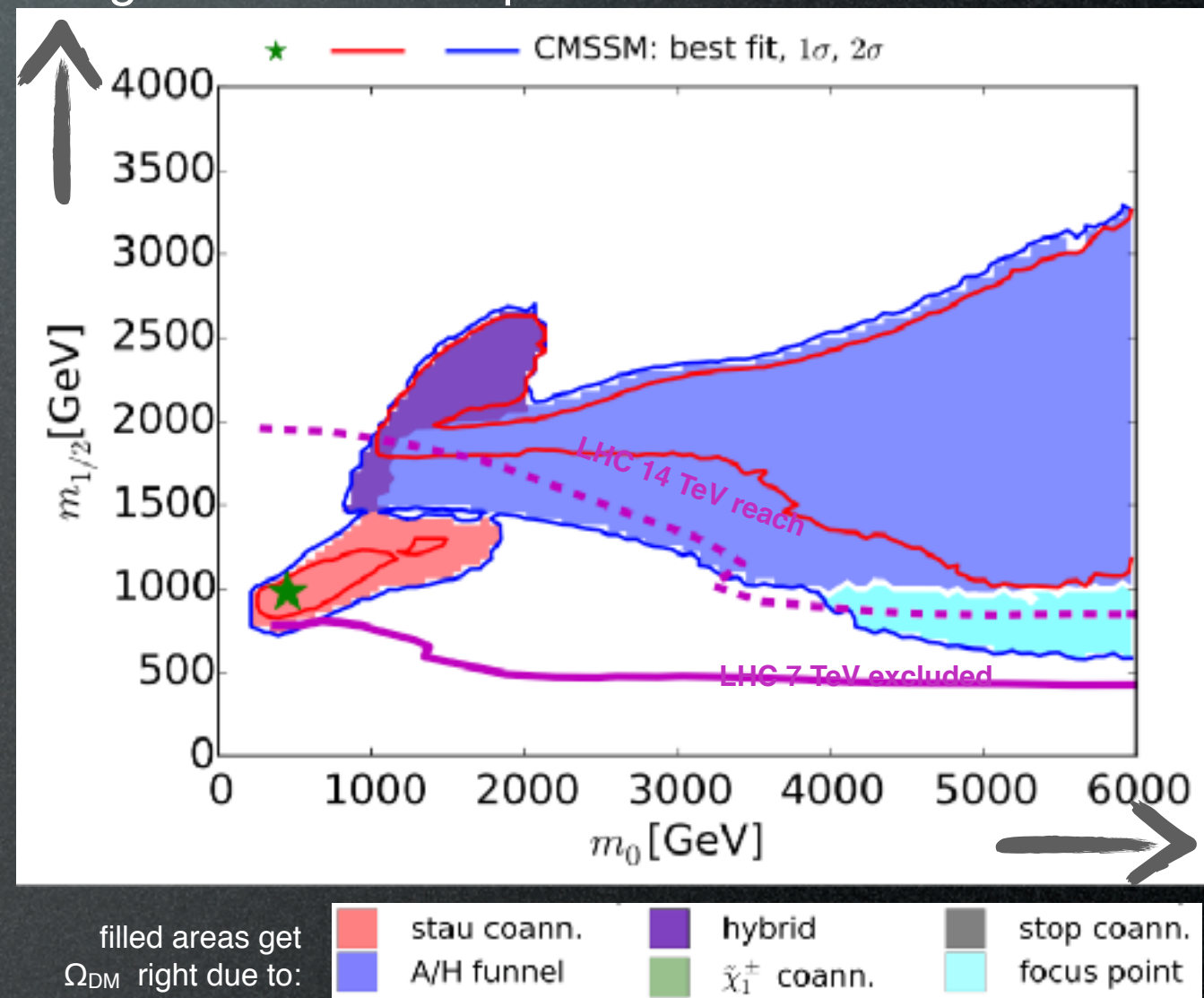
a simple model

a full theory

E.g. SUSY: cMSSM, pMSSM...

parameters: tens (or even hundreds!)

the glorious $m_0/m_{1/2}$ plane:



Production at colliders

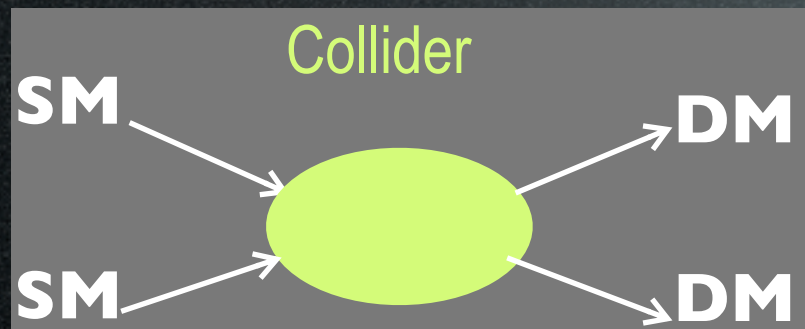
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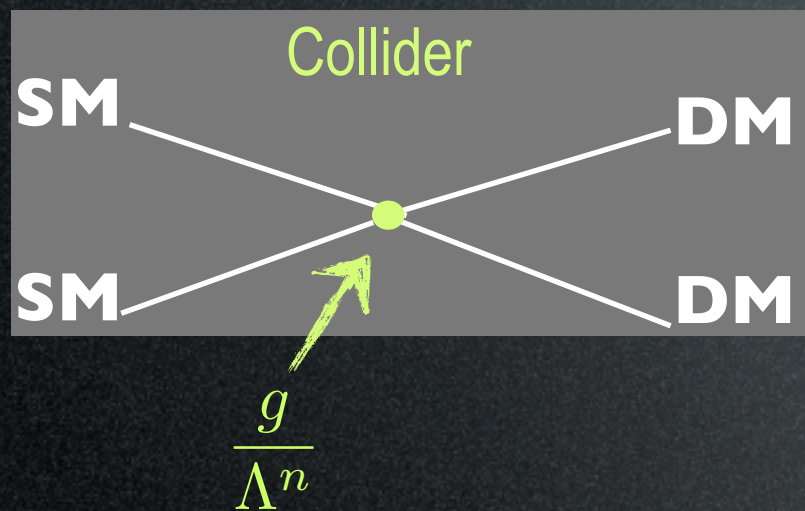
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$$\frac{1}{\Lambda_1^2} [q\bar{q}] [\chi\bar{\chi}]$$

$$\frac{1}{\Lambda_2^2} [q\gamma_\mu\bar{q}] [\chi\gamma^\mu\bar{\chi}]$$

parameters: Λ (assuming $g \simeq 1$)

Production at colliders

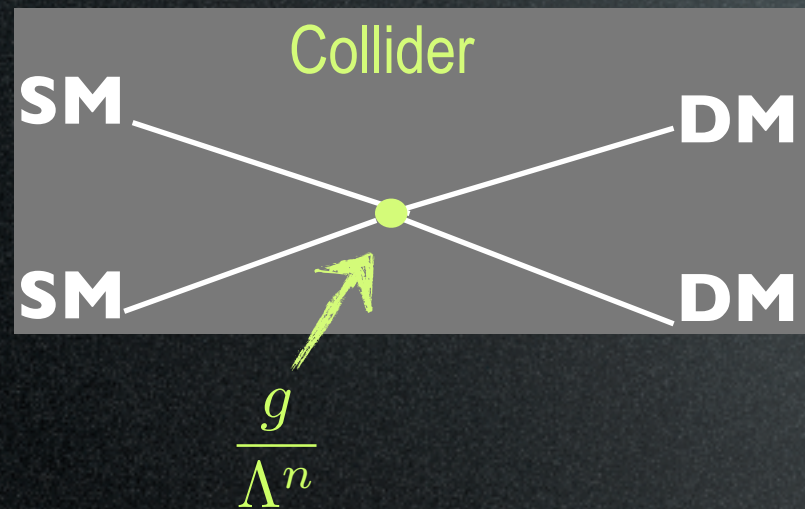
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(a) Operators for Dirac fermion DM

Name	Operator	Dimension	SI/SD
D1	$\frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} q$	7	SI
D2	$\frac{im_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} q$	7	N/A
D3	$\frac{im_q}{\Lambda^3} \bar{\chi} \chi \bar{q} \gamma^5 q$	7	N/A
D4	$\frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} \gamma^5 q$	7	N/A
D5	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$	6	SI
D6	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu q$	6	N/A
D7	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu \gamma^5 q$	6	N/A
D8	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$	6	SD
D9	$\frac{1}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$	6	SD
D10	$\frac{i}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \gamma^5 \chi \bar{q} \sigma_{\mu\nu} q$	6	N/A
D11	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \chi G^{\mu\nu} G_{\mu\nu}$	7	SI
D12	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \gamma^5 \chi G^{\mu\nu} G_{\mu\nu}$	7	N/A
D13	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \chi G^{\mu\nu} \tilde{G}_{\mu\nu}$	7	N/A
D14	$\frac{\alpha_s}{\Lambda^3} \bar{\chi} \gamma^5 \chi G^{\mu\nu} \tilde{G}_{\mu\nu}$	7	N/A

(b) Operators for Complex scalar DM

Name	Operator	Dimension	SI/SD
C1	$\frac{m_q}{\Lambda^2} \phi^\dagger \phi \bar{q} q$	6	SI
C2	$\frac{m_q}{\Lambda^2} \phi^\dagger \phi \bar{q} \gamma^5 q$	6	N/A
C3	$\frac{1}{\Lambda^2} \phi^\dagger \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu q$	6	SI
C4	$\frac{1}{\Lambda^2} \phi^\dagger \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu \gamma^5 q$	6	N/A
C5	$\frac{\alpha_s}{\Lambda^3} \phi^\dagger \phi G^{\mu\nu} G_{\mu\nu}$	6	SI
C6	$\frac{\alpha_s}{\Lambda^3} \phi^\dagger \phi G^{\mu\nu} \tilde{G}_{\mu\nu}$	6	N/A

Tim Tait, 2010+

and many many many others

Production at colliders

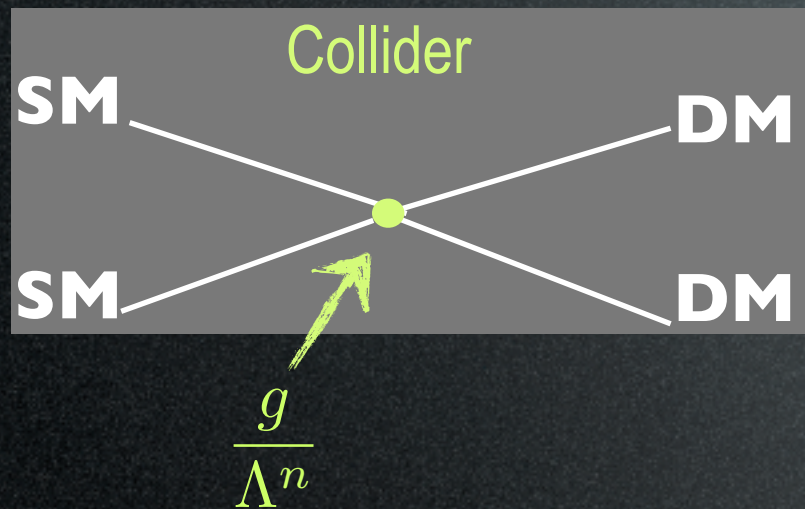
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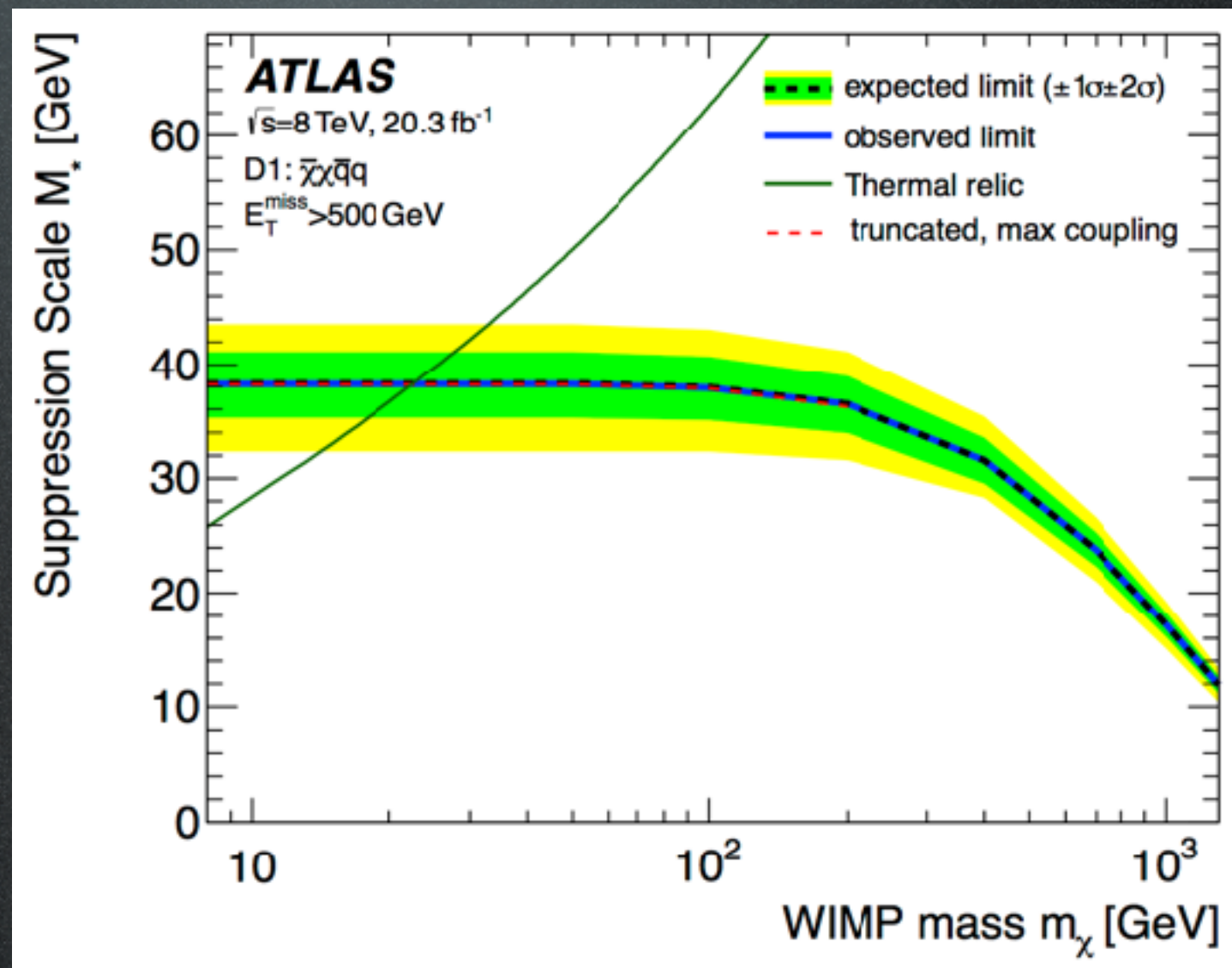
simplified models

a simple model

a full theory



limits on Λ



Production at colliders

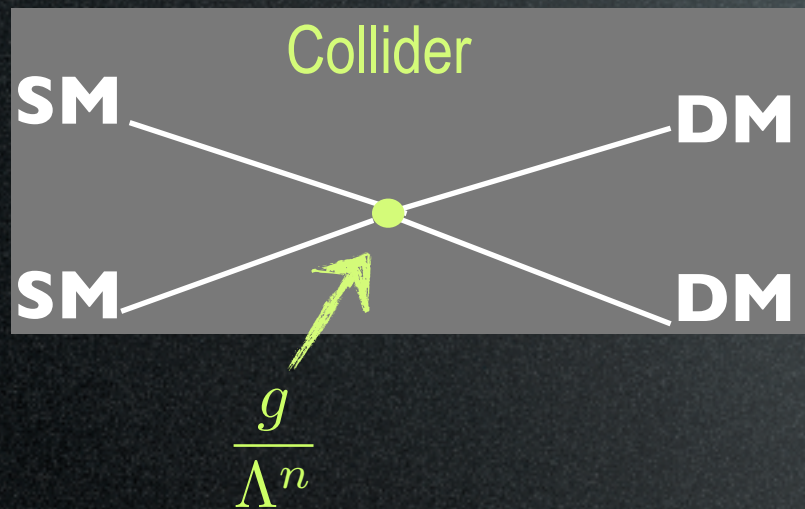
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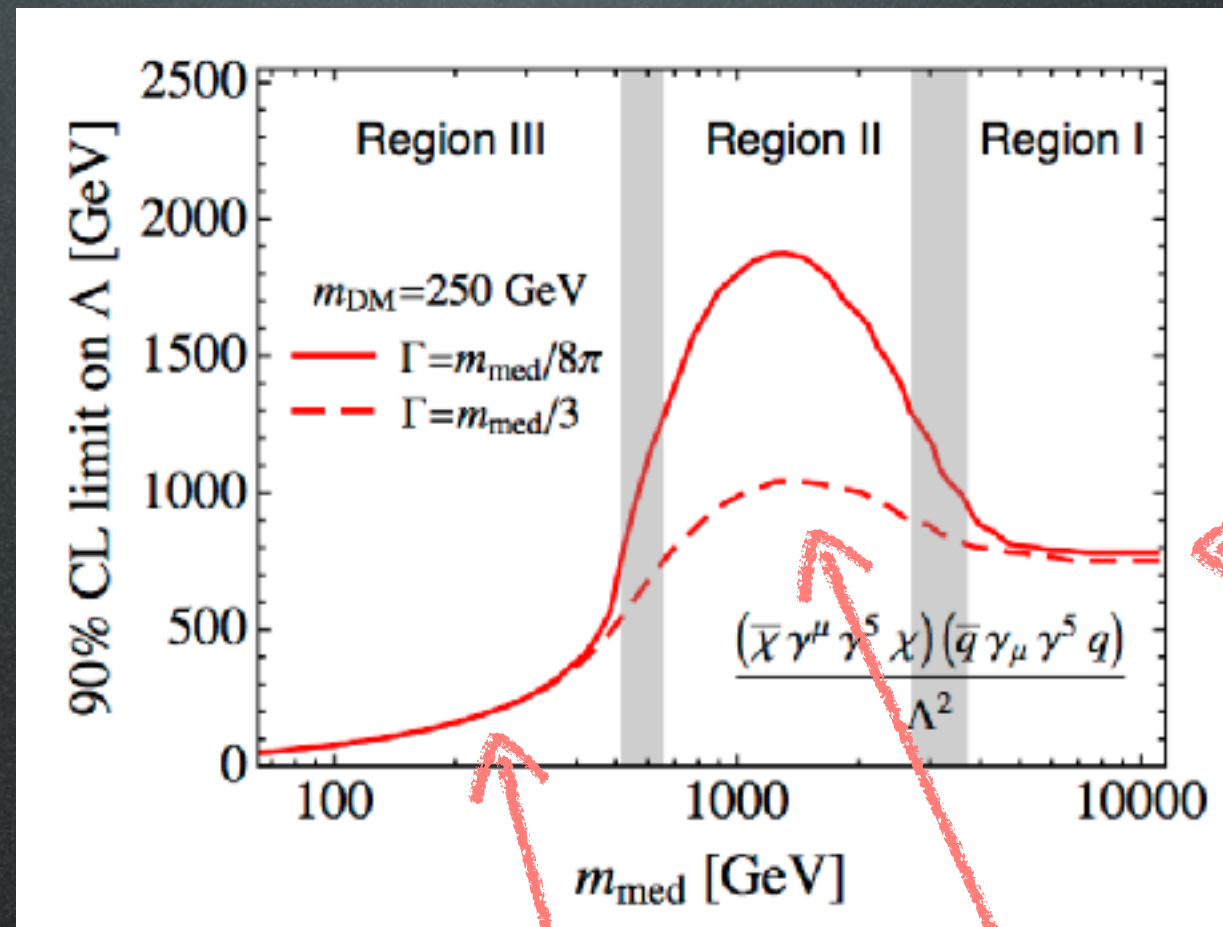
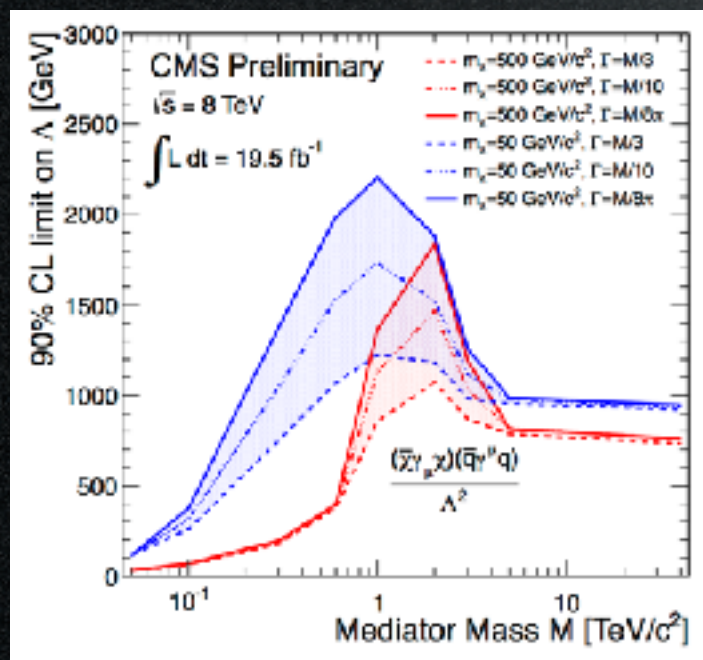
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but how 'sensible' are they?



EFT puts bounds on Λ

but Λ is actually $\Lambda = \frac{m_{\text{med}}}{\sqrt{g_q g_\chi}}$
at large m_{med} , the two agree

resolving, at small m_{med} ,
what is the discrepancy?
i.e. what is the limit deduced on Λ ,
defined as above?

limits agree

Buchmuller, Dolan,
McCabe 1308.6799

EFT bound is
much larger

EFT underestimates bound
by some % up to factor 2 or more

Production at colliders

OK, given what we (do not) see, how do we interpret?

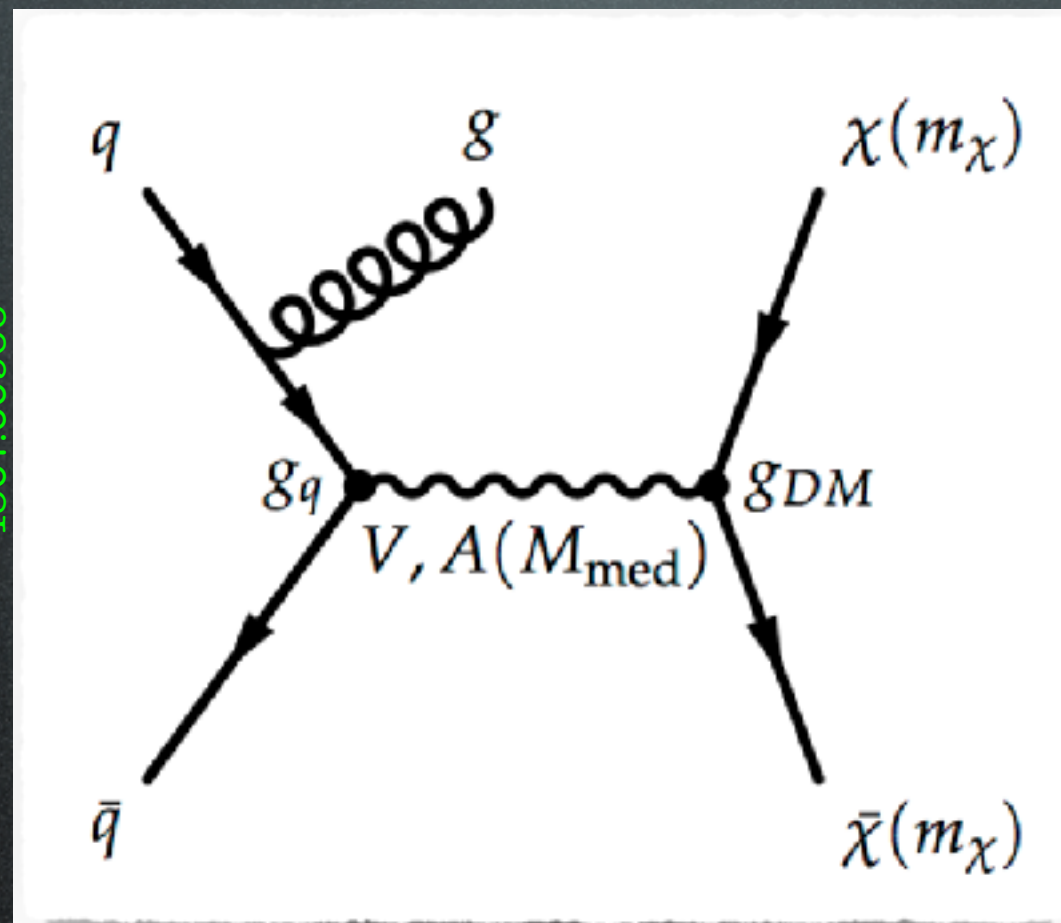
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ATLAS+CMS DM forum,
1507.00966



parameters:

$m_\chi, m_{\text{med}}, g_q, g_{\text{DM}}$

$$\mathcal{L}_{\text{vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

$$\mathcal{L}_{\text{axial-vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma^5 q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi.$$

(and similarly for t-channel, scalar mediator, scalar DM etc...)

Production at colliders

OK, given what we (do not) see, how do we interpret?

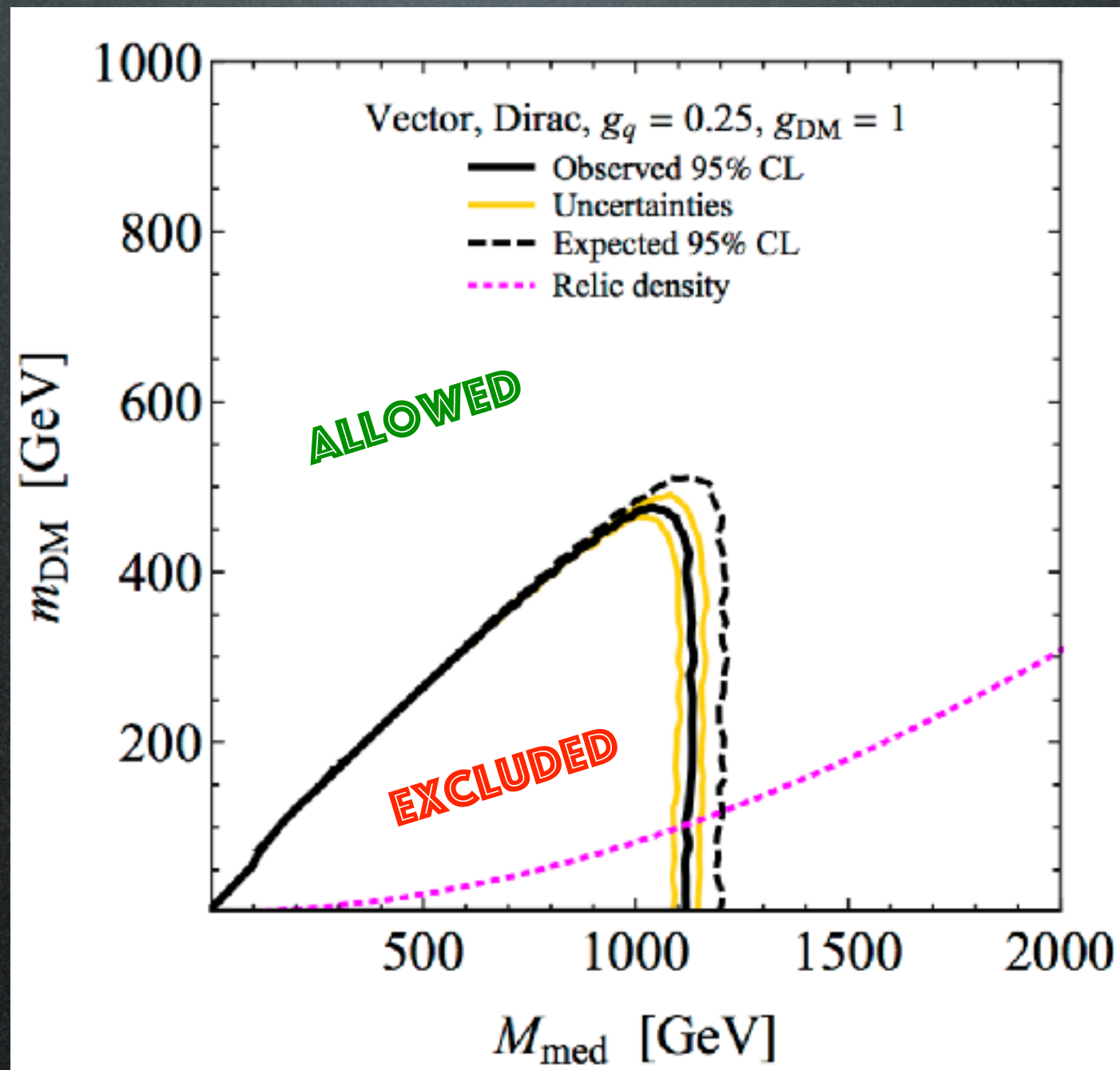
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mock plot:



Production at colliders

OK, given what we (do not) see, how do we interpret?

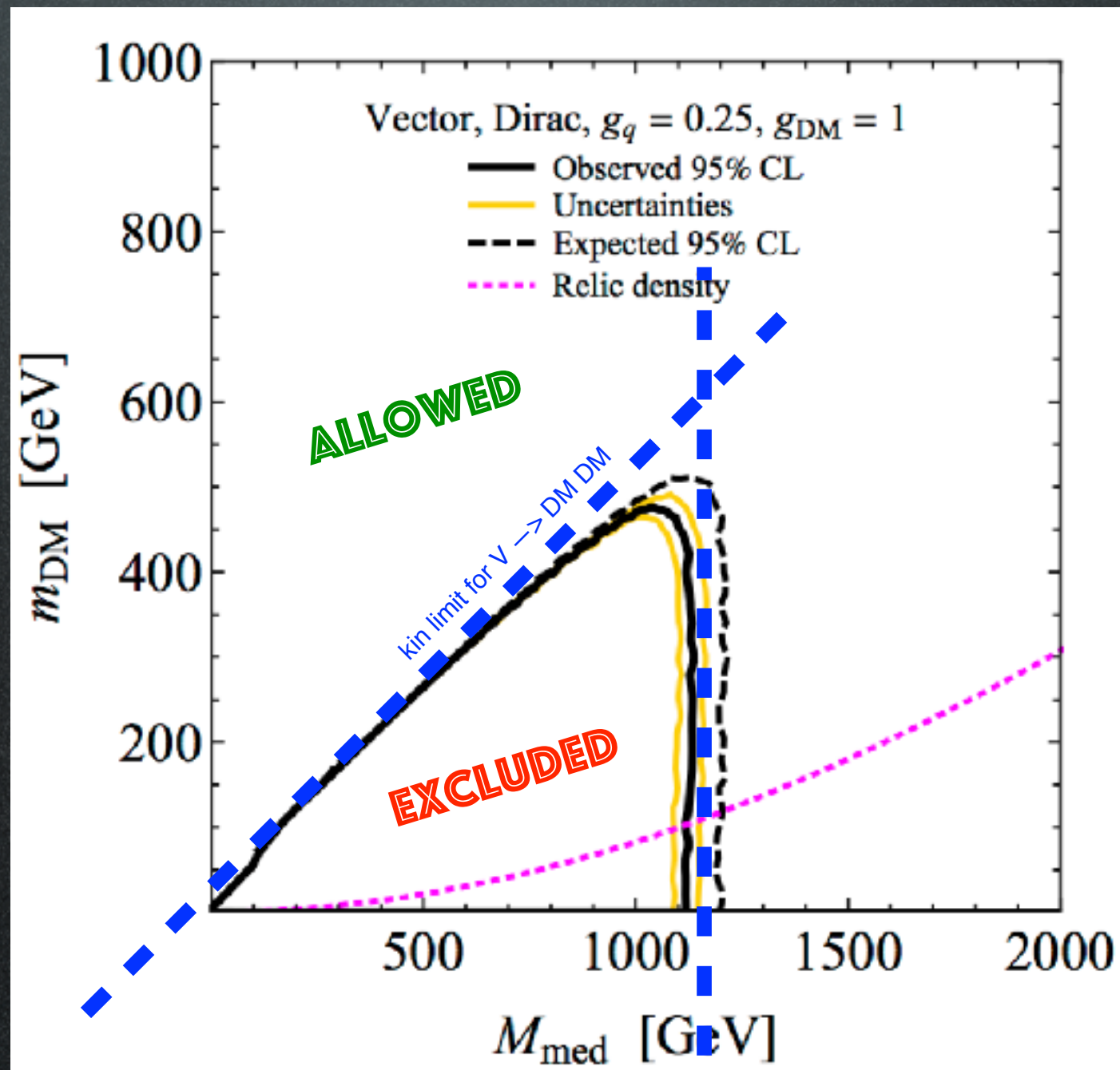
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Production at colliders

OK, given what we (do not) see, how do we interpret?

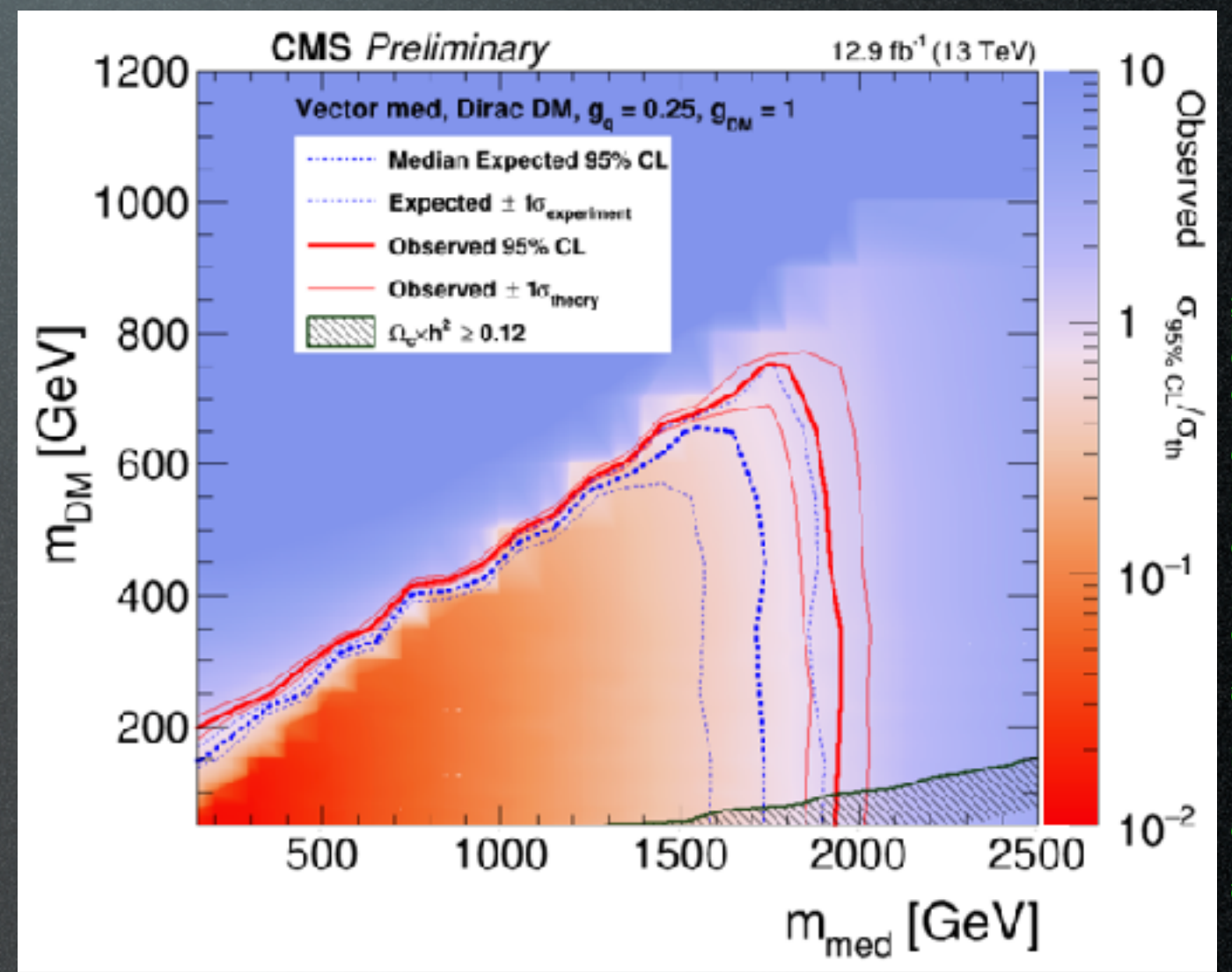
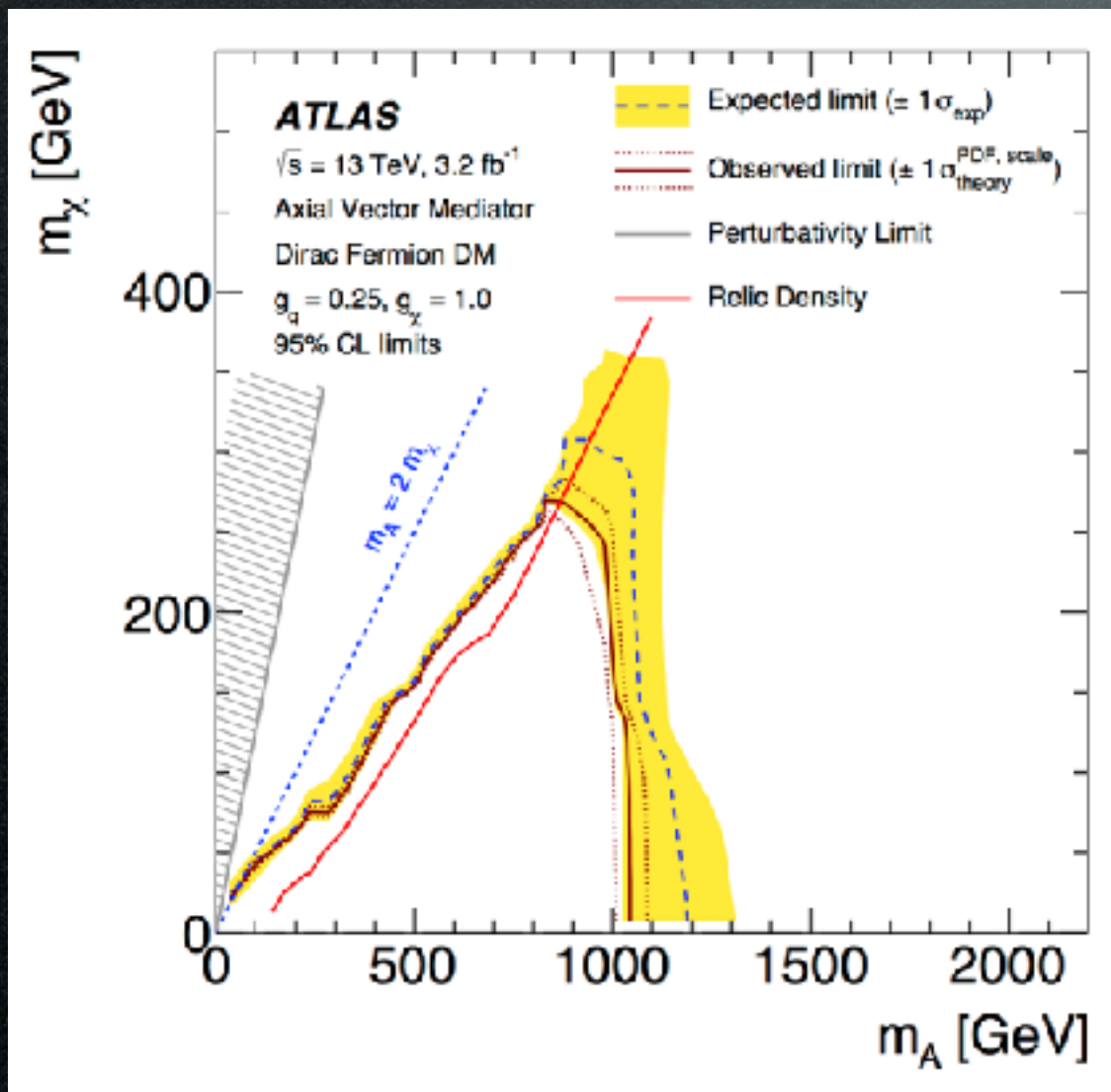
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actual plots:



Production at colliders

OK, given what we (do not) see, how do we interpret?

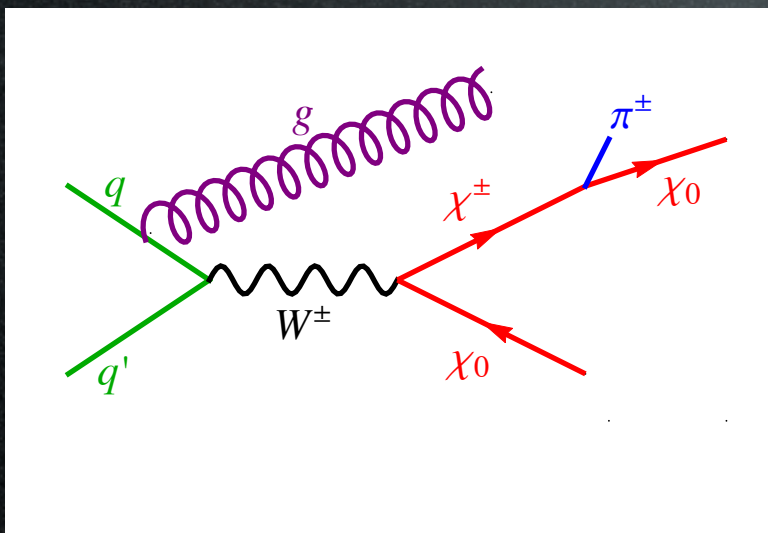
effectively

simplified models

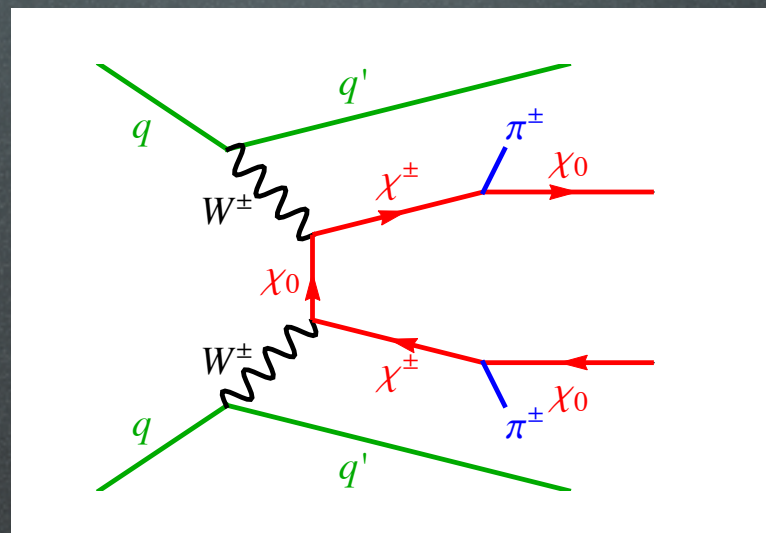
a simple model

a full theory

e.g. pure WIMP model (a.k.a. 'Minimal DM')



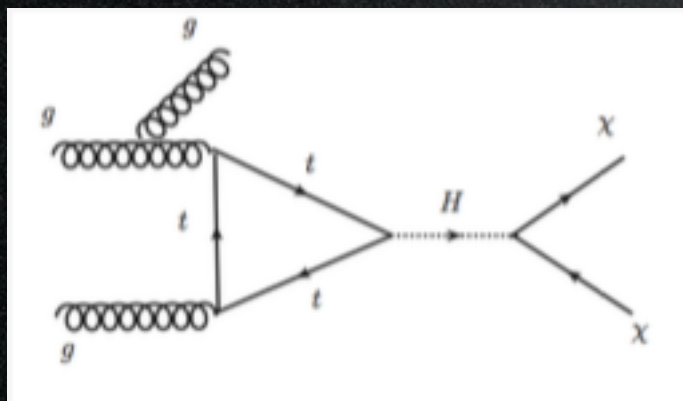
mono-jet



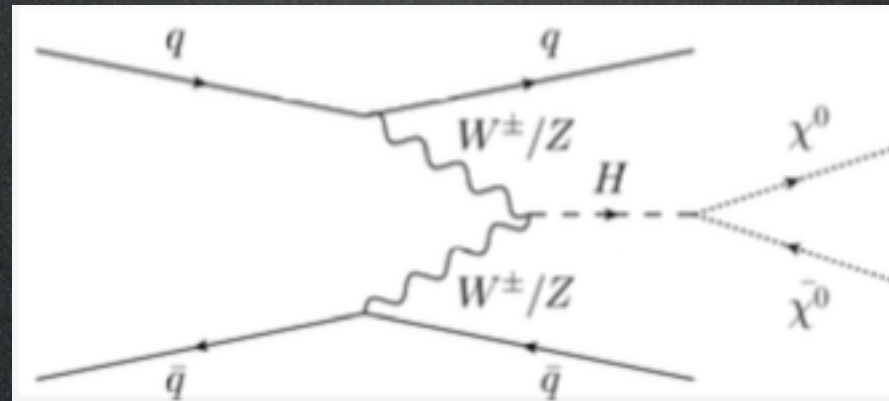
Vector Boson Fusion (VBF)

parameters: M_{DM}

e.g. higgs portal models



mono-jet



Vector Boson Fusion (VBF)

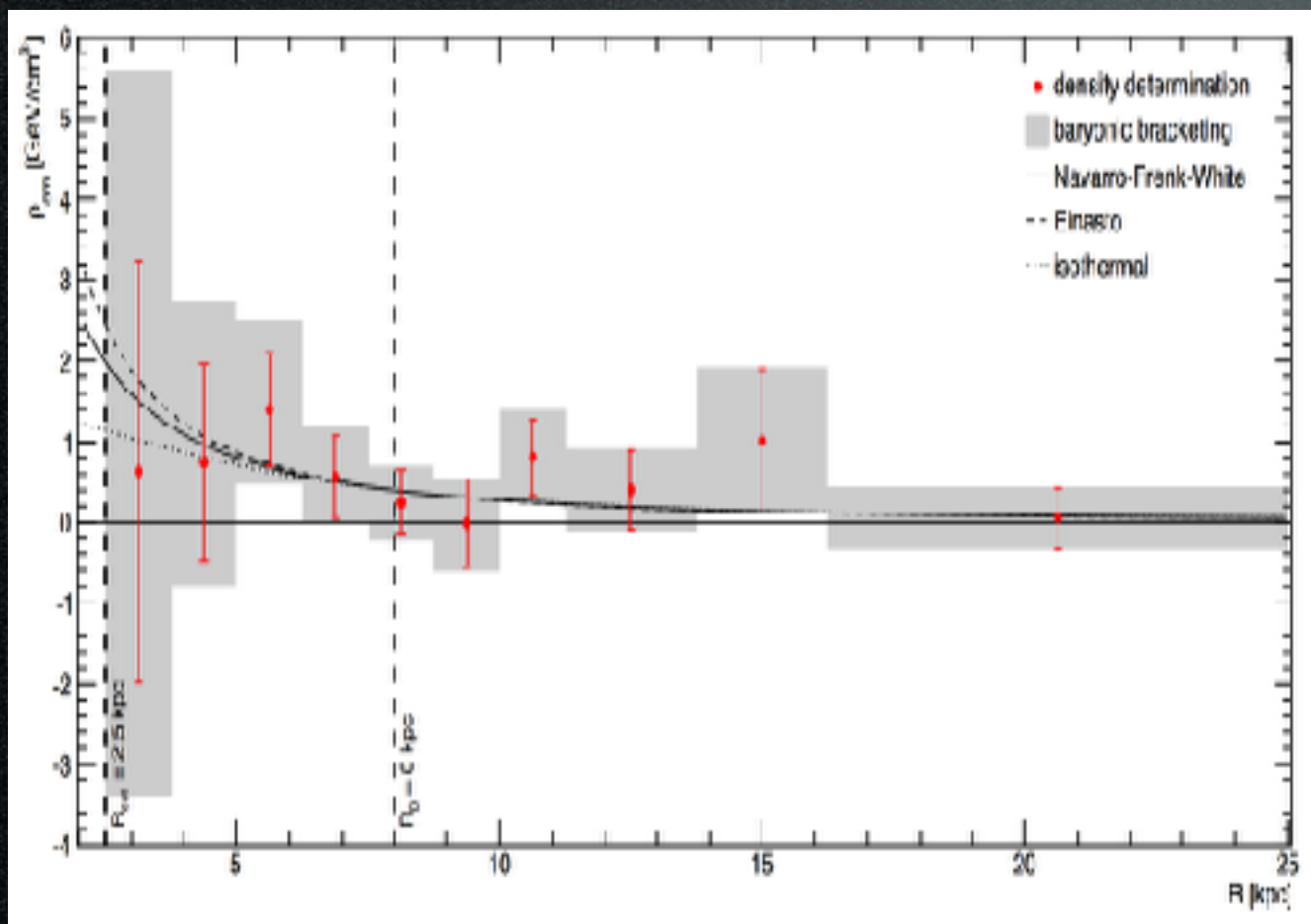
parameters: $g_{h\chi}$

3. Indirect Detection

So is the **Milky Way** profile **peaked** or **cored**?

Observations:

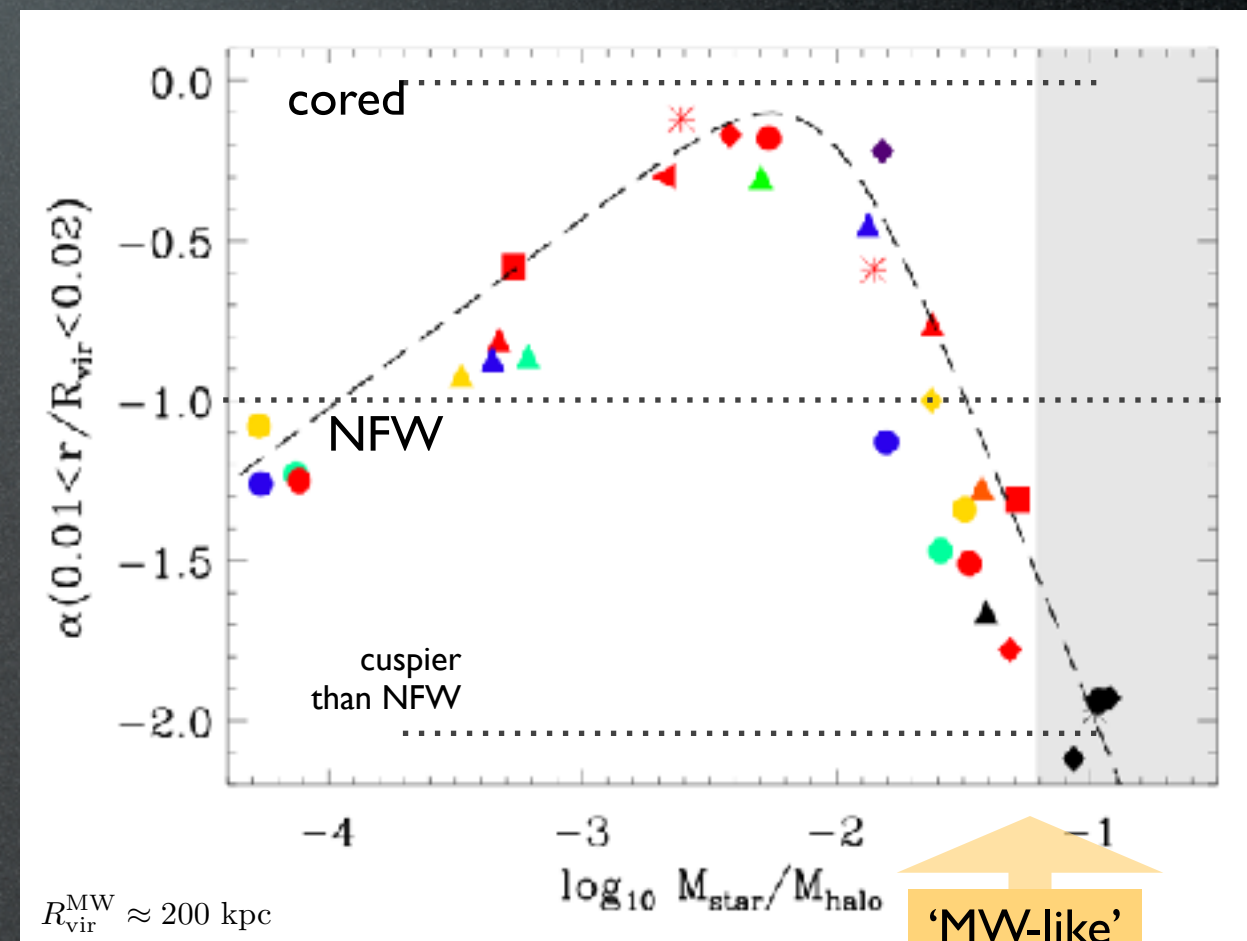
- (difficult from inside)
- no discriminating power
- anyway, no resolution < 2.5 kpc



Pato, Iocco 1504.03317

Simulations:

- (still open debate, but)
- prefer cuspy
- but: no resolution < 2 kpc



Di Cintio ... Macciò1306.0898

3. Indirect Detection

MW center area, search for γ -ray lines:

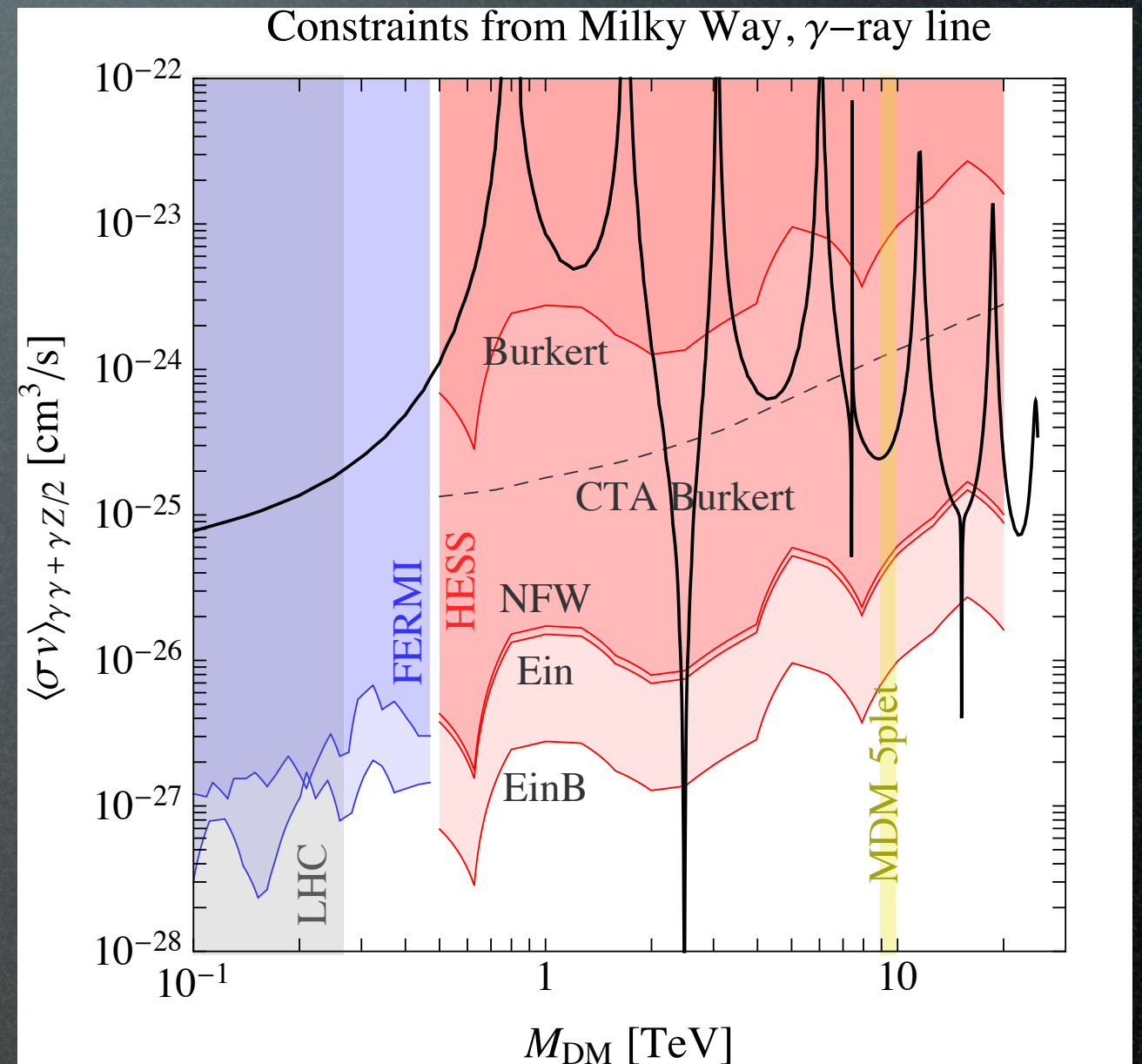
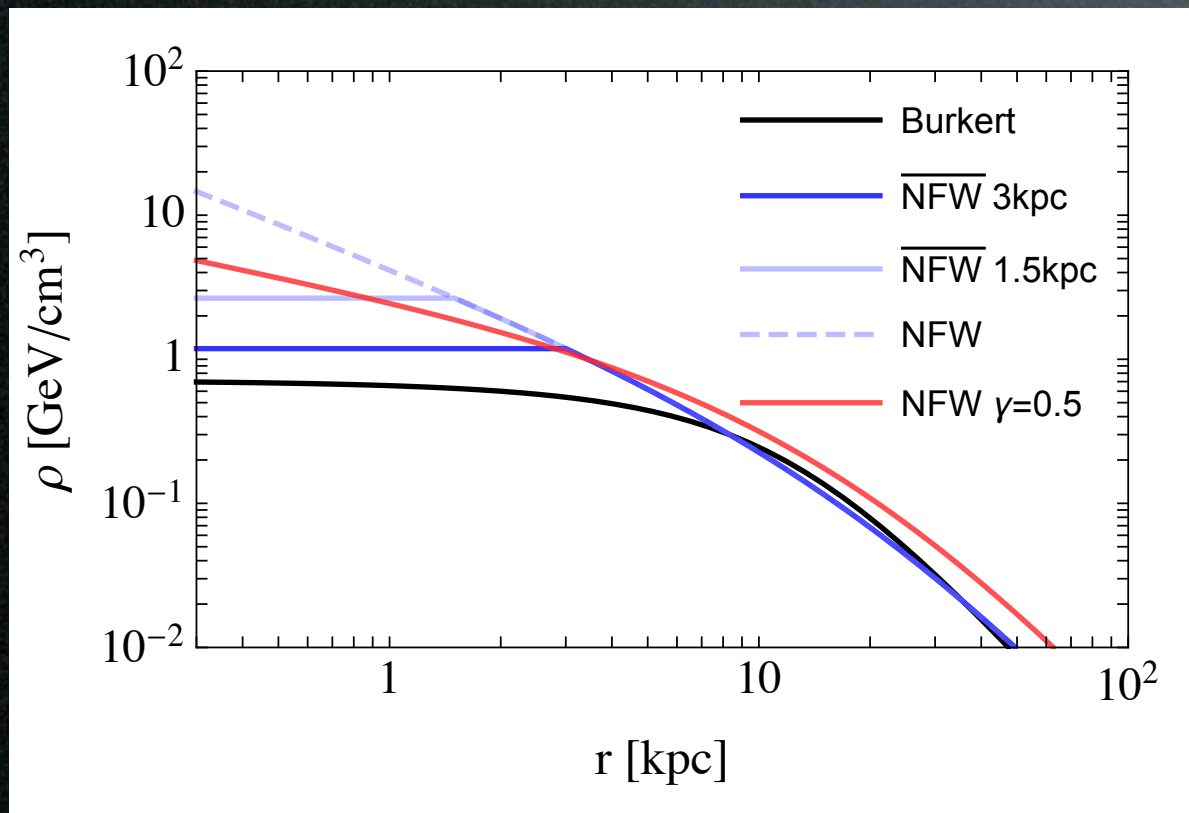
Simulations and observations
do not resolve $\lesssim 2$ kpc

3. Indirect Detection

MW center area, search for γ -ray lines:

Simulations and observations
do not resolve $\lesssim 2$ kpc

Uncertainties in DM profile:

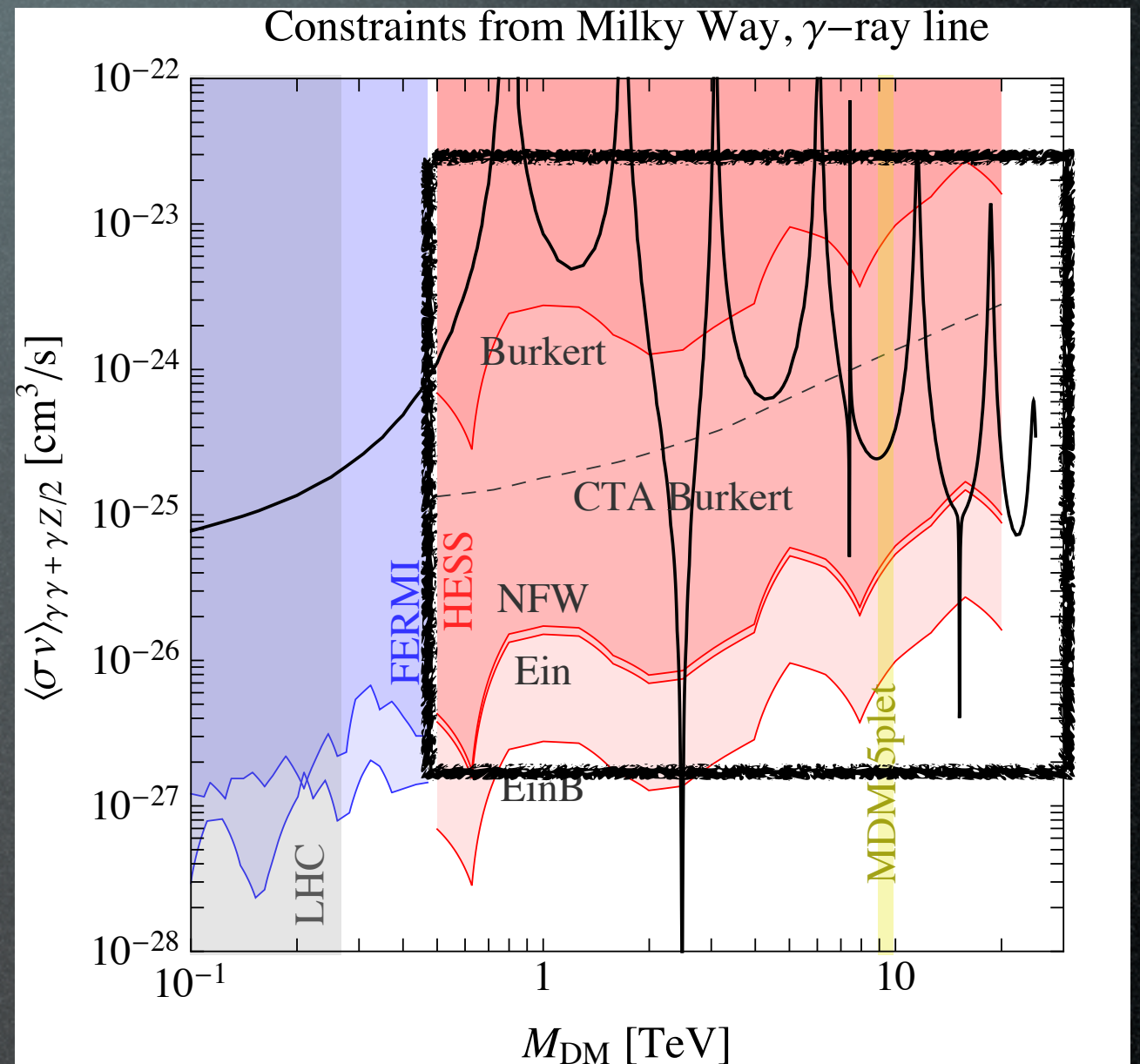
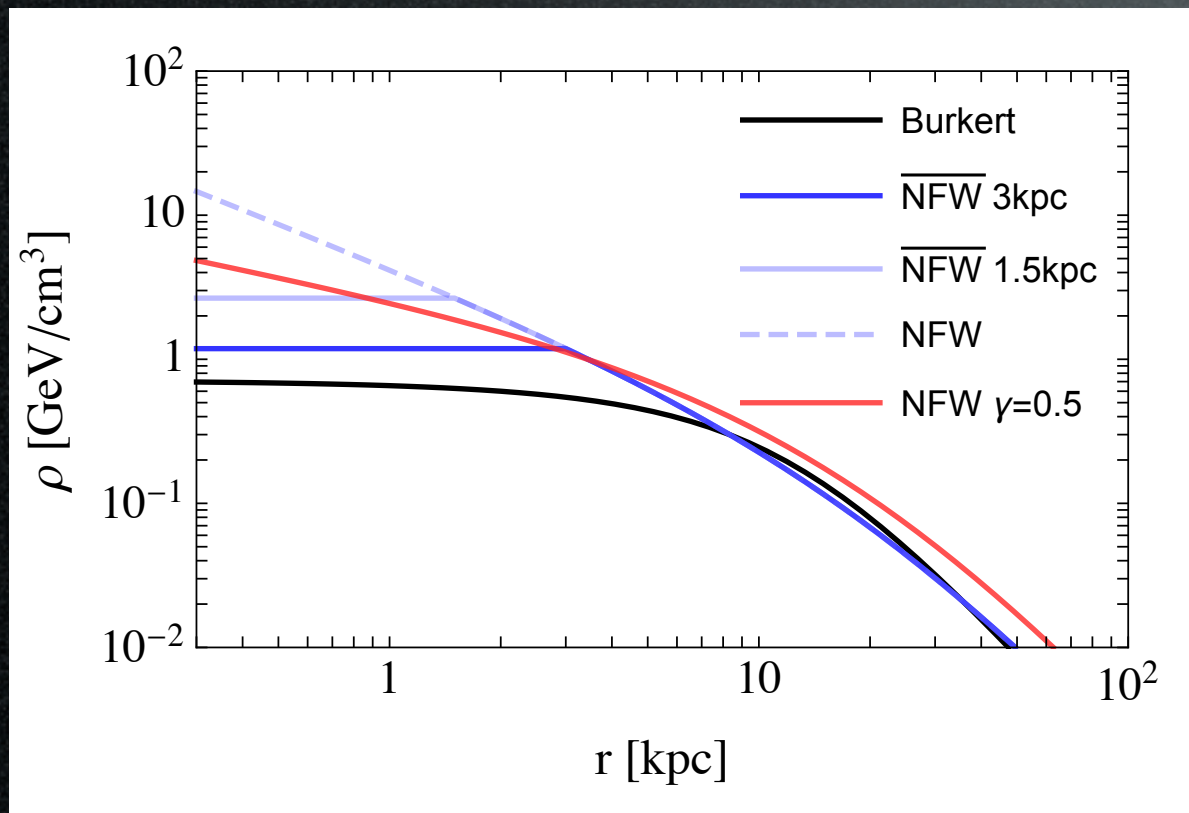


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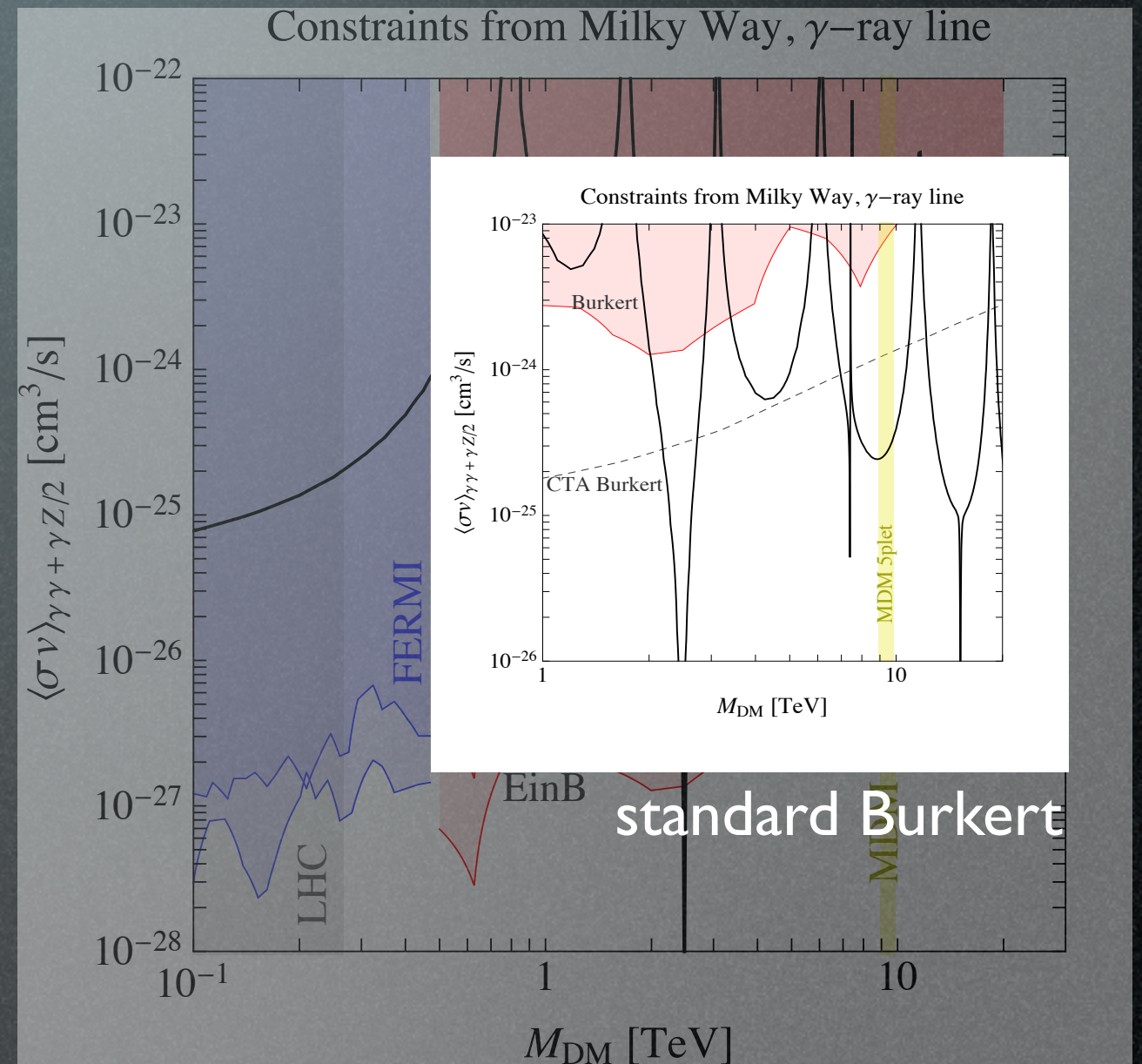
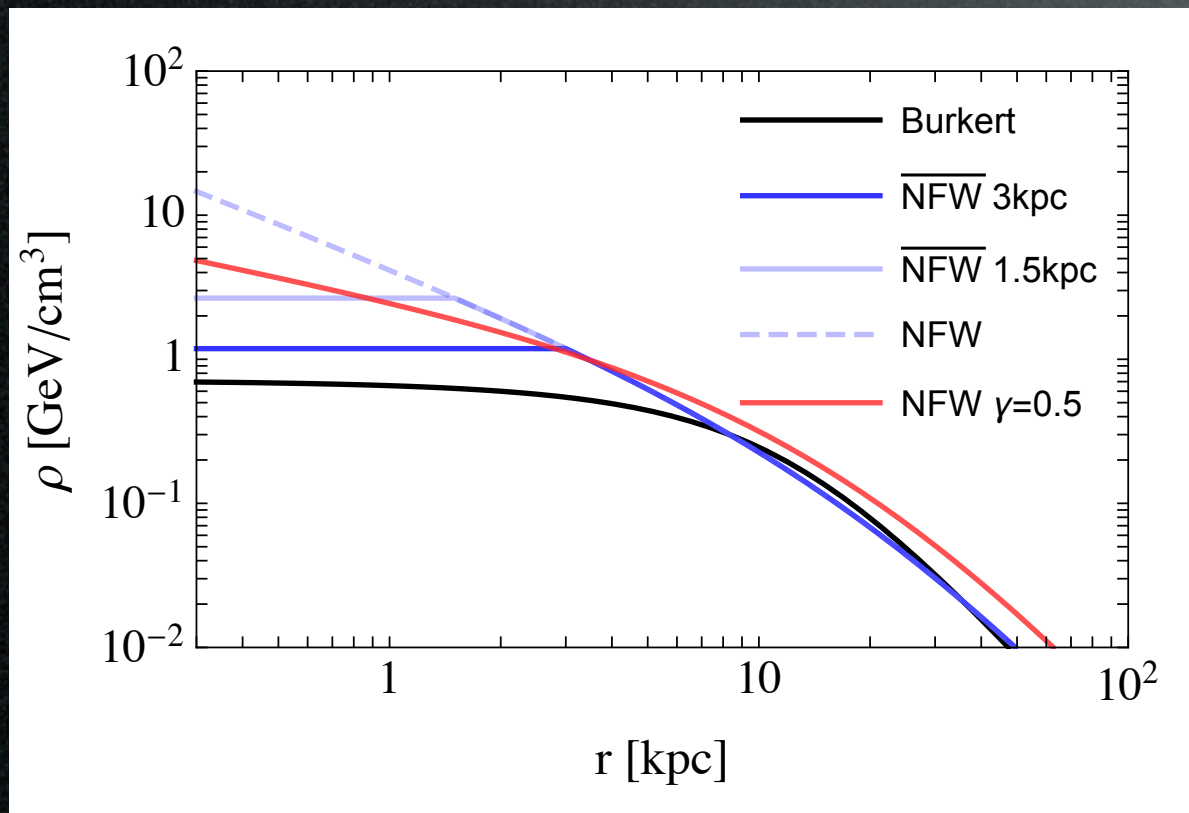


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Simulations and observations do not resolve $\lesssim 2$ kpc

Uncertainties in DM profile:

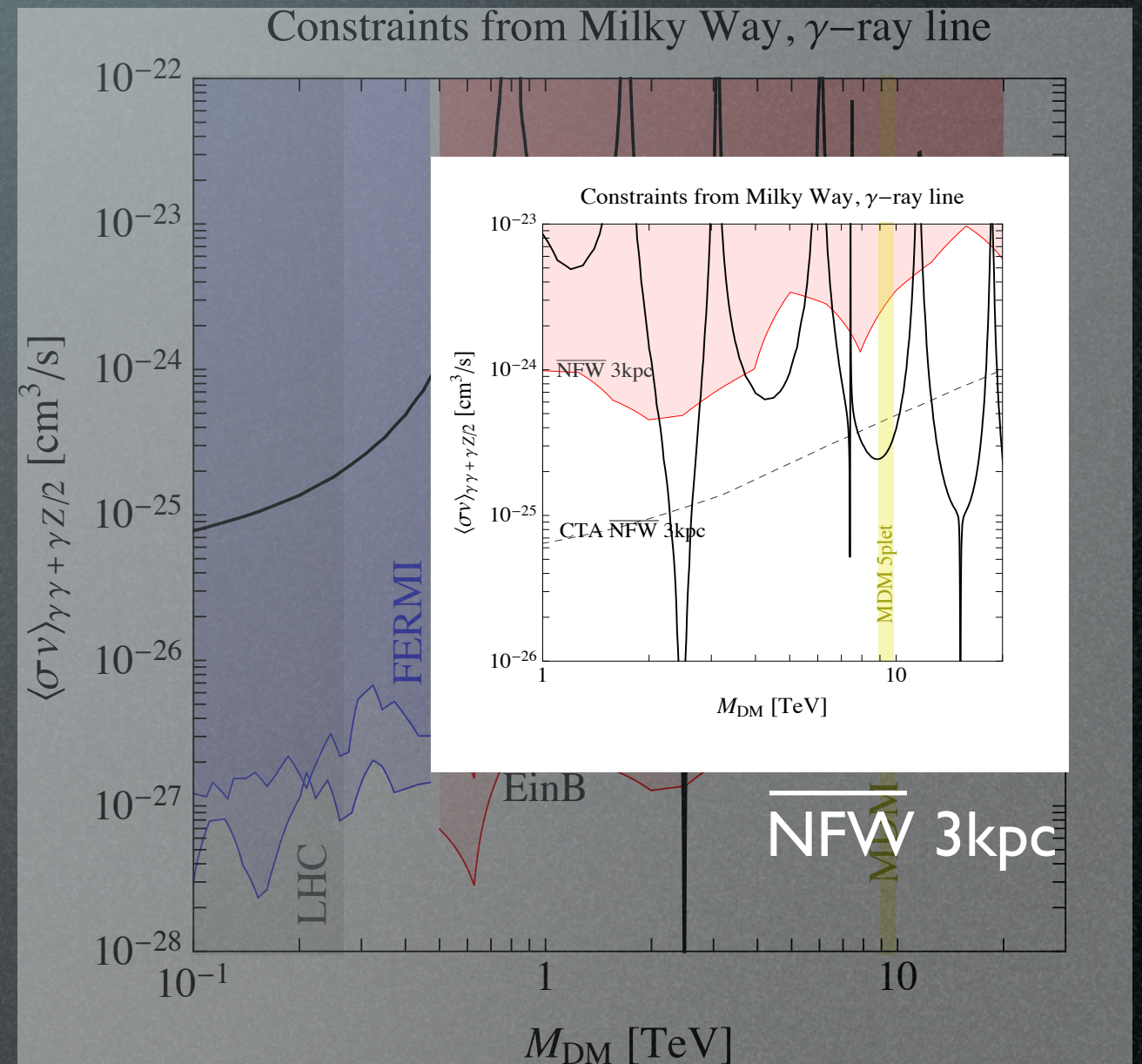
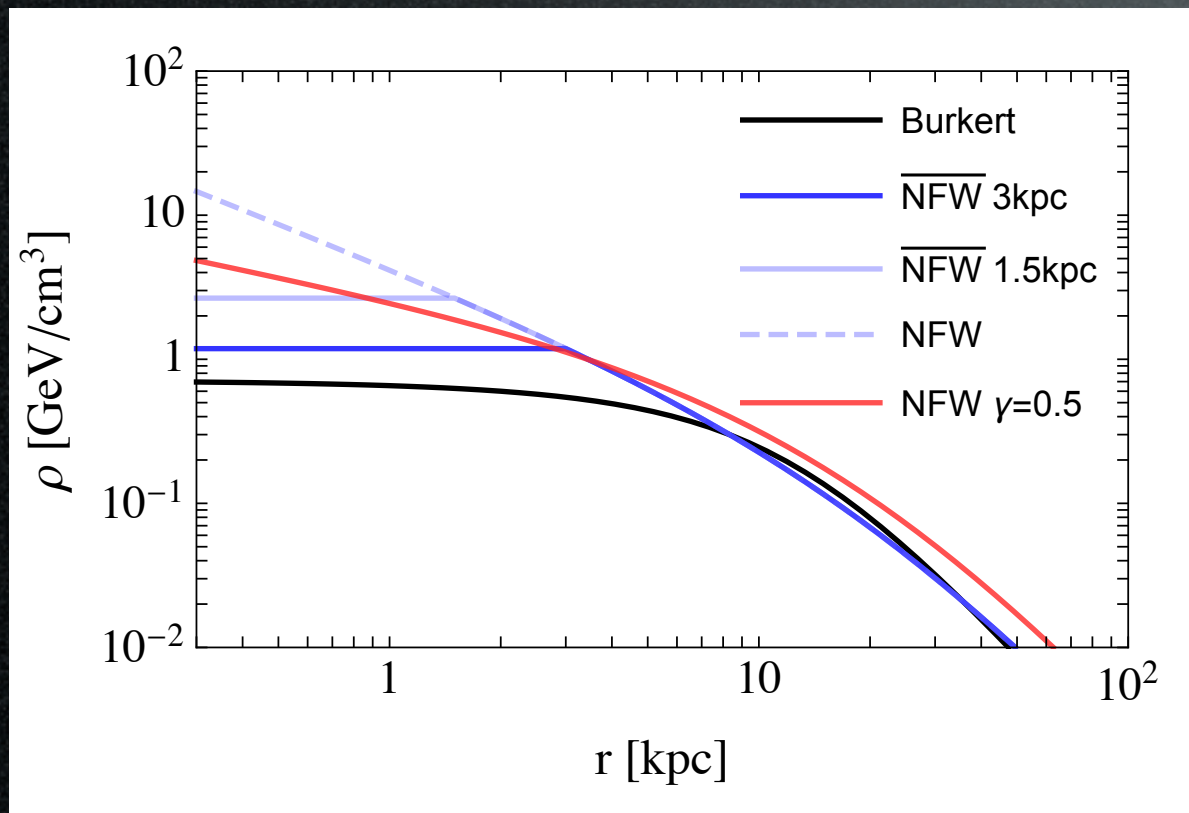


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Uncertainties in DM profile:

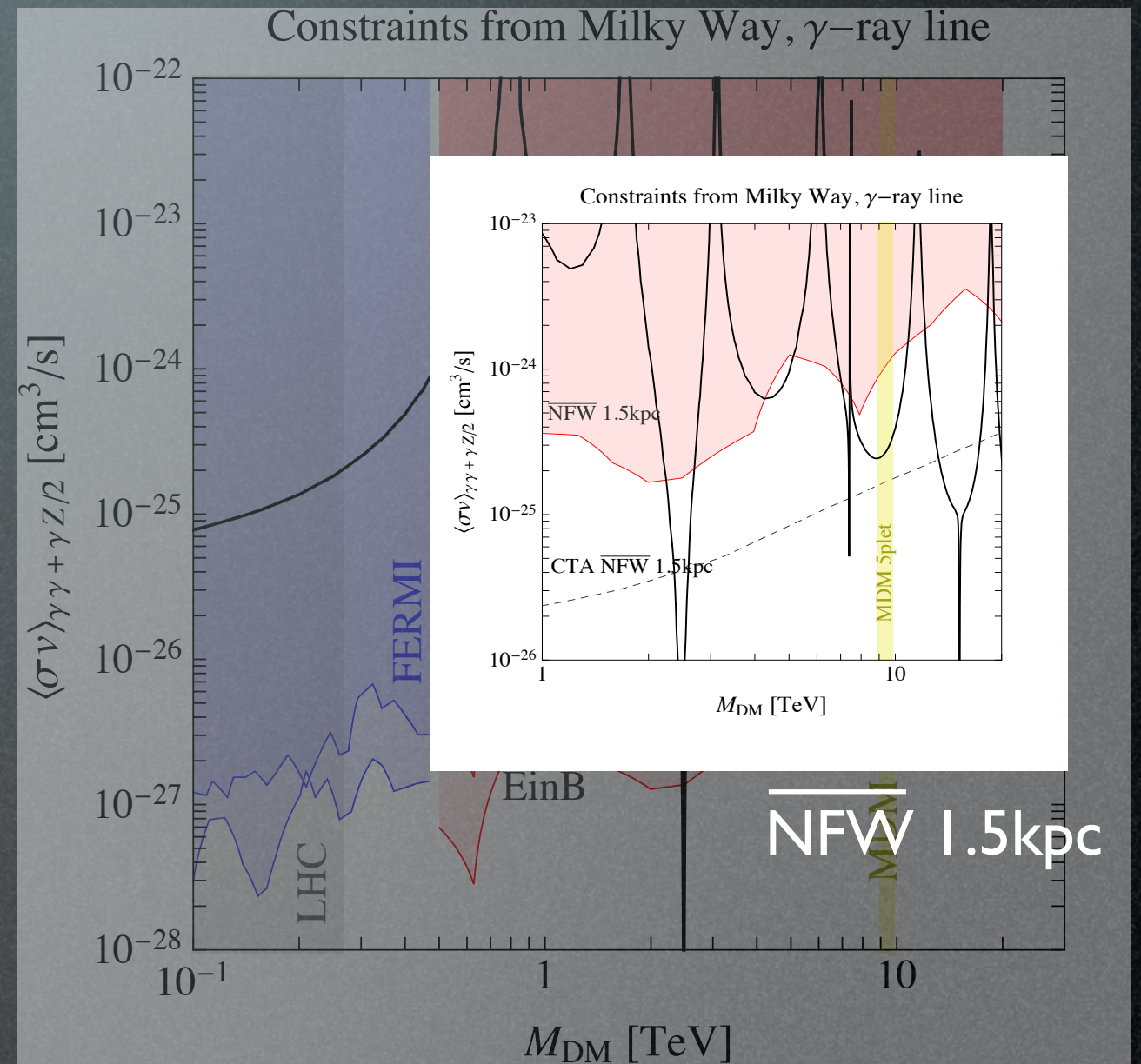
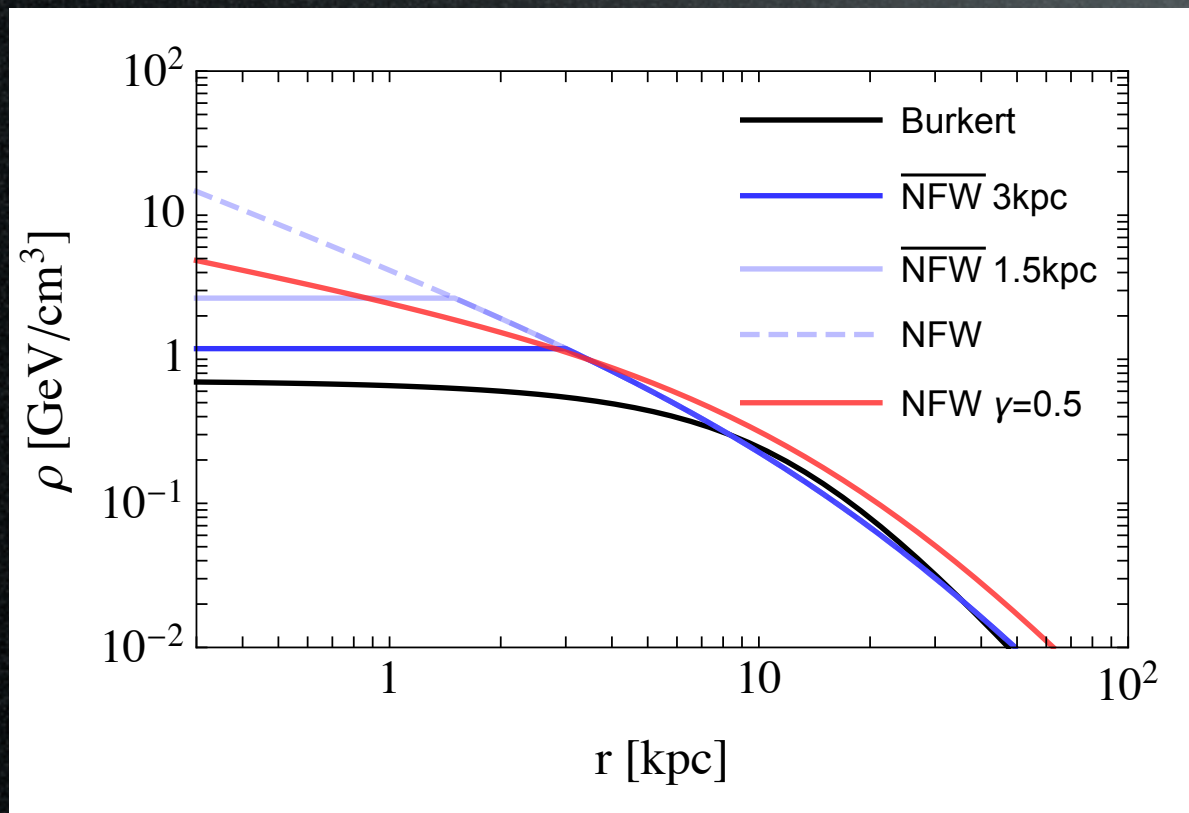


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Simulations and observations do not resolve $\lesssim 2$ kpc

Uncertainties in DM profile:

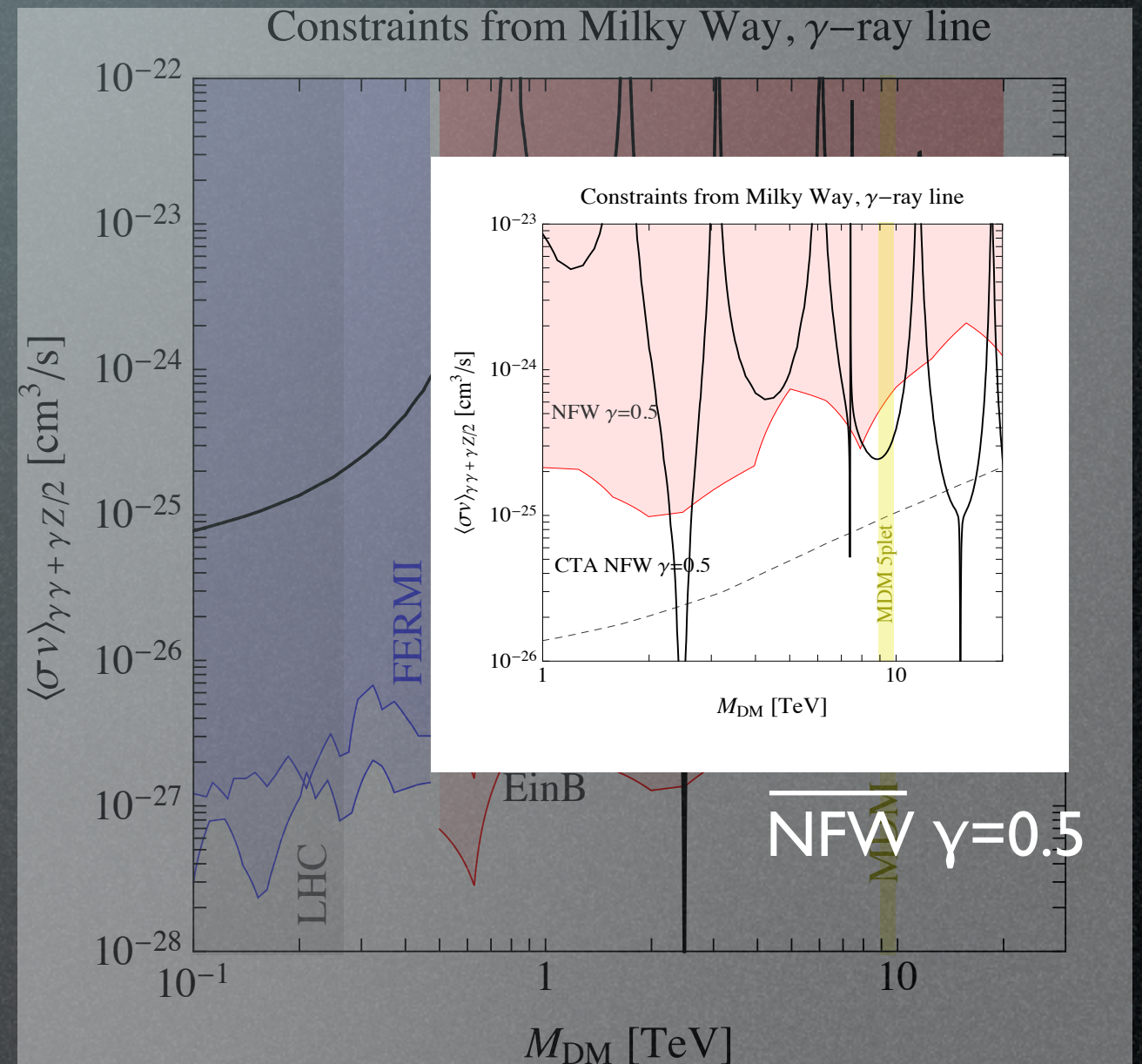
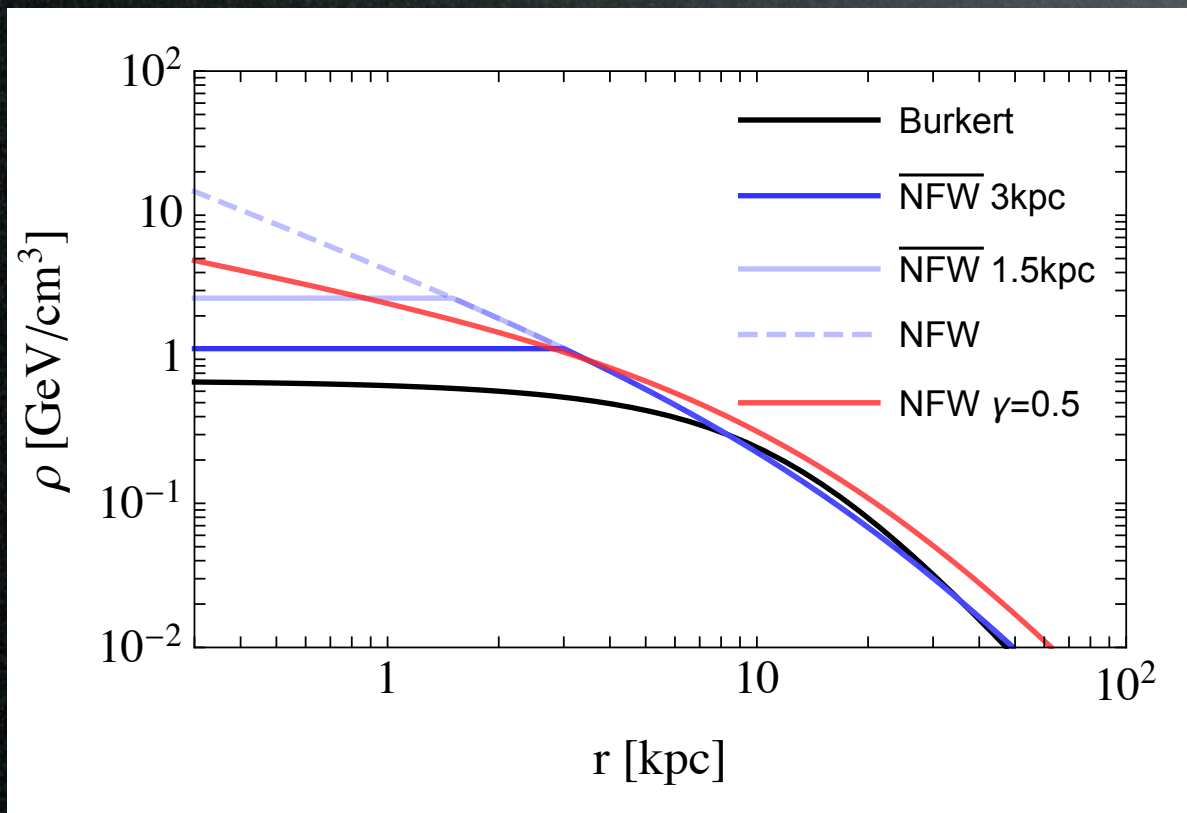


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MW center area, search for γ -ray lines:

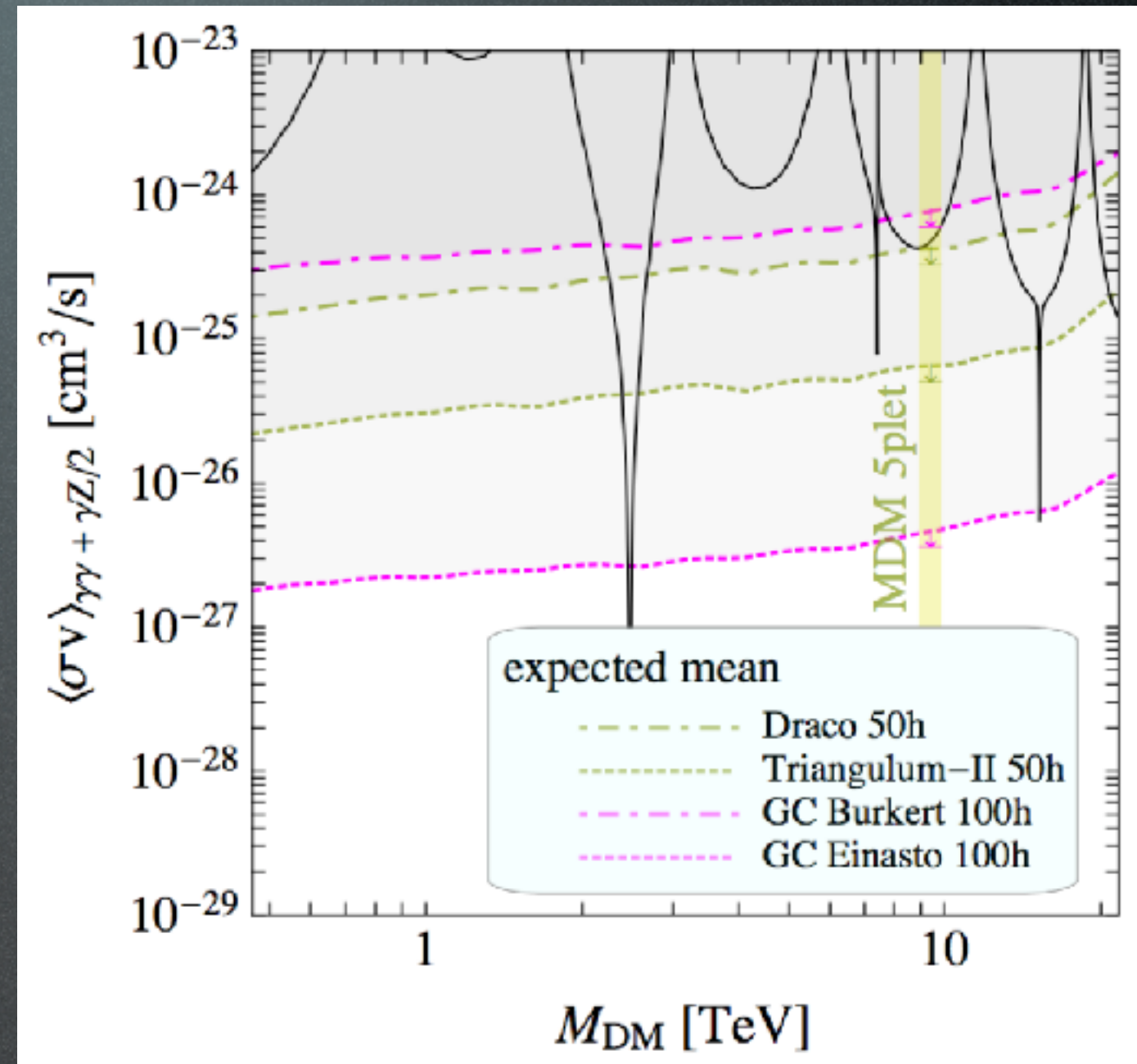
Simulations and observations do not resolve $\lesssim 2$ kpc

Uncertainties in DM profile:



3. Indirect Detection

Updated future **sensitivity** with **CTA**,
search for γ -ray **lines**:



Btw, dwarfs do better than GC
(if GC is cored).